

Lecture Topics

- Cobordism Hypothesis - see Dawd

Key Example: Finite Htpy types \rightarrow n-Alg [various]?

Correspondence cat & Quantization functors

Induction recipe, up to dim 4 mb.?

Boundary Cond. Dirichlet cond.

EM duality in Abelian case, examples of groups

[Self-dual theories in dim $(4k-1)$]

- EM duality. Freudenthal spectra. - stable htpy range
Example
Htpy type \rightarrow QFTs not "full" (more symmetries)
Flipping half the space down to a minimal version

Dirichlet cond: gives back a htpy type

Symmetries. Symmetries of QFTs.

Extended sps from defects. Defect calculus.

FH symmetries from group homs.

Normal bdr role in defect calculus.

Anomalous symmetries & examples (Zesting)

- Topological Condensate.

1 fold def. Rippling space theory

Higher cond (conjectural)

Embedded Dirichlet cond: need for defects

Cond to top dimension - examples

- Lattice models from TQFTs & defects

• Ising theo. Landau paradigm. KW duality in 2 of EFT

• Toric Code. Other TV models. RT models

I TQFTs : fully local \Rightarrow valued in a linear n -category
 with duals
 $+ \mapsto Z(+)$
 $- \mapsto Z(-) \cong Z(+)^{\vee}$

duality units & co-units have duals
 all the way to top dim.

$$X \otimes X^* \xrightarrow{ev} 1, 1 \xrightarrow{u} X^* \otimes X, \text{ relation}$$

In invertible case : dual = inverse

Caution at top : if asking for duals \Rightarrow inverse

\Rightarrow all objects below were invertible

[Useful in proving invertibility].

Main example (in higher dim) - FHT

X finite htop type $\xrightarrow[\text{path integral}]{\text{finite}}$ $\mathcal{Q}_{m-1}(X, \tau)$ on $(m-1)$ category
 $m \geq 0$
 $Z \in \mathbb{R}^n(X; \mathbb{C}^*)$ generates an m dim TQFT

Plan: explain relation betw. htop types and \mathcal{Q}_m

[Key ingredient: EM duality]

Illustrate symmetry & defect calculus

Condensation of defects

Applications to lattice models & questions.

Key open question: good target for TQFTs & \mathcal{Q}_m
 Conflicting desirable properties!

Some wishes:

- linearity
- algebra objects, bimodules should exist cat and
- finite co-invariants should exist and be invariant

(finite colimits? "Idempotent completion")

- later: units. Generated by its topological condition

Eg $G \hookrightarrow \text{Vect}$ (trivial)

coinv. cat: crossed product
same objects, add iso $x \xrightarrow{\varphi_g} gx + \text{rels.}$

Result: free $\mathbb{C}[G]$ -modules

Invar: $(\alpha, \varphi_g: x \rightarrow gx + \text{coherence})$:

Result: $\text{Rep}(G)$

Have map $\text{Vect}_G \xrightarrow{\text{Norm}} \text{Vect}^G$

Need Idempotent compl to get iso.

1. \mathcal{Q}_m by induction:

$$m=0: \mathcal{Q}_{-1}(X) = \sum_{p \in \pi_0 X} \prod_i \# \pi_i(X, p)^{(-1)^i} \cdot \tau(p) \quad (\text{Kontsevich})$$

[Only possible functional invariant]

X : built from homotopy groups "stacked up"

$\pi_0, \pi_1, \pi_2 \dots$

Recipe: iterated crossed product.

$$\mathcal{Q}_m(X) = \bigoplus_{p \in \pi_0 X} \mathcal{Q}_m(X_p)$$

$$Q_m(\text{connected } X) : Q_0(X) = \bigoplus \mathbb{C}$$

hey it's a group \Rightarrow algebra object
 \Rightarrow linear higher cat

Alg. Alg. Alg. . .

$$\text{Eg } m=0 \quad Q_0(X) = \bigoplus_{p \in \pi_0 X} \mathbb{C}_p = \mathbb{C}[\pi_0 X]$$

$$Q_1(X) : Q_0(\Omega X) = \bigoplus_{p \in \pi_1} \mathbb{C}_p$$

hey it's a group $\rightarrow \mathbb{C}\langle \pi_1 X \rangle$
 group m)

$$\Gamma(B\pi_1; \text{Vect}) = \text{Vect}^{\pi_1} = \text{Rep}(\pi_1)$$

$$Q_2(X) : Q_1(\Omega X) = \bigoplus_{p \in \pi_1 X} \mathbb{C}\langle \pi_2 X \rangle = \bigoplus_{p \in \pi_1 X} \text{Rep}(\pi_2 X)$$

