# RESONANCES AS VISCOSITY LIMITS FOR EXTERIOR DILATION ANALYTIC POTENTIALS 

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#### Abstract

For exterior dilation analytic potential, $V$, we use the method of complex scaling to show that the resonances of $-\Delta+V$, in a conic neighbourhood of the real axis, are limits of eigenvalues of $-\Delta+V-i \varepsilon x^{2}$ as $\varepsilon \rightarrow 0+$, if $V$ can be analytically extended from $\mathbb{R}^{n}$ to a truncated cone in $\mathbb{C}^{n}$.


## 1. Introduction and statement of results

We extend the results of [Z2], when $V \in L_{\text {comp }}^{\infty}$, to the case of exterior dilation analytic potentials. For motivation and pointers to the literature we refer to [Z2].

Thus, we consider

$$
H:=-\Delta+V,
$$

where $V$ is a real-valued potential which can be analytically extended from $\left\{x \in \mathbb{R}^{n}\right.$ : $|x|>R\}$, for some $R>0$, to a truncated cone

$$
\mathcal{C}_{\beta_{0}}^{R}:=\left\{z \in \mathbb{C}^{n}:|\operatorname{Im} z|<\tan \beta_{0}|\operatorname{Re} z| \text { and }|\operatorname{Re} z|>R\right\}, \quad \beta_{0} \leq \pi / 8
$$

We still denote the analytic extension by $V$ and assume that

$$
\begin{equation*}
\lim _{\mathcal{C}_{\beta_{0}}^{R} \ni|z| \rightarrow \infty} V(z)=0 . \tag{1.1}
\end{equation*}
$$

The resonances of $H$ are defined by the Aguiliar-Balslev-Combes-Simon theory, see [HS, §16, §18], [DyZ2, §4.5] and a review in §3.

We now introduce a regularized operator,

$$
\begin{equation*}
H_{\varepsilon}:=-\Delta-i \varepsilon x^{2}+V, \quad \varepsilon>0 . \tag{1.2}
\end{equation*}
$$

(We write $x^{2}:=x_{1}^{2}+\cdots+x_{n}^{2}$.) It is easy to see, with details reviewed in $\S 4$, that $H_{\varepsilon}$ is a non-normal unbounded operator on $L^{2}\left(\mathbb{R}^{n}\right)$ with a discrete spectrum. We have

Theorem 1. Suppose that $\left\{z_{j}(\varepsilon)\right\}_{j=1}^{\infty}$ are the eigenvalues of $H_{\varepsilon}$. Then, uniformly on any compact subsets of $\left\{z:-2 \beta_{0}<\arg z<3 \pi / 2+2 \beta_{0}\right\}$,

$$
z_{j}(\varepsilon) \rightarrow z_{j}, \quad \varepsilon \rightarrow 0+
$$

where $z_{j}$ are the resonances of $H$.

Notation. We use the following notation: $f=\mathcal{O}_{\ell}(g)_{H}$ means that $\|f\|_{H} \leq C_{\ell} g$ where the norm (or any seminorm) is in the space $H$, and the constant $C_{\ell}$ depends on $\ell$. When either $\ell$ or $H$ are absent then the constant is universal or the estimate is scalar, respectively. When $G=\mathcal{O}_{\ell}(g): H_{1} \rightarrow H_{2}$ then the operator $G: H_{1} \rightarrow H_{2}$ has its norm bounded by $C_{\ell} g$. Also when no confusion is likely to result, we denote the operator $f \mapsto g f$ where $g$ is a function by $g$.

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## 2. SPECTRAL DEFORMATION AND ANALYTIC VECTORS

We will review several basic concepts in the Aguilar-Balslev-Combes-Simon theory, such as spectral deformation and analytic vectors. For a detailed introduction, we refer to $[H S, \S 17]$ and the references given there.

Let $h \in \mathcal{C}^{\infty}(\mathbb{R})$ be a non-decreasing function which satisfies

$$
\begin{cases}h(t)=0, & t<2 R  \tag{2.1}\\ h(t)=1, & t>8 R .\end{cases}
$$

Moreover, we assume that

$$
\begin{equation*}
\sup _{t \in \mathbb{R}} h(t)+t h^{\prime}(t) \leq 3 / 2 \tag{2.2}
\end{equation*}
$$

We define $g: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ as a smooth mapping by

$$
g(x):=h(|x|) x= \begin{cases}0, & |x|<2 R  \tag{2.3}\\ x, & |x|>8 R\end{cases}
$$

and consider, for $\theta \in \mathbb{R}$, the related family of maps $\phi_{\theta}: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ defined by

$$
\begin{equation*}
\phi_{\theta}(x)=x+\theta g(x) \tag{2.4}
\end{equation*}
$$

We let $D f$ denote the derivative of a map $f: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$, then

$$
D g(x)=h(|x|) I+|x|^{-1} h^{\prime}(|x|) x \cdot x^{T} .
$$

Using diagonalization, It is easy to see that

$$
\begin{equation*}
0 \leq h(|x|) I \leq D g(x) \leq\left(h(|x|)+|x| h^{\prime}(|x|)\right) I \leq 3 / 2 I \tag{2.5}
\end{equation*}
$$

where $A \leq B$ means $B-A$ is positive semi-definite and the last inequality is implied by (2.2). Hence $\sup _{x \in \mathbb{R}^{n}}\|D g(x)\| \leq 3 / 2$, where $\|\cdot\|$ denotes the operator norm on the
set of linear transformation on $\mathbb{R}^{n}$. We note that $D \phi_{\theta}(x)=I+\theta(D g)(x)$, if $|\theta|<2 / 3$, then $D \phi_{\theta}$ is invertible by a Neumann series argument,

$$
\left(D \phi_{\theta}\right)^{-1}=\sum_{j=0}^{\infty}(-1)^{j} \theta^{j}(D g)^{j}
$$

Hence $\phi_{\theta}$ is a diffeomorphism of $\mathbb{R}^{n}$ for $|\theta|<2 / 3$ by the inverse function theorem.
We should remark that all the above argument is valid when we extend the definition (2.4) of $\phi_{\theta}$ to $\theta \in \mathbb{C}$. We have $\phi_{\theta}: \mathbb{R}^{n} \rightarrow \mathbb{C}^{n}$ is a diffeomorphism provided $|\theta|<2 / 3$.

We now introduce the behavior of functions under the action of the maps $\phi_{\theta}$. We first define $U_{\theta}$ for $\theta \in \mathbb{R}$ by

$$
\begin{equation*}
\left(U_{\theta} f\right)(x)=J_{\theta}(x)^{1 / 2} f\left(\phi_{\theta}(x)\right) \tag{2.6}
\end{equation*}
$$

where $J_{\theta}(x)$ is the Jacobian of $\phi_{\theta}$,

$$
\begin{equation*}
J_{\theta}(x)=\operatorname{det} D \phi_{\theta}(x)=\operatorname{det}(I+\theta(D g)(x)) \tag{2.7}
\end{equation*}
$$

It is east to see that $U_{\theta}, \theta \in \mathbb{R}$ is unitary on $L^{2}\left(\mathbb{R}^{n}\right)$ with the inverse $U_{\theta}^{-1}$ given by

$$
\begin{equation*}
\left(U_{\theta}^{-1} f\right)(x)=J_{\theta}\left(\phi_{\theta}^{-1}(x)\right)^{-1 / 2} f\left(\phi_{\theta}^{-1}(x)\right) \tag{2.8}
\end{equation*}
$$

(2.5) and (2.7) show that $J_{\theta}(x)^{1 / 2}$ extends analytically to complex $\theta$ provided $\theta<2 / 3$. Hence, to extend the operators $U_{\theta}$ from $\theta \in \mathbb{R}$ to $\theta \in \mathbb{C}$, at least for small $|\theta|$, we need to find a dense set of functions $f$ in $L^{2}\left(\mathbb{R}^{n}\right)$ that can be analytically extended on a small complex neighborhood of $\mathbb{R}^{n}$ in $\mathbb{C}^{n}$ such that $f \circ \phi_{\theta} \in L^{2}\left(\mathbb{R}^{n}\right)$. For that we introduce the set of analytic vectors in $L^{2}\left(\mathbb{R}^{n}\right)$.

Definition 1. Let $\mathcal{A}$ be the linear space of all entire functions $f(z)$ having the property that in any conical region $\mathcal{C}_{\varepsilon}$,

$$
\mathcal{C}_{\varepsilon}:=\left\{z \in \mathbb{C}^{n}:|\operatorname{Im} z| \leq(1-\varepsilon) \operatorname{Re} z\right\}
$$

for any $\varepsilon>0$, we have for any $k \in \mathbb{N}$,

$$
\lim _{z \in \mathcal{C}_{\varepsilon} \rightarrow \infty}|z|^{k}|f(z)|=0
$$

The set of analytic vectors in $L^{2}\left(\mathbb{R}^{n}\right)$ are the restrictions to $\mathbb{R}^{n}$ of $\mathcal{A}$, which is also denoted by $\mathcal{A}$.

We define a domain $D_{\beta_{0}}$ in $\mathbb{C}$ by

$$
\begin{equation*}
D_{\beta_{0}}=\left\{\theta \in \mathbb{C}:|\operatorname{Re} \theta|+|\operatorname{Im} \theta|<\tan \beta_{0}\right\} \tag{2.9}
\end{equation*}
$$

Note that $D_{\beta_{0}} \subset\{z \in \mathbb{C}:|z|<1 / 2\}$ since $\beta_{0} \leq \pi / 8$, (2.5) and (2.7) guarantee that the Jacobian $J_{\theta}$ is uniformly bounded for $\theta \in D_{\beta_{0}}$. Then, we recall the following results in [HS, Proposition 17.10]:

Proposition 2. Let $\mathcal{U} \equiv\left\{U_{\theta}: \theta \in D_{\beta_{0}}\right\}$ be a spectral deformation family associated with vector field $g$ defined by (2.3). Then,

- the $\operatorname{map}(\theta, f) \in D_{\beta_{0}} \times \mathcal{A} \rightarrow U_{\theta} f$ is an $L^{2}$-analytic map ;
- for any $\theta \in D_{\beta_{0}}, U_{\theta} \mathcal{A}$ is dense in $L^{2}\left(\mathbb{R}^{n}\right)$.

We conclude this section with some properties about the deformation of $\mathbb{R}^{n} \subset \mathbb{C}^{n}$ under the map $\phi_{\theta}$ provided $\theta \in D_{\beta_{0}}$. We recall that $\phi_{\theta}: \mathbb{R}^{n} \rightarrow \mathbb{C}^{n}$ is injective with the Jacobian $J_{\theta} \neq 0$ provided $|\theta|<2 / 3$. Hence $\phi_{\theta}\left(\mathbb{R}^{n}\right) \subset \mathbb{C}^{n}$ is an $n$-dimensional totally real submanifolds, see $[\mathrm{DyZ2}, \S 4.5]$. Let $\Gamma_{a(\theta)}=\phi_{\theta}\left(\mathbb{R}^{n}\right)$, where $a(\theta) \in(-\pi / 2, \pi / 2)$ is defined by

$$
\begin{equation*}
a(\theta)=\arg (1+\theta) . \tag{2.10}
\end{equation*}
$$

In the literature about complex scaling, one can define $L^{2}\left(\Gamma_{a(\theta)}\right)$ with volume element $|d w|=\left|J_{\theta}(x)\right| d x$ where $w=\phi_{\theta}(x)$ are the coordinates on $\Gamma_{a(\theta)}$, see [DyZ2, §2.7, §4.5] for details. Then we have the following:

Proposition 3. For any $\theta \in D_{\beta_{0}}, \Gamma_{a(\theta)}$ satisfies

$$
\begin{gather*}
\Gamma_{a(\theta)} \cap B_{\mathbb{C}^{n}}(0,2 R)=B_{\mathbb{R}^{n}}(0,2 R), \\
\Gamma_{a(\theta)} \cap \mathbb{C}^{n} \backslash B_{\mathbb{C}^{n}}(0,12 R)=e^{i a(\theta)} \mathbb{R}^{n} \cap \mathbb{C}^{n} \backslash B_{\mathbb{C}^{n}}(0,12 R),  \tag{2.11}\\
\Gamma_{a(\theta)} \subset \mathbb{R}^{n} \cup \mathcal{C}_{\beta_{0}}^{R}
\end{gather*}
$$

Furthermore, the spectral deformation operator $U_{\theta}$ extends to an isometry:

$$
U_{\theta}: L^{2}\left(\Gamma_{a(\theta)}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)
$$

Proof. In view of (2.3) and (2.4), it is easy to see that $\Gamma_{a(\theta)}=\phi_{\theta}\left(\mathbb{R}^{n}\right)$ satisfies the first two equations of (2.11). For $\theta \in D_{\beta_{0}}$, we have

$$
\frac{\left|\operatorname{Im} \phi_{\theta}(x)\right|}{\left|\operatorname{Re} \phi_{\theta}(x)\right|}=\frac{|\operatorname{Im} \theta||\chi(|x|)|}{|1+\operatorname{Re} \theta \chi(|x|)|} \leq \frac{|\operatorname{Im} \theta|}{1-|\operatorname{Re} \theta|}<\tan \beta_{0}
$$

where the last inequality is implied by (2.9). Moreover, $\phi_{\theta}(x)=x$ for $|x|<2 R$, and $\left|\operatorname{Re} \phi_{\theta}(x)\right| \geq(1-|\operatorname{Re} \theta|)|x|>\left(1-\tan \beta_{0}\right)|x|>|x| / 2 \geq R$ provided $|x| \geq 2 R$, since $\beta_{0} \leq \pi / 8$. Hence $\Gamma_{a(\theta)} \subset \mathbb{R}^{n} \cup \mathcal{C}_{\beta_{0}}^{R}$.

Now we assume that $\theta \in D_{\beta_{0}}$, for any $f \in L^{2}\left(\Gamma_{a(\theta)}\right)$, we can define $U_{\theta} f$ on $\mathbb{R}^{n}$ by (2.6). To see $U_{\theta} f \in L^{2}\left(\mathbb{R}^{n}\right)$, we compute directly:

$$
\begin{align*}
\int_{\mathbb{R}^{n}}\left|U_{\theta} f(x)\right|^{2} d x & =\int_{\mathbb{R}^{n}}\left|J_{\theta}(x)^{1 / 2} f\left(\phi_{\theta}(x)\right)\right|^{2} d x \\
& =\int_{\mathbb{R}^{n}}\left|f\left(\phi_{\theta}(x)\right)\right|^{2}\left|J_{\theta}(x)\right| d x  \tag{2.12}\\
& =\int_{\Gamma_{a(\theta)}}|f(w)|^{2}|d w|<\infty
\end{align*}
$$

which also shows that $\left\|U_{\theta} f\right\|_{L^{2}\left(\mathbb{R}^{n}\right)}=\|f\|_{L^{2}\left(\Gamma_{a(\theta)}\right)}$ and $U_{\theta}$ is one-to-one. It remains to show that $U_{\theta}$ is onto. For $g \in L^{2}\left(\mathbb{R}^{n}\right)$, let $G(w)=J_{\theta}\left(\phi_{\theta}^{-1}(w)\right)^{-1 / 2} g\left(\phi_{\theta}^{-1}(w)\right), w \in \Gamma_{a(\theta)}$. We can follow (2.12) to derive

$$
\int_{\Gamma_{a(\theta)}}|G(w)|^{2}|d w|=\int_{\mathbb{R}^{n}}|g(x)|^{2} d x
$$

Hence $G \in L^{2}\left(\Gamma_{a(\theta)}\right)$, then we conclude that $U_{\theta}$ is onto since $U_{\theta} G=g$.

## 3. RESONANCES

We will follow Aguilar-Balslev-Combes-Simon theory to define the resonances of $H \equiv-\Delta+V$, see $[H S, \S 16, \S 18]$ and those resonances in a conic neighborhood of the real axis can be identified with the eigenvalues of certain non-self-adjoint operators associated with $H$. Using the analytic vectors $\mathcal{A}$, we recall the definition:

Definition 4. The resonances of $H$ associated with analytic vectors $\mathcal{A}$ are the poles of the meromorphic continuations of all matrix elements $\left\langle f, R_{H}(z) g\right\rangle\left(R_{H}(z)\right.$ denotes the resolvent of $H), f, g \in \mathcal{A}$, from $\{z \in \mathbb{C}: \operatorname{Im} z>0\}$ to $\{z \in \mathbb{C}: \operatorname{Im} z \leq 0\}$.

First we introduce the spectral deformed Schrödinger operators $H(\theta)$ of $H$ associated with the spectral deformation family $\mathcal{U}=\left\{U_{\theta}: \theta \in D_{\beta_{0}}\right\}$. Consider, for $\theta \in D_{\beta_{0}} \cap \mathbb{R}$,

$$
\begin{equation*}
H(\theta):=U_{\theta} H U_{\theta}^{-1}=p_{\theta}^{2}+V\left(\phi_{\theta}(x)\right) \tag{3.1}
\end{equation*}
$$

where

$$
\begin{equation*}
p_{\theta}^{2}=U_{\theta} p^{2} U_{\theta}^{-1}, \quad p_{j} \equiv \frac{1}{i} \frac{\partial}{\partial x_{j}} . \tag{3.2}
\end{equation*}
$$

In view of Proposition 3, we can extend $H(\theta)$ to $\theta \in D_{\beta_{0}}$. We recall the following basic facts about $p_{\theta}^{2}, \theta \in D_{\beta_{0}}$ in [HS, §18]:

Proposition 5. Let $p_{\theta}^{2}$ be as defined in (3.2), then $p_{\theta}^{2}, \theta \in D_{\beta_{0}}$ is an analytic family of operators with domain $D\left(p_{\theta}^{2}\right)=H^{2}\left(\mathbb{R}^{n}\right)$. For the spectrum, we have $\sigma\left(p_{\theta}^{2}\right)=\sigma_{\text {ess }}\left(p_{\theta}^{2}\right)=$ $e^{-2 i a(\theta)}[0, \infty)$.