Hey it's a group \Rightarrow new multiplication
 convolution n components

$$\text{Rep } \pi_2 X = \text{Vect}[\pi_2 X^r], \text{ yhmh.}$$

$$\pi_1 \times \text{Vect}[\pi_2 X^r]^{\otimes}$$

$$Q_3(X) : \pi_1 \times \left(\begin{array}{l} \text{Braided fusion cat from } \pi_2, \pi_3 \\ \mathbb{C}\langle \pi_2 \rangle \langle \pi_3 \rangle, * \\ k\text{-chars } H^*(B^2\pi_2; \pi_3) \end{array} \right)$$

$$\hookrightarrow \pi_2 \rightarrow \pi_3 \text{ quadratic}$$

Symmetric

$K(\pi_1, n) \rightarrow$ En object in (m-n) catynics
 . π with conn.

Composed
 category
 with τ

Twisted examples:

$m=0$ $\tau \in H^1(X; \mathbb{C}^x) \rightarrow$ line bundle
Kleinian $\mathbb{C}(\pi_0 X)$ w/ flat sections

$m=1$ $\tau \in H^2(X; \mathbb{C}^x) \rightarrow$ 2-cocycle

$m=2$ $\tau \in H^3(X; \mathbb{C}^x)$ comes from π_1
 $\pi_1^T \times (\text{loc. } \pi_1 X^v)$

$m=3$ $H^4(X; \mathbb{C}^x)$ contrib to braiding & DW for π_1 .

Caution: π_n & higher do not contribute to \mathcal{Q}_{n-1}
unless $\tau: \pi_n \rightarrow \mathbb{C}^x$ is nontrivial
then 0.

π_{n-1} can be replaced by a collect of DW from
[EM duality]

TAFT: $\mathcal{Q}: \text{Loc}_m \rightarrow (m\text{-catyng})$

$M^k \mapsto \mathcal{Q}_{m-k-1}(\text{Map}(M^k; X))$

II. EM duality - Strictly Commutative Case

Strictly commutative case: in dim n ,

$$Q_{n-1}(\Sigma^p A) \cong Q_{n-1}(\Sigma^{n-p-1} A^\vee)$$

by

$$\begin{array}{ccc} & (\Sigma^p A \times \Sigma^{n-p-1} A^\vee, b) & \\ & \swarrow \quad \searrow & \\ \Sigma^p A & & \Sigma^{n-p-1} A^\vee. \end{array}$$

Eg $n=1, p=0$: Fourier transform $\mathcal{C}(A) \cong \mathcal{C}(A^\vee)$

$n=2, p=1$: $BA \leftrightarrow A^\vee$

$$(\mathcal{C}(A)\text{-Mod}) \leftrightarrow \text{Vect}(A^\vee).$$

$n=3, p=1$: Gauge theories for A and A^\vee interchanged

$$(\text{Rep}(A^\vee), \otimes =) \text{Vect}(A) \cong \underset{\text{Vect}}{\text{Vect}(A^\vee)} (= \text{Rep}(A), \otimes)$$

Vect: trivial as left & right A, A^\vee module sep

NONTRIVIAL as $A \times A^\vee$ bimodule

[Heisenberg extension]

$$\text{Vect}(A) \stackrel{\otimes}{=} \text{End}_{\text{Vect}(A^\vee)}(\text{Vect}).$$

Remark this holds in nonab. case with $\text{Vect}(G), \text{Rep}(G)$

In general: $Q_{n-1}(\Sigma^{n-p-1} A^\vee) = \text{End}_{Q_{n-1}(\Sigma^p A)}(\text{Vect})$