And for the resolvent $R_{\theta}(z):=\left(p_{\theta}^{2}-z\right)^{-1}$ we have:
Proposition 6. For $\delta>0$, we have

$$
\begin{equation*}
R_{\theta}(z)=\mathcal{O}_{\delta}(1 /|z|): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right), \quad-2 a(\theta)+\delta<\arg z<2 \pi-2 a(\theta)-\delta \tag{3.3}
\end{equation*}
$$

Proof. We note that in the notation of Proposition 3,

$$
\begin{equation*}
p_{\theta}^{2}=U_{\theta}\left(-\Delta_{a(\theta)}\right) U_{\theta}^{-1}, \tag{3.4}
\end{equation*}
$$

where $-\Delta_{a(\theta)}: H^{2}\left(\Gamma_{a(\theta)}\right) \rightarrow L^{2}\left(\Gamma_{a(\theta)}\right)$ is defined as the restriction of $\Delta_{z}$ to the totally real submanifold $\Gamma_{a(\theta)}$, see [DyZ2, §4.5]. Since $U_{\theta}: L^{2}\left(\Gamma_{a(\theta)}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ and $U_{\theta}^{-1}$ : $L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\Gamma_{a(\theta)}\right)$ are both isometries, we have

$$
\left\|\left(p_{\theta}^{2}-z\right)^{-1}\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)}=\left\|\left(-\Delta_{a(\theta)}-z\right)^{-1}\right\|_{L^{2}\left(\Gamma_{a(\theta)}\right) \rightarrow L^{2}\left(\Gamma_{a(\theta)}\right)}, \quad z \notin e^{-2 i a(\theta)}[0, \infty)
$$

and thus (3.3) is a direct consequence of [DyZ2, Theorem 4.35].
Then we introduce some preliminary properties of the spectrum of $H(\theta)$ :
Proposition 7. There exists $R>0$ such that for any $\theta \in D_{\beta_{0}}$, we have

$$
\sigma(H(\theta)) \cap i(R, \infty)=\emptyset
$$

As for the essential spectrum $\sigma_{\text {ess }}(H(\theta))$, we have more precisely,

$$
\sigma_{e s s}(H(\theta))=e^{-2 i a(\theta)}[0, \infty)
$$

Remark: In fact, $\sigma(H(\theta)) \cap\{z: 0<\arg z<2 \pi-2 a(\theta)\}$ is discrete and lies in $(-\infty, 0)$, which is a consequence of the following Lemma 1.

Proof. For $\theta \in D_{\beta_{0}}$, we have

$$
\begin{equation*}
(H(\theta)-z) R_{\theta}(z)=I+V\left(\phi_{\theta}(x)\right) R_{\theta}(z) \tag{3.5}
\end{equation*}
$$

Now assume $z \in i(R, \infty)$, note that $\theta \in D_{\beta_{0}} \Longrightarrow-\beta_{0}<a(\theta)<\beta_{0}$, we have

$$
-2 a(\theta)+\pi / 4<\arg z<2 \pi-2 a(\theta)-\pi / 4, \text { for all } \theta \in D_{\beta_{0}} .
$$

Using (3.3), we see that $R_{\theta}(z)=\mathcal{O}(1 /|z|): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ for all $z \in i(R, \infty)$ and $\theta \in D_{\beta_{0}}$. Recalling $\phi_{\theta}\left(\mathbb{R}^{n}\right) \subset \mathbb{R}^{n} \cup C_{\beta_{0}}^{R}$ and $V \in L^{\infty}\left(\mathbb{R}^{n} \cup C_{\beta_{0}}^{R}\right)$, we conclude that

$$
\sup _{z \in i(R, \infty)}\left\|V\left(\phi_{\theta}(x)\right) R_{\theta}(z)\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)}=\mathcal{O}\left(R^{-1}\right), \quad \text { uniformly for } \theta \in D_{\beta_{0}}
$$

Then for $R \gg 1, I+V\left(\phi_{\theta}(x)\right) R_{\theta}(z)$ is invertible using the Neumann series:

$$
\left(I+V\left(\phi_{\theta}(x)\right) R_{\theta}(z)\right)^{-1}=\sum_{j=0}^{\infty}\left(V\left(\phi_{\theta}(x)\right) R_{\theta}(z)\right)^{j}
$$

Hence $H(\theta)-z$ is invertible by (3.5), for all $z \in i(R, \infty)$.
For the essential spectrum $\sigma_{\text {ess }}(H(\theta))$, note that $\sigma_{\text {ess }}\left(p_{\theta}^{2}\right)=e^{-2 i a(\theta)}[0, \infty)$ in Proposition 5 , by the invariance under compact perturbations, it suffices to show that $V\left(\phi_{\theta}(x)\right)$ is $p_{\theta}^{2}$-compact, i.e. $V\left(\phi_{\theta}(x)\right): D\left(p_{\theta}^{2}\right)=H^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ is compact. Since $H^{2}\left(B_{\mathbb{R}^{n}}(0, R)\right) \Subset L^{2}\left(B_{\mathbb{R}^{n}}(0, R)\right), \forall R>0$, and $V \circ \phi_{\theta} \in L^{\infty}\left(\mathbb{R}^{n}\right), \quad V\left(\phi_{\theta}(x)\right) \rightarrow 0, x \rightarrow$ $\infty$ by (1.1), it is easy to see the compactness of $V\left(\phi_{\theta}(x)\right): H^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$.

Now we state the main result in this section, in which we identify the resonances defined in Definition 4 as the eigenvalues of certain spectral deformed operators $H(\theta)$.

Lemma 1. Let $H=-\Delta+V$ be a self-adjoint Schrödinger operator with a real-valued potential $V$ satisfying our assumptions as in $\S 1$. Then for any $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$, we have:

- For $f, g \in \mathcal{A}$, the function

$$
\begin{equation*}
F_{f, g}(z) \equiv\left\langle f, R_{H}(z) g\right\rangle \tag{3.6}
\end{equation*}
$$

defined for $\operatorname{Im} z>0$, has a meromorphic continuation across $[0, \infty)$ into $S_{\theta}^{-} \equiv$ $\mathbb{C} \backslash e^{-2 i a(\theta)}[0, \infty)$.

- The poles of the meromorphic continuations of all matrix elements $F_{f, g}(z)$ into $S_{\theta}^{-}$are eigenvalues of the operator $H(\theta)$.

Proof. With $F_{f, g}(z)$ defined in (3.6), the assumption on $V$ implies that $F_{f, g}$ is analytic on $\mathbb{C}^{+} \equiv\{z \in \mathbb{C}: \operatorname{Im} z>0\}$. Fix $z \in \mathbb{C}^{+}$. For $\theta \in D_{\beta_{0}} \cap \mathbb{R}, U_{\theta}$ is unitary and thus we can write

$$
\begin{equation*}
F_{f, g}(z)=\left\langle U_{\theta} f,\left(U_{\theta} R_{H}(z) U_{\theta}^{-1}\right) U_{\theta} g\right\rangle=\left\langle U_{\theta} f, R_{H(\theta)}(z) U_{\theta} g\right\rangle . \tag{3.7}
\end{equation*}
$$

Proposition 7 implies that $\theta \in D_{\beta_{0}} \rightarrow R_{H(\theta)}(z)$ is an analytic map provided $z \in$ $i(R, \infty)$. Since we can write $U_{\bar{\theta}} f$ instead of $U_{\theta} f$ in (3.7), we have

$$
\begin{equation*}
\theta \in D_{\beta_{0}} \rightarrow F_{f, g}(z ; \theta) \equiv\left\langle U_{\bar{\theta}} f, R_{H(\theta)}(z) U_{\theta} g\right\rangle \tag{3.8}
\end{equation*}
$$

is an analytic map provided $z \in i(R, \infty)$. Hence for any $z \in i(R, \infty)$, we have

$$
F_{f, g}(z ; \theta)=F_{f, g}(z), \quad \forall \theta \in D_{\beta_{0}},
$$

since this is true for all $\theta \in D_{\beta_{0}} \cap \mathbb{R}$. Now fix any $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$, Proposition 7 guarantees that $F_{f, g}(z ; \theta)$ can be meromorphically continued from $i(R, \infty)$ to $S_{\theta}^{-}$since $\sigma_{\text {ess }}(H(\theta)) \cap S_{\theta}^{-}=\emptyset$. We have shown that $F_{f, g}(z ; \theta)=F_{f, g}(z), z \in i(R, \infty)$, then by the identity principle for meromorphic functions, we conclude that $F_{f, g}(z ; \theta)$ is a meromorphic continuation of $F_{f, g}(z)$ from $\mathbb{C}^{+}$to $S_{\theta}^{-}$.

Recalling that $F_{f, g}(z ; \theta)=\left\langle U_{\bar{\theta}} f,(H(\theta)-z)^{-1} U_{\theta} g\right\rangle$ and that $U_{\theta} \mathcal{A}, U_{\bar{\theta}} \mathcal{A}$ are both dense in $L^{2}\left(\mathbb{R}^{n}\right)$, thus if $H(\theta)$ has an eigenvalue at $\lambda_{\theta} \in S_{\theta}^{-}$, there must exist $f, g \in \mathcal{A}$ such that $\lambda_{\theta}$ is a pole of $F_{f, g}(z ; \theta)$. Conversely, if $F_{f, g}(z ; \theta)$ has a pole $\lambda_{\theta} \in S_{\theta}^{-}$, then it must be an eigenvalue of $H(\theta)$.

Remark: For nonzero resonance $\lambda$ of $H$, we can define its multiplicity as the (algebraic) multiplicity of $\lambda$ as an eigenvalue of some $H(\theta)$. More precisely, let $\lambda \in\{z$ : $\left.-2 \beta_{0}<\arg z<3 \pi / 2+2 \beta_{0}\right\}$ be a resonance of $H$, there exists $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$such that $-2 a(\theta)<\arg \lambda$. Lemma 1 implies that $\lambda$ is also an eigenvalue of $H(\theta)$, then we define the multiplicity of resonance $\lambda$ as follows:

$$
\begin{equation*}
m(\lambda):=m_{\theta}(\lambda) \equiv-\frac{1}{2 \pi i} \operatorname{tr} \oint_{\lambda}(H(\theta)-z)^{-1} d z \tag{3.9}
\end{equation*}
$$

where the integral is over a positively oriented circle enclosing $\lambda$ and containing no eigenvalues of $H(\theta)$ other than $\lambda$. To see that the multiplicity $m(\lambda)$ is well-defined,
we need to show that $m(\lambda)$ does not depend on the choice of $\theta$. Assume $\theta_{0}, \theta_{1} \in D_{\beta_{0}}$ satisfy $-2 a\left(\theta_{0}\right) \leq-2 a\left(\theta_{1}\right)<\arg \lambda$, let $\theta_{t}=(1-t) \theta_{0}+t \theta_{1}$ then $-2 a\left(\theta_{t}\right)<\arg \lambda$ for all $t \in[0,1]$. Let $C_{\lambda}$ be a positively oriented circle enclosing $\lambda$ with sufficiently small radius such that $C_{\lambda} \subset\left\{z: \arg z>-2 a\left(\theta_{1}\right)\right\}$ and contains no resonances of $H$ other than $\lambda$. Therefore, $C_{\lambda}$ contains no eigenvalues of $H\left(\theta_{t}\right)$ other than $\lambda$ for all $t \in[0,1]$ as a consequence of Lemma 1. Now we have

$$
m_{\theta_{t}}(\lambda)=-\frac{1}{2 \pi i} \operatorname{tr} \int_{C_{\lambda}}\left(H\left(\theta_{t}\right)-z\right)^{-1} d z, \quad t \in[0,1] .
$$

Hence $m_{\theta_{t}}(\lambda)$ depends continuously on $t$ which implies that $m_{\theta_{t}}(\lambda)$ must be a constant as it is integer-valued. In particular, we have $m_{\theta_{0}}(\lambda)=m_{\theta_{1}}(\lambda)$, thus $m(\lambda)$ is welldefined.