Examples in 3D

$\text{Vect}(G)$ alg object $\mathbb{1}_g \leftrightarrow$ reg. module

alg object $\bigoplus_g \mathbb{C}_g = \mathbb{A}$

$$\mathbb{A} \otimes \mathbb{A} = \bigoplus_{g,h} \mathbb{C}_{gh} \xrightarrow{\text{sum}} \bigoplus_k \mathbb{C}_k$$

modules = G vect bds $\sim G, \text{Vect}, \text{Newman}$

$H \subset G$

$\bigoplus_{h \in H} \mathbb{C}_h$ convolution: module = $\text{Vect}(G/H)$

generating object: trivial bdd

Chal ext $\tilde{H}: \text{Vect}^*(G/H)$, trivial bdd still ok

$\text{Rep}(G)$

alg object $\mathbb{1} \leftrightarrow$ reg module

alg object $\mathbb{C}[G]$ as left G -module, pointwise
 $\hookrightarrow G$ equivariant vect bds $\sim G, \text{Vect}$

alg object $\mathbb{C}[G/H]$ as left G -module, pointwise

$\rightarrow \text{Rep}(H)$. Generatr: trivial rep.

Chal extnsm $\tilde{H} \hookrightarrow \text{Rep}(\tilde{H})$

Choose proj rep E of H , $G \times_H \text{End}(E)$ algebra

Exercise (1) Work out statements for twisting

(2) For Abelian group: $(\text{Rep}(A), \otimes) = (\text{Vect}(A^*), *)$

Find correspondence betw. centrally extended subgroups!

EM duality for Spectra

← postponed to III

Orthogonal group action on category of spectra:

naturally trivialized;

Second trivialization of action = action on Idl functor

= J-action of O on spectra

\Rightarrow two different oriented theory structures
interchanged by EM duality \checkmark

Boundary conditions on spectra:

$$F \hookrightarrow Y \rightarrow X$$

$$\Sigma^{n-2} F^{\vee} \rightarrow \Sigma^{n-1} X^{\vee} \rightarrow \Sigma^{n-1} Y^{\vee}$$

Dirichlet: $Y = X \Leftrightarrow F = \Sigma^1 X \Leftrightarrow \Sigma^{n-2} F^{\vee} = \Sigma^{n-1} X^{\vee}$ Neumann

Caution with connectivity:

$Y \rightarrow X$ "Dirichlet" \Leftrightarrow
(generating)

$\pi_0 Y \rightarrow \pi_0 X$ onto

$\pi_{n-2} Y \rightarrow \pi_{n-1} X$ zero

\Downarrow

$(\pi_{n-2} X)^{\vee} \rightarrow (\pi_{n-2} Y)^{\vee}$ zero

$(\pi_0 X)^{\vee} \rightarrow (\pi_0 Y)^{\vee}$ into

\Downarrow

$\pi_0 \Sigma^{n-2} F^{\vee} \rightarrow \pi_0 \Sigma^{n-1} X^{\vee}$ onto

$\pi_{n-1} \Sigma^{n-2} F^{\vee} \rightarrow \pi_{n-1} \Sigma^{n-2} X^{\vee}$ zero

Thm Every $Y \rightarrow X$ satisfying \uparrow
is a 'generating' boundary
condition: the algebra object

$\mathcal{Q}_n(Y \times X)$ is $\cong \mathcal{Q}_n(X)$.

Ex explain π_{n-2} condition.

III. Recap, Complements & Examples
 Representations. Duality for spectra reinterpreted.
 Boundary conditions, duality & symmetries

$$\Sigma^p A \leftrightarrow \Sigma^{d-p-1} A^\vee \quad \text{give equivalent TFTs}$$

\mathcal{Q}_{d-1} : algebra objects in a (dri) category, Mod $\mathcal{A} \cong$.

More general:

$$A' \hookrightarrow A \rightarrow A'', e \in \text{Ext}^1(A'', A') \Leftrightarrow b: A'' \times (A')^\vee \rightarrow \Sigma \mathbb{C}^x$$

\Downarrow

$$\Sigma^p A \leftrightarrow \Sigma^p A'' \times \Sigma^{d-p-1} (A')^\vee, \quad \tau: \Sigma^p A'' \times \Sigma^{d-p-1} (A')^\vee \rightarrow \Sigma^d \mathbb{C}^x$$

equivalent TFTs

Spaces X : $d = \begin{matrix} 2n \\ 2n-1 \end{matrix} \quad X_{\geq n} \hookrightarrow X \twoheadrightarrow X_{< n}$

$\left[\begin{array}{l} \infty \text{ loop space, except possibly when } n \text{ even} \\ \text{and a quadratic } k^{2n} \in H^{2n}(\Sigma^n \pi_n; \pi_{2n-1}) \end{array} \right.$

\rightarrow Requires SPECIAL HANDLING when n even:

$\Sigma^n \pi_n$, τ nondeg. quadratic to \mathbb{C}^x
 gives invertible TFTs. Can be formally trivialized.