## 4. Eigenvalues and complex scaling

In this section we will show that the eigenvalues of $H_{\varepsilon} \equiv-\Delta-i \varepsilon x^{2}+V$ are invariant under complex scaling, in other words, these eigenvalues are the same as the eigenvalues of

$$
H_{\varepsilon}(\theta):=U_{\theta} H_{\varepsilon} U_{\theta}^{-1}=p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}+V\left(\phi_{\theta}(x)\right), \quad \theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}
$$

First we recall some basic properties about the Davies harmonic oscillator and its deformation, see $[\mathrm{Z} 2, \S 3]$ for details. The operator $H_{\varepsilon, \gamma}:=-\Delta+e^{-i \gamma} \varepsilon x^{2}, \varepsilon>0$, $0 \leq \gamma<\pi$, was used by Davies [Da1] to illustrate properties of non-normal differential operators. We are more interested in the deformations of $H_{\varepsilon, \gamma}$ under complex scaling. Let

$$
Q_{\varepsilon, \theta}=-\Delta_{\theta}-i \varepsilon x_{\theta}^{2}, \quad \text { where } x_{\theta}=\left.z\right|_{\Gamma_{\theta}}
$$

be a deformed operator on $\Gamma_{\theta}$ as in [Z2, §3]. In view of (3.4), we have

$$
\begin{equation*}
p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}=U_{\theta} Q_{\varepsilon, a(\theta)} U_{\theta}^{-1}, \quad \theta \in D_{\beta_{0}} . \tag{4.1}
\end{equation*}
$$

Hence we can study the spectrum and the resolvents of $p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}$ using the relevant results about $Q_{\varepsilon, a(\theta)}$. We recall [Z2, Lemma 4.] that $\sigma\left(Q_{\varepsilon, a(\theta)}\right)=\sqrt{\varepsilon} e^{-i \pi / 4}\left(n+2\left|\mathbb{N}_{0}^{n}\right|\right)$ for $\theta \in D_{\beta_{0}}$, then by (4.1) we have

Proposition 8. For $\theta \in D_{\beta_{0}}, \varepsilon>0$, the spectrum of $p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}$ is independent of $\theta$ and given by $\sqrt{\varepsilon} e^{-i \pi / 4}\left(n+2\left|\mathbb{N}_{0}^{n}\right|\right)$.

For the resolvents of $p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}$ :

$$
\begin{equation*}
R_{\varepsilon, \theta}(z):=\left(p_{\theta}^{2}-i \varepsilon \phi_{\theta}(x)^{2}-z\right)^{-1}, \quad \theta \in D_{\beta_{0}} \tag{4.2}
\end{equation*}
$$

we recall [Z2, Lemma 5.] that for $\delta>0,-\pi / 8<\theta<\pi / 8$, we have

$$
\left(Q_{\varepsilon, \theta}-z\right)^{-1}=\mathcal{O}_{\delta}(1 /|z|): L^{2}\left(\Gamma_{\theta}\right) \rightarrow L^{2}\left(\Gamma_{\theta}\right), \quad-2 \theta+\delta<\arg z<3 \pi / 2+2 \theta-\delta,
$$

uniformly for $0<\varepsilon<\varepsilon_{0}$, where $\varepsilon_{0}>0$ is a constant. Using (4.1), we have

Proposition 9. Let $\theta \in D_{\beta_{0}}, \delta>0$, then uniformly for $0<\varepsilon<\varepsilon_{0}$, we have

$$
\begin{equation*}
R_{\varepsilon, \theta}(z)=\mathcal{O}_{\delta}(1 /|z|): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right), \quad-2 a(\theta)+\delta<\arg z<3 \pi / 2+2 a(\theta)-\delta \tag{4.3}
\end{equation*}
$$

Now we state the main result about the eigenvalues of $H_{\varepsilon}$ :
Lemma 2. For any $\theta \in D_{\beta_{0}}, 0<\varepsilon<\varepsilon_{0}$,

$$
z \mapsto R_{H_{\varepsilon}(\theta)}(z) \equiv\left(H_{\varepsilon}(\theta)-z\right)^{-1}, \quad-\pi / 4<\arg z<7 \pi / 4,
$$

is a meromorphic family of operators on $L^{2}\left(\mathbb{R}^{n}\right)$ with poles of finite rank. Furthermore, the poles of $\left(H_{\varepsilon}(\theta)-z\right)^{-1}$ do not depend on $\theta \in D_{\beta_{0}}$ and coincide, with agreement of multiplicities, with the poles of $\left(H_{\varepsilon}-z\right)^{-1}$.

Proof. For fixed $\theta \in D_{\beta_{0}}$, one can compute

$$
\left(H_{\varepsilon}(\theta)-z\right) R_{\varepsilon, \theta}(z)=I+V\left(\phi_{\theta}(x)\right) R_{\varepsilon, \theta}(z),
$$

then we obtain from (4.3) that

$$
\begin{gather*}
R_{H_{\varepsilon}(\theta)}(z)=R_{\varepsilon, \theta}(z)\left(I+V\left(\phi_{\theta}(x)\right) R_{\varepsilon, \theta}(z)\right)^{-1}, \\
-2 a(\theta)+\delta<\arg z<3 \pi / 2+2 a(\theta)-\delta, \quad|z| \gg 1, \tag{4.4}
\end{gather*}
$$

where for large $|z|, I+V\left(\phi_{\theta}(x)\right) R_{\varepsilon, \theta}(z)$ is invertible by a Neumann series argument. Note that $R_{\varepsilon, \theta}(z): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow H^{2}\left(\mathbb{R}^{n}\right), \quad \arg z \neq-\pi / 4$ by Proposition 8 , recalling that $V\left(\phi_{\theta}(x)\right): H^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ is compact (see the proof of Proposition 7 ), we conclude that $z \mapsto V\left(\phi_{\theta}(x)\right) R_{\varepsilon, \theta}(z)$ is an analytic family of compact operators for $-\pi / 4<z<7 \pi / 4$. Hence $z \mapsto\left(I+V\left(\phi_{\theta}(x)\right) R_{\varepsilon, \theta}(z)\right)^{-1}$ is a meromorphic family of operators in the same range of $z$. In particular $z \mapsto R_{H_{\varepsilon}(\theta)}(z), \quad-\pi / 4<z<7 \pi / 4$ is a meromorphic family of operators on $L^{2}\left(\mathbb{R}^{n}\right)$ with poles of finite rank.

The poles and their multiplicities are independent of $\theta$. For that we modify the proof of Lemma 1 and define matrix elements:

$$
G_{f, g}(z)=\left\langle f,\left(H_{\varepsilon}-z\right)^{-1} g\right\rangle,
$$

and

$$
G_{f, g}(z ; \theta)=\left\langle U_{\bar{\theta}} f,\left(H_{\varepsilon}(\theta)-z\right)^{-1} U_{\theta} g\right\rangle,
$$

for all $f, g \in \mathcal{A}$.
Note that $-2 a(\theta)+\pi / 4<\pi / 2<3 \pi / 2+2 a(\theta)-\pi / 4$ since $-\beta_{0}<a(\theta)<\beta_{0}, \theta \in D_{\beta_{0}}$, using (4.3) and Neumann series argument, $H_{\varepsilon}(\theta)-z$ is invertible at $z=i \rho, \rho \gg 1$ for each $\theta \in D_{\beta_{0}}$. Like (3.8), we have

$$
\theta \in D_{\beta_{0}} \rightarrow G_{f, g}(z ; \theta) \equiv\left\langle U_{\bar{\theta}} f, R_{H(\theta)}(z) U_{\theta} g\right\rangle
$$

is an analytic map provided $z=i \rho, \rho \gg 1$. Hence we have

$$
\begin{equation*}
G_{f, g}(z ; \theta)=G_{f, g}(z), \quad \forall \theta \in D_{\beta_{0}}, \quad z=i \rho, \rho \gg 1, \tag{4.5}
\end{equation*}
$$

since this is true for all $\theta \in D_{\beta_{0}} \cap \mathbb{R}$. Now fix any $\theta \in D_{\beta_{0}}$, note that $G_{f, g}(z)$ and $G_{f, g}(z ; \theta)$ are both meromorphic in $-\pi / 4<z<7 \pi / 4$, we conclude that

$$
\begin{equation*}
G_{f, g}(z ; \theta)=G_{f, g}(z), \quad-\pi / 4<z<7 \pi / 4 \tag{4.6}
\end{equation*}
$$

by (4.5) and the identity principle of meromorphic functions.
Now argue as in the end of the proof of Lemma 1: if $\left(H_{\varepsilon}-z\right)^{-1}$ has a pole at $\lambda_{\theta} \in \mathbb{C} \backslash e^{-i \pi / 4}[0, \infty)$, then there must exist $f, g \in \mathcal{A}$ such that $\lambda_{\theta}$ is a pole of $G_{f, g}(z ; \theta)$, by (4.6), $\lambda_{\theta}$ is also a pole of $G_{f, g}(z)$ thus $\left(H_{\varepsilon}(\theta)-z\right)^{-1}$ must have a pole at $\lambda_{\theta}$ and vise versa. Hence for any $\theta \in D_{\beta_{0}}$, the poles of $\left(H_{\varepsilon}(\theta)-z\right)^{-1}$ in $\mathbb{C} \backslash e^{-i \pi / 4}[0, \infty)$ coincide the poles of $\left(H_{\varepsilon}-z\right)^{-1}$ in $\mathbb{C} \backslash e^{-i \pi / 4}[0, \infty)$.