This defines a new class of "self-dual theories"
 in dim $2n-1$ which have no top. ∂ conditions.

Assuming that taken care of,

flip to $X_{< n} \times \Sigma^{d-1}(X_{\geq n}^\vee)$, τ

τ valued in $\Sigma^d \mathbb{I}_{\mathbb{C}^x}$

linear on Σ^{d-1}

from k -invariant $[X_{< n}, \Sigma X_{\geq n}]$

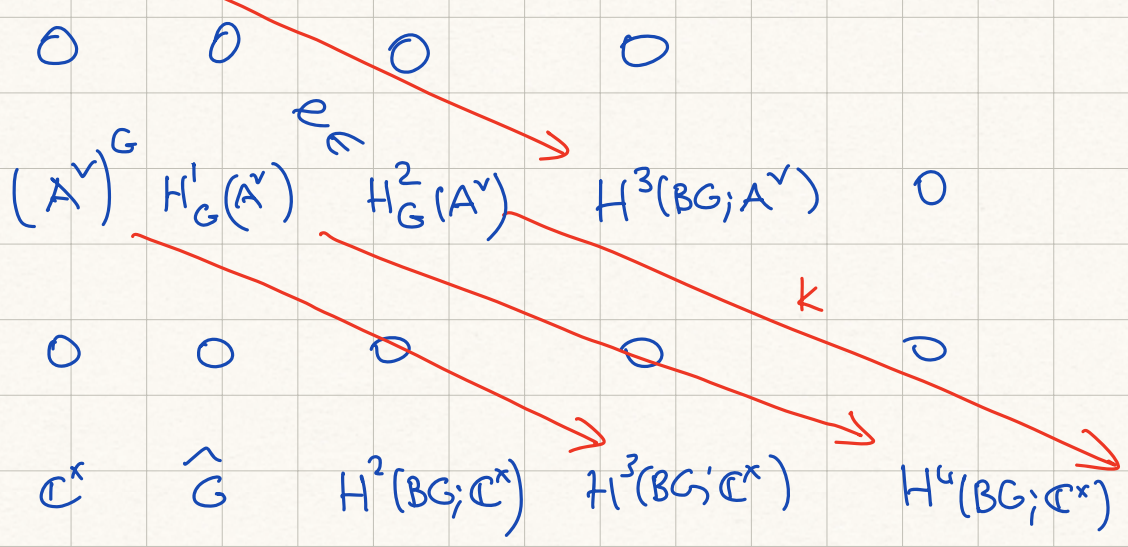
$n=4$: Partial duality : 2-group \rightarrow pure gauge theory

2-group: $B^2A \hookrightarrow X \rightarrow BG$

$\pi_2 = A$
 $\pi_1 = G$

- G acts on A
- $k \in H^3(BG; A)$
- $\tau \in H^4(X; \mathbb{C}^x)$

$H^4(B^2A; \mathbb{C}^x)$ ← Quadratic part:
assume 0 (else an invertible factor)



Partial EM dual:

extension $1 \rightarrow A^v \hookrightarrow \tilde{G} \xrightarrow{\sim} G \rightarrow 1$

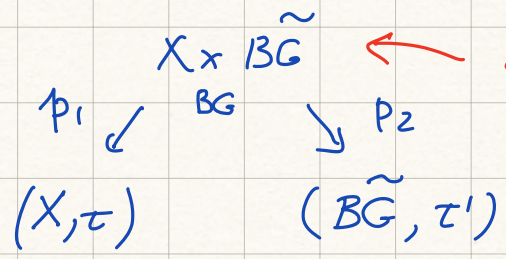
$e \in H^2_G(A^v)$

$\tau' \in H^4_{\tilde{G}}(\mathbb{C}^x)$

$\left\{ \begin{array}{l} k \in H^3(BG; \overbrace{H^1(BA^v; \mathbb{C}^x)}^A)) \\ H^4(BG; \mathbb{C}^x) \text{ from } \tau \end{array} \right.$

Exercise:

EM duality induced by
Correspondence with
suitable τ'



$\beta_2^* \tau' - \beta_1^* \tau$
trivializable

Representations of $\mathcal{Q}_{d-1}(X)$

On objects in the same $(d-1)$ category:

Formally, Local systems of such objects over X
= Fixed points (invariants) of ΩX on
said $(d-1)$ category
= "representation of ΩX " in same

$\mathcal{Q}_{d-1}(X) \rightarrow \text{Reps}$ is meant to be a form of
"idempotent completion".