To show the agreement of multiplicities, for any pole $\lambda$ of $\left(H_{\varepsilon}(\theta)-z\right)^{-1}$, the multiplicity of $\lambda$ is defined by

$$
m_{\varepsilon, \theta}(\lambda)=-\frac{1}{2 \pi i} \operatorname{tr} \oint_{\lambda}\left(H_{\varepsilon}(\theta)-z\right)^{-1} d z
$$

where the integral is over a positively oriented circle independent of $\theta$ enclosing $\lambda$ and containing no poles other than $\lambda$. Since $m_{\varepsilon, \theta}(\lambda)$ is continuous on $\theta \in D_{\beta_{0}}$ and integer-valued, it must be independent of $\theta \in D_{\beta_{0}}$. Hence we have

$$
m_{\varepsilon, \theta}(\lambda)=m_{\varepsilon, 0}(\lambda)=-\frac{1}{2 \pi i} \operatorname{tr} \oint_{\lambda}\left(H_{\varepsilon}-z\right)^{-1} d z
$$

which is the multiplicity of $\lambda$ as a pole of $\left(H_{\varepsilon}-z\right)^{-1}$.

## 5. Meromorphic continuation

In this section we will introduce a new way to express the meromorphic continuations of resolvents $R_{H(\theta)}(z)$ and $R_{H_{\varepsilon}(\theta)}(z)$ in a given region $\Omega \Subset\{z:-2 a(\theta)<\arg z<$ $3 \pi / 2+2 a(\theta)\}$, which is crucial in the proof of Theorem 1. For that we will first review some properties about $R_{\theta}(z)$ and the weighted $L^{2}$ space, $\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$.

Lemma 3. Let $\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$ be a weighted $L^{2}$ space with the norm

$$
\begin{equation*}
\|u\|_{\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)}=\left\|\langle x\rangle^{2} u\right\|_{L^{2}\left(\mathbb{R}^{n}\right)} . \tag{5.1}
\end{equation*}
$$

Then $H^{2}\left(\mathbb{R}^{n}\right) \cap\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$ is compactly embedded in $L^{2}\left(\mathbb{R}^{n}\right)$.
Proof. Let $u_{n} \in H^{2}\left(\mathbb{R}^{n}\right) \cap\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$ with $\left\|u_{n}\right\|_{H^{2}\left(\mathbb{R}^{n}\right)} \leq 1$ and $\left\|\langle x\rangle^{2} u_{n}\right\|_{L^{2}\left(\mathbb{R}^{n}\right)} \leq 1$. For some $r>0$ to be decided, we have

$$
\int_{|x| \geq r}\left|u_{n}(x)\right|^{2} d x \leq\langle r\rangle^{-4} \int_{|x| \geq r}\langle x\rangle^{4}\left|u_{n}(x)\right|^{2} d x \leq\langle r\rangle^{-4}\left\|\langle x\rangle^{2} u_{n}\right\|_{L^{2}\left(\mathbb{R}^{n}\right)}^{2}=\langle r\rangle^{-4} .
$$

Then we choose $r$ sufficientlt large such that $\int_{|x| \geq r}\left|u_{n}(x)\right|^{2} d x<1 / 8$ for all $n$. Since $H^{2}\left(B(0, r) \Subset L^{2}\left(B(0, r)\right.\right.$, there exists subsequence $\left\{u_{n}^{(1)}\right\} \subset\left\{u_{n}\right\}$ satisfying

$$
\int_{B(0, r)}\left|u_{n}^{(1)}(x)-u_{m}^{(1)}(x)\right|^{2} d x<1 / 2, \quad \text { for all } n, m
$$

Hence we have

$$
\begin{aligned}
\left\|u_{n}^{(1)}-u_{m}^{(1)}\right\|_{L^{2}\left(\mathbb{R}^{n}\right)}^{2} & =\int_{B(0, r)}\left|u_{n}^{(1)}(x)-u_{m}^{(1)}(x)\right|^{2} d x+\int_{|x| \geq r}\left|u_{n}^{(1)}(x)-u_{m}^{(1)}(x)\right|^{2} d x \\
& <1 / 2+\int_{|x| \geq r}\left(2\left|u_{n}^{(1)}(x)\right|^{2}+2\left|u_{m}^{(1)}(x)\right|^{2}\right) d x \\
& <1 / 2+2 / 8+2 / 8=1 .
\end{aligned}
$$

By the same argument, we can find $\left\{u_{n}^{(1)}\right\} \supset \cdots \supset\left\{u_{n}^{(j)}\right\} \supset \cdots$ with

$$
\left\|u_{n}^{(j)}-u_{m}^{(j)}\right\|_{L^{2}\left(\mathbb{R}^{n}\right)}<1 / j, \quad \text { for all } n, m
$$

Then the subsequence $\left\{u_{j}^{(j)}\right\} \subset\left\{u_{n}\right\}$ is a Cauchy sequence in $L^{2}\left(\mathbb{R}^{n}\right)$.
Lemma 4. Fix $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}, R_{\theta}(z)$ is an analytic family of operators $\langle x\rangle^{-2} L^{2} \rightarrow$ $\langle x\rangle^{-2} L^{2}$ for $-2 a(\theta)<\arg z<2 \pi-2 a(\theta)$. Furthermore, if $\Omega \Subset\{z:-2 a(\theta)<\arg z<$ $2 \pi-2 a(\theta)\}$ then there exists $C=C_{\Omega, \theta}$ such that

$$
\left\|R_{\theta}(z)\right\|_{\langle x\rangle^{-2} L^{2} \rightarrow\langle x\rangle^{-2} L^{2}} \leq C, \quad z \in \Omega
$$

Proof. In view of (2.1) and (2.9), we have

$$
|x| / 2<\left|\phi_{\theta}(x)\right|=|1+\theta h(|x|)||x|<3|x| / 2 \quad \Longrightarrow \quad\langle x\rangle / 2<\left\langle\phi_{\theta}(x)\right\rangle<3\langle x\rangle / 2 .
$$

Then it is equivalent to prove the lemma with $\left\langle\phi_{\theta}(x)\right\rangle$ replacing $\langle x\rangle$. We recall Proposition 3 to write

$$
\left\langle\phi_{\theta}(x)\right\rangle^{2} R_{\theta}(z)\left\langle\phi_{\theta}(x)\right\rangle^{-2}=U_{a(\theta)}\left\langle x_{a(\theta)}\right\rangle^{2}\left(-\Delta_{a(\theta)}-z\right)^{-1}\left\langle x_{a(\theta)}\right\rangle^{-2} U_{a(\theta)}^{-1},
$$

where $x_{a(\theta)}$ is the coordinate on $\Gamma_{a(\theta)}$. Then it suffices to show that, for any $0<\alpha<\beta_{0}$,

$$
\begin{equation*}
\langle w\rangle^{2}\left(-\Delta_{\alpha}-\lambda^{2}\right)^{-1}\langle w\rangle^{-2}: L^{2}\left(\Gamma_{\alpha}\right) \rightarrow L^{2}\left(\Gamma_{\alpha}\right), \quad \operatorname{Im}\left(e^{i \alpha} \lambda\right)>0 \tag{5.2}
\end{equation*}
$$

is analytic with uniformly bounded norm provided $\lambda$ in any compact subset of $\{\lambda \in$ $\left.\mathbb{C}: \operatorname{Im}\left(e^{i \alpha} \lambda\right)\right\}$, where $w$ denotes the coordinate on $\Gamma_{\alpha}$. To prove (5.2), consider the integral kernel of that operator:

$$
\begin{equation*}
K\left(\lambda, w_{1}, w_{2}\right)=\left\langle w_{1}\right\rangle^{2} R_{0}\left(\lambda, w_{1}, w_{2}\right)\left\langle w_{2}\right\rangle^{-2}, \quad w_{1}, w_{2} \in \Gamma_{\alpha} \tag{5.3}
\end{equation*}
$$

where $R_{0}\left(\lambda, w_{1}, w_{2}\right)$ is the integral kernel of $\left(-\Delta_{\alpha}-\lambda^{2}\right)^{-1}: L^{2}\left(\Gamma_{\alpha}\right) \rightarrow L^{2}\left(\Gamma_{\alpha}\right)$. It is easy to see that

$$
\begin{align*}
\left|K\left(\lambda, w_{1}, w_{2}\right)\right| & \leq\left(1+\left|w_{1}\right|^{2}\right)\left|R_{0}\left(\lambda, w_{1}, w_{2}\right)\right|\left\langle w_{2}\right\rangle^{-2} \\
& \leq 2\left(1+\left|w_{1}-w_{2}\right|^{2}+\left|w_{2}\right|^{2}\right)\left(1+\left|w_{2}\right|^{2}\right)^{-1}\left|R_{0}\left(\lambda, w_{1}, w_{2}\right)\right|  \tag{5.4}\\
& \leq 2\left(1+\left|w_{1}-w_{2}\right|^{2}\right)\left|R_{0}\left(\lambda, w_{1}, w_{2}\right)\right| .
\end{align*}
$$