Commutative Case, $X = \Sigma^p A$: $\Sigma^{d-p-1} A^\vee$ is the
"moduli space" of strictly commutative reps of
 $\Sigma^p A$ on Vect as a $(d-2)$ category ($\Sigma^{d-3} \text{Vect}$)
if $p < d-1$, a unique rep with
automorphism g_p $\Sigma^{d-p-2} A^\vee$

from bicharacter $\Sigma^p A \times \Sigma^{d-p-2} A^\vee \rightarrow \Sigma^{d-1} \mathbb{C}^\times$
= central extension of \mathbb{I} by $\Sigma^{d-1} \mathbb{C}^\times = \text{Aut Id}(\text{Vect}_{\text{cat}}^{(d-2)})$

Problem Action of $\Sigma^p A$ on $\Sigma^{d-3} \text{Vect}$ up to iso
 $\longleftrightarrow H^{d-1}(\Sigma^p A; \mathbb{C}^\times)$

these are components of the moduli space of reps
MISMATCH!

Reason: Mixing metaphors.

Strictly abelian world = chain complexes

⇒ moduli space is $\text{RHom}(\Sigma^{p-1}A, \Sigma^{d-1}\mathbb{C}^x) = \Sigma^{d-p-2}A^\vee$.

Cohomology comes from the homological world,
replacement for injective module $\mathbb{Q}/\mathbb{Z}(\mathbb{C}^x)$ is $\mathbb{I}\mathbb{C}^x$
(Linear) replacement for RHom is

Stack Map $(\mathcal{X}; \mathbb{I}\mathbb{C}^\vee) = \mathcal{X}^\vee$

↳ infinite loop spaces (E_∞ groups)

So $\Sigma^{d-p-2}A^\vee$ is the moduli of E_∞ 1-dim reps of $\Sigma^{p-1}A$ on $\Sigma^{d-1}\mathbb{C}$

- Not true that all 1-dim reps are E_∞ (true characters)
- Not true that every rep splits into 1-dims.

Eg $d=5, x = B^2A$, reps on $\Sigma^4\mathbb{C}$ $A = \mathbb{Z}/2$

$H^4(B^2A; \mathbb{C}^x) = \mathbb{Z}/4$ (quadratic forms)

braided tensor

4 BTCs based on A

2-category

each is a "twisted Heurmann" bdy condenser

(braiding braid)

$\mathbb{Z}/2 \subset \mathbb{Z}/4$ is stable ($iS\mathbb{Z}^2$)

Conducts clarification?

NO! killed after inclusion $\mathbb{C}^x \hookrightarrow \mathbb{C}^x$
 $\mathbb{Z}/2$

(super vector spaces)

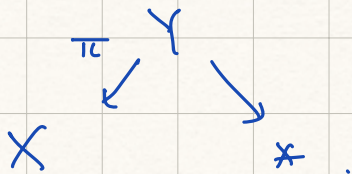
4d

Interesting representation: gauge theory BA (with DW twist)

Not a sum of 1d TFTs.

Representation as boundary conditions

Bdry conditions: interfaces with trivial theory



Associated local system on X : $R\pi_*(\Sigma^{d-1}\mathbb{C})$

Fiber = $(d-1)$ -quantization of fiber F

- has an ΩX action (in $\Sigma^{d-1}\mathbb{C}^x$)

$Y = X$, possibly with a cocycle: a Neumann cond

$Y =$ a point in each component

$F = (\mathbb{H}_r^{\pi_0}) \Omega_p X$ - with $\Omega_p X$ transl. action

"Dirichlet ∂ condition"

Note $\langle \text{Dirichlet} \mid \text{TQFT}_d \mid \text{Neumann} \rangle = \mathbb{1}$.