To introduce the explicit formula of $R_{0}\left(\lambda, w_{1}, w_{2}\right)$, we recall that one can define ( $\left(w_{1}-\right.$ $\left.\left.w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}$ for $w_{1}, w_{2} \in \Gamma_{\alpha}$, see [DyZ2, §4.5]. Then we can write
$R_{0}\left(\lambda, w_{1}, w_{2}\right)=C_{n} \lambda^{n-2}\left(\lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}\right)^{-\frac{n-2}{2}} H_{\frac{n}{2}-1}^{(1)}\left(\lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}\right)$
where $H_{k}^{(1)}$ denote the Hankel functions of the first kind, and we can estimate $\left|R_{0}\left(\lambda, w_{1}, w_{2}\right)\right|$ as follows:

$$
\begin{equation*}
\left|R_{0}\left(\lambda, w_{1}, w_{2}\right)\right| \leq \frac{P_{n}\left(\lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}\right)}{\left(\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}\right)^{n-2}} e^{-\operatorname{Im} \lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}} \tag{5.5}
\end{equation*}
$$

where $P_{n}$ is a polynomial of degree $(n-3) / 2$, see [GaSm, §2.2] and [DyZ2, §4.5] for details. Using (2.11), it is easy to see that for any $\delta$ small, there exists $C_{\delta}>0$ such that $\left|\arg \left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}-\alpha\right|<\delta$ provided $\left|w_{1}-w_{2}\right|>C_{\delta}$. Note that $0<\arg \lambda+\alpha<\pi$, for every $\lambda$, we can choose $\delta=\delta_{\lambda}$ such that $2 \delta<\arg \lambda+\alpha<\pi-2 \delta$, then for $|z-w|>C_{\lambda}$, we have

$$
\delta<\arg \lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}<\pi-\delta
$$

and thus

$$
\begin{equation*}
e^{-\operatorname{Im} \lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}}<e^{-c_{\lambda}\left|w_{1}-w_{2}\right|}, c_{\lambda}>0, \quad \text { if }\left|w_{1}-w_{2}\right|>C_{\lambda} . \tag{5.6}
\end{equation*}
$$

Then using (5.4), (5.5) and (5.6), we conclude that

$$
\sup _{w_{1} \in \Gamma_{\alpha}} \int_{\Gamma_{\alpha}}\left|K\left(\lambda, w_{1}, w_{2}\right)\right| d w_{1}<M_{\lambda}, \quad \sup _{w_{2} \in \Gamma_{\alpha}} \int_{\Gamma_{\alpha}}\left|K\left(\lambda, w_{1}, w_{2}\right)\right| d z<M_{\lambda} .
$$

By Schur criterion, we proved (5.2), the analyticity in $\lambda$ is easy to see using the explicit formula of $R_{0}\left(\lambda, w_{1}, w_{2}\right)$. If $\lambda \in K \Subset\left\{\lambda \in \mathbb{C}: \operatorname{Im}\left(e^{i \alpha} \lambda\right)\right\}$, then there exist $c_{K}$ and $C_{K}$ such that

$$
e^{-\operatorname{Im} \lambda\left(\left(w_{1}-w_{2}\right) \cdot\left(w_{1}-w_{2}\right)\right)^{1 / 2}}<e^{-c_{K}\left|w_{1}-w_{2}\right|}, c_{K}>0, \quad \text { if }\left|w_{1}-w_{2}\right|>C_{K} .
$$

Follow the above argument, there exists $M=M_{K}>0$ such that

$$
\left\|\langle w\rangle^{2}\left(-\Delta_{\alpha}-\lambda^{2}\right)^{-1}\langle w\rangle^{-2}\right\|_{L^{2}\left(\Gamma_{\alpha}\right) \rightarrow L^{2}\left(\Gamma_{\alpha}\right)}<M_{K}, \quad \text { for all } \lambda \in K,
$$

which completes the proof.
Now we state the main result of this section:

Lemma 5. Fix any $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$and $\Omega \Subset\{z:-2 a(\theta)<\arg z<3 \pi / 2+2 a(\theta)\}$, there exists $\chi \in \mathcal{C}_{\mathrm{c}}^{\infty}\left(\mathbb{R}^{n}\right)$, $\chi \equiv 1$ on $B(0, T)$ for some $T>0$ such that for $0 \leq \varepsilon<\varepsilon_{0}$, $H_{\varepsilon}(\theta)-\chi V-z$ is invertible in $\Omega$ and

$$
z \mapsto\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V\right)^{-1}, \quad z \in \Omega
$$

is a meromorphic family of operators on $L^{2}\left(\mathbb{R}^{n}\right)$ with poles of finite rank, where we write $R_{H_{\varepsilon}(\theta)-\chi V}(z)=\left(H_{\varepsilon}(\theta)-\chi V-z\right)^{-1}$ for simplicity. Moreover,

$$
\begin{equation*}
m_{\varepsilon, \theta}(z):=\frac{1}{2 \pi i} \operatorname{tr} \oint_{z}\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1} \partial_{w}\left(R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right) d w \tag{5.7}
\end{equation*}
$$

where the integral is over a positively oriented circle enclosing $z$ and containing no poles other than possibly $z$, satisfies

$$
\begin{equation*}
m_{\varepsilon, \theta}(z)=\frac{1}{2 \pi i} \operatorname{tr} \oint_{z}\left(w-H_{\varepsilon}(\theta)\right)^{-1} d w, \quad 0 \leq \varepsilon<\varepsilon_{0} \tag{5.8}
\end{equation*}
$$

where $H_{0}(\theta)=H(\theta)$.
Proof. We modify the argument in $[Z 2, \S 4]$ to our setting. First there exists $\delta=\delta_{\Omega}$ such that $\Omega \subset \mathcal{C}_{\delta}:=\{z:-2 a(\theta)+\delta<\arg z<3 \pi / 2+2 a(\theta)-\delta,|z|>\delta\}$, we recall (3.3) and (4.3) that uniformly for $0 \leq \varepsilon<\varepsilon_{0}$, we have

$$
\begin{equation*}
R_{\varepsilon, \theta}(z)=\mathcal{O}_{\delta}(1 /|z|): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right), \quad z \in \mathcal{C}_{\delta} \tag{5.9}
\end{equation*}
$$

Hence $\left\|R_{\varepsilon, \theta}(z)\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)}<C_{\delta}, \forall z \in \mathcal{C}_{\delta}$, for some $C_{\delta}>0$. In view of (1.1), for $T$ sufficiently large, we have $\|(1-\chi) V\|_{L^{\infty}}<1 / 2 C_{\delta}$ and thus

$$
\begin{equation*}
\left\|R_{\varepsilon, \theta}(z)(1-\chi) V\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)}<1 / 2, \quad \text { for all } z \in \mathcal{C}_{\delta} \tag{5.10}
\end{equation*}
$$

Then $\left(I+R_{\varepsilon, \theta}(z)(1-\chi) V\right)$ is invertible by the Neumann series argument, which implies that $H_{\varepsilon}(\theta)-\chi V-z$ is invertible and

$$
\begin{equation*}
R_{H_{\varepsilon}(\theta)-\chi V}(z)=\left(H_{\varepsilon}(\theta)-\chi V-z\right)^{-1}=\left(I+R_{\varepsilon, \theta}(z)(1-\chi) V\right)^{-1} R_{\varepsilon, \theta}(z), \quad \forall z \in \mathcal{C}_{\delta} . \tag{5.11}
\end{equation*}
$$

Since $\chi V \in L^{\infty}\left(\mathbb{R}^{n}\right)$, (5.9) and (5.11) imply that for $z \in \mathcal{C}_{\delta},|z| \gg 1$, both $I+$ $\chi V R_{H_{\varepsilon}(\theta)-\chi V}(z)$ and $I+R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V$ are invertible by the Neumann series argument. Hence we have

$$
\begin{align*}
R_{H_{\varepsilon}(\theta)}(z) & =R_{H_{\varepsilon}(\theta)-\chi V}(z)\left(I+\chi V R_{H_{\varepsilon}(\theta)-\chi V}(z)\right)^{-1} \\
& =R_{H_{\varepsilon}(\theta)-\chi V}(z) \sum_{j=0}^{\infty}(-1)^{j}\left(\chi V R_{H_{\varepsilon}(\theta)-\chi V}(z)\right)^{j} \\
& =R_{H_{\varepsilon}(\theta)-\chi V}(z)\left(I-\chi V \sum_{j=0}^{\infty}(-1)^{j}\left(R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V\right)^{j} R_{H_{\varepsilon}(\theta)-\chi V}(z)\right) \\
& =R_{H_{\varepsilon}(\theta)-\chi V}(z)\left[I-\chi V\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V\right)^{-1} R_{H_{\varepsilon}(\theta)-\chi V}(z)\right], \tag{5.12}
\end{align*}
$$

Using (5.11), we have

$$
R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V=\left(I+R_{\varepsilon, \theta}(z)(1-\chi) V\right)^{-1} R_{\varepsilon, \theta}(z) \chi V
$$

For $\varepsilon>0, R_{\varepsilon, \theta}(z): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow H^{2}\left(\mathbb{R}^{n}\right) \cap\left\langle\phi_{\theta}(x)\right\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$, then Lemma 3 implies that $R_{\varepsilon, \theta}(z): L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ is compact. For $\varepsilon=0$, note that $\chi V: L^{2}\left(\mathbb{R}^{n}\right) \rightarrow$ $\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$, by Lemma 4 we have $R_{\theta}(z) \chi V: L^{2}\left(\mathbb{R}^{n}\right) \rightarrow H^{2}\left(\mathbb{R}^{n}\right) \cap\langle x\rangle^{-2} L^{2}\left(\mathbb{R}^{n}\right)$, then Lemma 3 implies that $R_{\theta}(z) \chi V: L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$ is compact. Hence we can conclude that $z \mapsto R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V$ is an analytic family of compact operators for $z \in \mathcal{C}_{\delta}, \quad 0 \leq \varepsilon<\varepsilon_{0}$, and thus $z \mapsto\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V\right)^{-1}$ is a meromorphic family of operators in the same range of $z$.

Then we recall Lemma 1 and 2 that $R_{H_{\varepsilon}(\theta)}(z)$ is meromorphic in $-2 a(\theta)<\arg z<$ $3 \pi / 2+2 a(\theta)$, by the identity principle of meromorphic operators, we conclude that (5.12) holds for all $z \in \mathcal{C}_{\delta}$ in the sense of meromorphic family of operators.

To obtain the multiplicity formula, we assume that $z \in \Omega$, then there exists a neighborhood $z \in U \subset \Omega$ and finite rank operators $A_{j}, 1 \leq j \leq J$ such that

$$
\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1}-\sum_{j=1}^{J} \frac{A_{j}}{(w-z)^{j}} \quad \text { is holomorphic in } w \in U
$$

Let $C_{z} \subset U$ be a positively oriented circle enclosing $z$ and containing no poles of $\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1}$ other than possibly $z$, thus it also contains no poles of $(w-$ $\left.H_{\varepsilon}(\theta)\right)^{-1}$ other than possibly $z$ as a consequence of (5.12). On the one hand, we can compute

$$
\begin{align*}
m_{\varepsilon, \theta}(z) & =\frac{1}{2 \pi i} \operatorname{tr} \int_{C_{z}}\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1} \partial_{w}\left(R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right) d w \\
& =\frac{1}{2 \pi i} \operatorname{tr} \int_{C_{z}}\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1} R_{H_{\varepsilon}(\theta)-\chi V}(w)^{2} \chi V d w \\
& =\frac{1}{2 \pi i} \operatorname{tr} \int_{C_{z}} \sum_{j=1}^{J} \frac{A_{j} R_{H_{\varepsilon}(\theta)-\chi V}(w)^{2} \chi V}{(w-z)^{j}} d w  \tag{5.13}\\
& =\sum_{j=1}^{J} \frac{1}{(j-1)!} \operatorname{tr} \partial_{z}^{j-1}\left(A_{j} R_{H_{\varepsilon}(\theta)-\chi V}(z)^{2} \chi V\right) \\
& =\sum_{j=1}^{J} \sum_{k=0}^{j-1} \frac{1}{k!(j-1-k)!} \operatorname{tr} A_{j} \partial_{z}^{k} R_{H_{\varepsilon}(\theta)-\chi V}(z) \partial_{z}^{j-1-k} R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V .
\end{align*}
$$

## RESONANCES AS VISCOSITY LIMITS FOR EXTERIOR DILATION ANALYTIC POTENTIALS15

On the other hand, by (5.12) we have

$$
\begin{align*}
& \frac{1}{2 \pi i} \operatorname{tr} \oint_{z}\left(w-H_{\varepsilon}(\theta)\right)^{-1} d w \\
= & \frac{1}{2 \pi i} \operatorname{tr} \int_{C_{z}} R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1} R_{H_{\varepsilon}(\theta)-\chi V}(w) d w \\
= & \frac{1}{2 \pi i} \operatorname{tr} \int_{C_{z}} \sum_{j=1}^{J} \frac{R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V A_{j} R_{H_{\varepsilon}(\theta)-\chi V}(w)}{(w-z)^{j}} d w  \tag{5.14}\\
= & \sum_{j=1}^{J} \frac{1}{(j-1)!} \operatorname{tr} \partial_{z}^{j-1}\left(R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V A_{j} R_{H_{\varepsilon}(\theta)-\chi V}(z)\right) \\
= & \sum_{j=1}^{J} \sum_{k=0}^{j-1} \frac{1}{k!(j-1-k)!} \operatorname{tr} \partial_{z}^{j-1-k} R_{H_{\varepsilon}(\theta)-\chi V}(z) \chi V A_{j} \partial_{z}^{k} R_{H_{\varepsilon}(\theta)-\chi V}(z) .
\end{align*}
$$

Now we compare (5.13) and (5.14). Since $A_{j}$ factors have finite rank, we can apply cyclicity of the trace to obtain the multiplicity formula (5.8).

## 6. Proof of convergence

The proof of convergence is based on Lemma 1, Lemma 2, Lemma 5 and the following lemma:

Lemma 6. Fix any $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$and $\Omega \Subset\{z:-2 a(\theta)<\arg z<3 \pi / 2+2 a(\theta)\}$, there exists $\chi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$, $\chi \equiv 1$ on $B(0, T)$ for some $T>0$ such that for $0<\varepsilon<\varepsilon_{0}$,

$$
T_{\varepsilon, \theta}(z):=\left(H_{\varepsilon}(\theta)-\chi V-z\right)^{-1} \phi_{\theta}(x)^{2}(H(\theta)-\chi V-z)^{-1} \chi V
$$

is an analytic family of operators: $L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)$. Furthermore, there exists $C=$ $C_{\Omega, \theta}$ such that

$$
\begin{equation*}
\left\|T_{\varepsilon, \theta}(z)\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)} \leq C, \quad z \in \Omega, \quad \text { uniformly for } 0<\varepsilon<\varepsilon_{0} \tag{6.1}
\end{equation*}
$$

Proof. We recall the proof of Lemma 5 that for $T$ sufficiently large, $H_{\varepsilon}(\theta)-\chi V-z$ is invertible, then (5.10) and (5.11) imply that

$$
\left\|\left(H_{\varepsilon}(\theta)-\chi V-z\right)^{-1}\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)} \leq C_{\Omega}, \quad z \in \Omega
$$

for some $C_{\Omega}>0$. Hence it suffices to prove

$$
\begin{equation*}
\left\|\phi_{\theta}(x)^{2}(H(\theta)-\chi V-z)^{-1} \chi V\right\|_{L^{2}\left(\mathbb{R}^{n}\right) \rightarrow L^{2}\left(\mathbb{R}^{n}\right)} \leq C_{\Omega}, \quad z \in \Omega . \tag{6.2}
\end{equation*}
$$

By Lemma 4, we have $\left\|R_{\theta}(z)\right\|_{\langle x\rangle^{-2} L^{2} \rightarrow\langle x\rangle^{-2} L^{2}} \leq C$. We can choose $T$ sufficiently large such that (5.10) still holds and $\|(1-\chi) V\|_{L^{\infty}}<1 / 2 C$, then we have

$$
\left\|R_{\theta}(z)(1-\chi V)\right\|_{\langle x\rangle^{-2} L^{2} \rightarrow\langle x\rangle^{-2} L^{2}}<1 / 2, \quad z \in \Omega
$$

Hence $\left(I+R_{\theta}(z)(1-\chi) V\right)^{-1}: L^{2} \rightarrow L^{2}$ defined by the Neumann series in the proof of Lemma 5 also maps $\langle x\rangle^{-2} L^{2}$ to $\langle x\rangle^{-2} L^{2}$ by the same Neumann series and we have

$$
\begin{equation*}
\left\|\left(I+R_{\theta}(z)(1-\chi) V\right)^{-1}\right\|_{\langle x\rangle^{-2} L^{2} \rightarrow\langle x\rangle^{-2} L^{2}}<2, \quad z \in \Omega . \tag{6.3}
\end{equation*}
$$

Since $\chi V: L^{2} \rightarrow\langle x\rangle^{-2} L^{2}$ with the operator norm bounded by $\left\|\langle x\rangle^{2} \chi V\right\|_{L^{\infty}}=C_{\Omega}$, by Lemma 4, (5.11) and (6.3) we conclude that

$$
\begin{aligned}
& \left\|(H(\theta)-\chi V-z)^{-1} \chi V\right\|_{L^{2} \rightarrow\langle x\rangle^{-2} L^{2}} \\
= & \left\|\left(I+R_{\theta}(z)(1-\chi) V\right)^{-1} R_{\theta}(z) \chi V\right\|_{L^{2} \rightarrow\langle x\rangle^{-2} L^{2}} \\
\leq & \left\|\left(I+R_{\theta}(z)(1-\chi) V\right)^{-1} R_{\theta}(z)\right\|_{\langle x\rangle^{-2} L^{2} \rightarrow\langle x\rangle^{-2} L^{2}}\|\chi V\|_{L^{2} \rightarrow\langle x\rangle-2} L^{2} \\
\leq & C_{\Omega},
\end{aligned}
$$

which implies (6.2).
Now we state the result about the convergence of eigenvalues of the deformed operator $H_{\varepsilon}(\theta)$ :

Theorem 2. Fix $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$and $\Omega \Subset\{z:-2 a(\theta)<\arg z<3 \pi / 2+2 a(\theta)\}$, there exists $\delta_{0}=\delta_{0}(\Omega)$ satisfying the following:

For any $0<\delta<\delta_{0}$ there exists $\varepsilon^{\prime}>0$ such that for any $z \in \Omega$ with $m_{\theta}(z)>0$ and $0<\varepsilon<\varepsilon^{\prime}, H_{\varepsilon}(\theta)$ has $m_{\theta}(z)$ eigenvalues in $B(z, \delta)$, where $m_{\theta}(z)$ is the multiplicity of the eigenvalue of $H(\theta)$ at $z$ - see (3.9).

Proof. Since the eigenvalues of $H(\theta)$ are isolated and $\bar{\Omega}$ is compact, there are finite many $z \in \Omega$ with $m_{\theta}(z)>0$, we denote them by $z_{1}, \ldots, z_{J}$. Then we can choose $\delta_{0}$ such that $B\left(z_{j}, \delta_{0}\right), j=1, \ldots, J$ are disjoint.

Now we fix $\delta<\delta_{0}$, by Lemma $5, I+R_{H(\theta)-\chi V}(w) \chi V$ is invertible in $\Omega \backslash\left\{z_{1}, \ldots, z_{J}\right\}$, thus we have

$$
\left\|\left(I+R_{H(\theta)-\chi V}(w) \chi V\right)^{-1}\right\|_{L^{2} \rightarrow L^{2}}<C(\delta), \quad w \in \partial B(z, \delta), \text { for all } z \in\left\{z_{1}, \ldots, z_{J}\right\}
$$

for some $C(\delta)>0$. We note that in the notation of Lemma 6,

$$
I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V-\left(I+R_{H(\theta)-\chi V}(w) \chi V\right)=i \varepsilon T_{\varepsilon, \theta}(w) .
$$

Hence there exists $0<\varepsilon^{\prime}<\varepsilon_{0}$ such that for any $\varepsilon<\varepsilon^{\prime}$,

$$
\left\|\left(I+R_{H(\theta)-\chi V}(w) \chi V\right)^{-1}\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V-\left(I+R_{H(\theta)-\chi V}(w) \chi V\right)\right)\right\|<1
$$

on $\partial B(z, \delta)$. Now we apply Gohberg-Sigal-Rouché theorem, see [GS] and [DyZ2, Appendix C.] to obtain that

$$
\begin{aligned}
& \frac{1}{2 \pi i} \operatorname{tr} \int_{\partial B(z, \delta)}\left(I+R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right)^{-1} \partial_{w}\left(R_{H_{\varepsilon}(\theta)-\chi V}(w) \chi V\right) d w \\
= & \frac{1}{2 \pi i} \operatorname{tr} \int_{\partial B(z, \delta)}\left(I+R_{H(\theta)-\chi V}(w) \chi V\right)^{-1} \partial_{w}\left(R_{H(\theta)-\chi V}(w) \chi V\right) d w .
\end{aligned}
$$

Then we recall (5.8) to conclude that

$$
\frac{1}{2 \pi i} \operatorname{tr} \int_{\partial B(z, \delta)}\left(w-H_{\varepsilon}(\theta)\right)^{-1} d w=m_{\theta}(z)
$$

which implies that $H_{\varepsilon}(\theta)$ has $m_{\theta}(z)$ eigenvalues in $B(0, \delta)$.
Finally, we can give the proof of Theorem 1.
Proof. We assume from now on that $\varepsilon<\varepsilon_{0}$. Fix any $\Omega \Subset\left\{z:-2 \beta_{0}<\arg z<\right.$ $\left.3 \pi / 2+2 \beta_{0}\right\}$, we can choose $\theta \in D_{\beta_{0}} \cap \mathbb{C}^{+}$such that

$$
\Omega \Subset\{z:-2 a(\theta)<\arg z<3 \pi / 2+2 a(\theta)\} .
$$

In view of Lemma 1 , we see that $\left\{z_{j}\right\}_{j=1}^{\infty}$, the resonances of $H$ in $\Omega$, can be identified as the eigenvalues of $H(\theta)$, denoted by $\left\{z_{\theta, j}\right\}_{j=1}^{\infty}$. Similarly, Lemma 2 guarantees that $\left\{z_{j}(\varepsilon)\right\}_{j=1}^{\infty}$, the eigenvalues of $H_{\varepsilon}$ in $\{z:-2 a(\theta)<\arg z<3 \pi / 2+2 a(\theta)\}$, are the eigenvalues of $H_{\varepsilon}(\theta)$, denoted by $\left\{z_{\theta, j}(\varepsilon)\right\}_{j=1}^{\infty}$. Hence it suffices to show

$$
\begin{equation*}
z_{\theta, j}(\varepsilon) \rightarrow z_{\theta, j}, \quad \varepsilon \rightarrow 0+, \quad \text { uniformly on } \Omega \tag{6.4}
\end{equation*}
$$

which is a direct result of Theorem 2.

## References

[AgCo] J. Aguilar and J.M. Combes, A class of analytic perturbations for one-body Schrödinger Hamiltonians, Comm. Math. Phys. 22(1971), 269-279.
[BaCo] E. Balslev and J.M. Combes, Spectral properties of many-body Schrödinger operators wth dilation analytic interactions, Comm. Math. Phys. 22(1971), 280-294.
[BiZ] D. Bindel and M. Zworski, Theory and computation of resonances in $1 d$ scattering, online presentation, including MATLAB codes, http://www.cims.nyu.edu/~\{\}dbindel/resonant1d
[Da1] E.B. Davies, Pseudospectra, the harmonic oscillator and complex resonances, Proc. R. Soc. Lond. A 455(1999), 585-599.
[Da2] E Brian Davies. Pseudospectra of differential operators. Journal of Operator Theory, pages 243-262, 2000.
[DeSZ] N. Dencker, J. Sjöstrand, and M. Zworski, Pseudospectra of semiclassical differential operators, Comm. Pure Appl. Math 57 (2004), 384-415.
[DyZ1] S. Dyatlov and M. Zworski, Stochastic stability of Pollicott-Ruelle resonances, preprint, arXiv:1407.8531.
[DyZ2] Semyon Dyatlov and Maciej Zworski. Mathematical theory of scattering resonances. American Mathematical Society, 2019.
[GaSm] Jeffrey Galkowski and Hart F Smith. Restriction bounds for the free resolvent and resonances in lossy scattering. International Mathematics Research Notices, 2015(16):7473-7509, 2014.
[GS] IC u Gohberg and EI Sigal. An operator generalization of the logarithmic residue theorem and the theorem of rouché. Mathematics of the USSR-Sbornik, 13(4):603, 1971.
[HeSj] B. Helffer and J. Sjöstrand, Resonances en limite semiclassique, Bull. Soc. Math. France 114, no. 24-25, 1986.
[HS] Peter D Hislop and Israel Michael Sigal. Introduction to spectral theory: With applications to Schrödinger operators, volume 113. Springer Science \& Business Media, 2012.
[HSV] M. Hitrik, J. Sjöstrand, and J. Viola, Resolvent Estimates for Elliptic Quadratic Differential Operators, Analysis \& PDE 6(2013), 181-196.
[JZBRK] T-C. Jagau, D. Zuev, K. B. Bravaya, E. Epifanovsky, and A.I. Krylov, A Fresh Look at Resonances and Complex Absorbing Potentials: Density Matrix-Based Approach, J. Phys. Chem. Lett. 5(2014), 310-315.
[Ka] Tosio Kato. Perturbation theory for linear operators, volume 132. Springer Science \& Business Media, 2013.
[Le] G. Lebeau, Fonctions harmoniques et spectre singulier, Ann. Sci. École Norm.Sup. 13(1980), 269-291.
[Ma] A. Martinez Prolongement des solution holomorphe des problèmes aux limits, Ann.Inst. Fourier, 35(1985), 93-116.
[NZ1] S. Nonnenmacher and M. Zworski, Quantum decay rates in chaotic scattering, Acta Math. 203(2009), 149-233.
[NZ2] S. Nonnenmacher and M. Zworski, Decay of correlations for normally hyperbolic trapping, Invent. Math. $\mathbf{2 0 0}(2015)$, to appear.
[Rei] W. P. Reinhardt, Complex Scaling in Atomic and Molecular Physics, In and Out of External Fields, AMS Proceedings Series: Proceedings of Symposia in Pure Mathematics 76(2007), 357377.
[RiMe] U.V. Riss and H.D. Meyer, Reflection-Free Complex Absorbing Potentials, J. Phys. B 28 (1995), 1475-1493.
[SeMi] T. Seideman and W.H. Miller, Calculation of the cumulative reaction probability via a discrete variable representation with absorbing boundary conditions, J. Chem. Phys. 96(1992), 4412-4422.
[Si] B. Simon, The definition of molecular resonance curves by the method of exterior complex scaling, Phys. Lett. A 71(1979), 211-214.
[Sj1] J. Sjöstrand, Singularités analytiques microlocales, Astérisque, volume 95, 1982.
[Sj2] J. Sjöstrand, Geometric bounds on the density of resonances for semiclassical problems, Duke Math. J. 60(1990), 1-57.
[Sj3] J. Sjöstrand, Lectures on resonances, version préliminaire, printemps 2002.
[SjZ] J. Sjöstrand and M. Zworski, Complex scaling and the distribution of scattering poles, J. Amer. Math. Soc. 4(1991), 729-769.
[St1] P. Stefanov, Approximating resonances with the complex absorbing potential method, Comm. Partial Differential Equations 30(2005), 1843-1862.
[St2] Stanly Steinberg. Meromorphic families of compact operators. Archive for Rational Mechanics and Analysis, 31(5):372-379, 1968.
[V] A. Vasy, Microlocal analysis of asymptotically hyperbolic and Kerr-de Sitter spaces, with an appendix by Semyon Dyatlov. arXiv:1012.4391, Invent. Math., 194(2013), 381-513.
[Z1] Maciej Zworski, Semiclassical analysis, Graduate Studies in Mathematics 138, AMS, 2012.
[Z2] Maciej Zworski. Scattering resonances as viscosity limits. In Algebraic and Analytic Microlocal Analysis, pages 635-654. Springer, 2013.

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