Can characterize Neumann conds in terms of Dirichlet.

Q What characterizes Dirichlet?

A: $\text{End}_{\mathcal{Q}_{d-1}(X)}(\mathbb{D}) \cong \mathcal{Q}_{d-1}(X)$ Morita, via \mathbb{D} .

Thm (Ostrik, Etingof et al) For an indecomposable fusion cat,
every non-zero module cat is Dirichlet.

Thm $Y \rightarrow X$ is Dirichlet iff:

* $\pi_0 Y \rightarrow \pi_0 X$ is onto

* $\pi_{d-1} Y \rightarrow \pi_{d-1} X$ is ZERO.

Eg: $BG \rightarrow BG$ NOT

Dirichlet in dim 2.

EM duality and 2 conditions

Orthogonal group action on category of spectra:

naturally trivialized;

Second trivialization of action = action on Idl functor

= J-action of O on spectra

\Rightarrow two different oriented Hoop structures
intertwined by EM duality \checkmark

Boundary conditions on spectra:

$$F \hookrightarrow \boxed{Y \rightarrow X} \quad \boxed{\Sigma^{n-2} F^{\vee} \rightarrow \Sigma^{n-1} X^{\vee}} \rightarrow \Sigma^{n-1} Y^{\vee}$$

Dirichlet: $Y = X \Leftrightarrow F = \Sigma^{-1} X \Leftrightarrow \Sigma^{n-2} F^{\vee} = \Sigma^{n-1} X^{\vee}$ Neumann

Caution with connectivity:

$$Y \rightarrow X \text{ "Dirichlet" } \Leftrightarrow \text{(generating)}$$

$$\begin{aligned} \pi_0 Y &\rightarrow \pi_0 X \text{ onto} \\ \pi_{n-1} Y &\rightarrow \pi_{n-1} X \text{ zero} \end{aligned}$$

Thm Every $Y \rightarrow X$ satisfying \uparrow
is a 'generating' boundary
condition: the algebraic object

$$\begin{aligned} &\Downarrow \\ (\pi_{n-1} X)^{\vee} &\rightarrow (\pi_{n-1} Y)^{\vee} \text{ zero} \\ (\pi_0 X)^{\vee} &\rightarrow (\pi_0 Y)^{\vee} \text{ onto} \end{aligned}$$

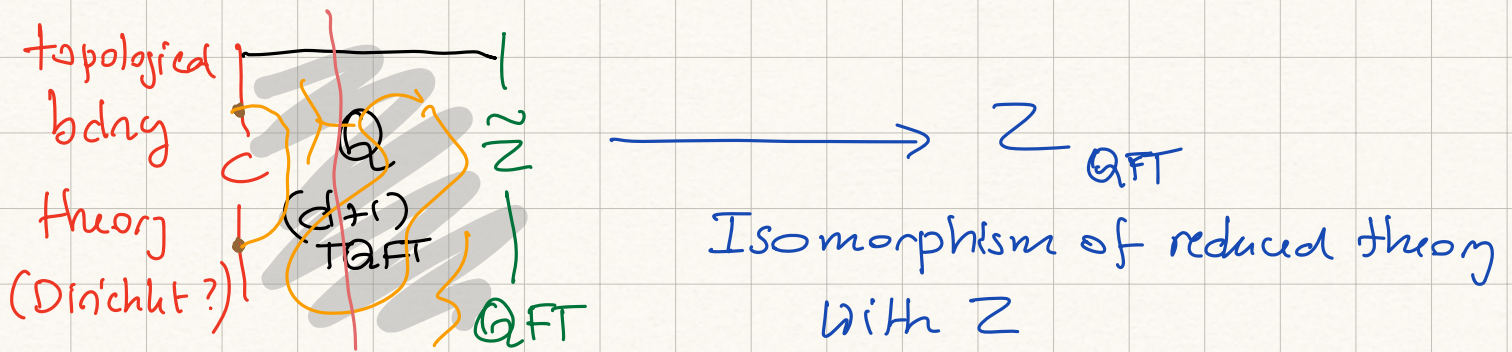
$$\mathcal{Q}_{n-1}(Y \times X) \text{ is } \cong \mathcal{Q}_{n-1}(X).$$

$$\begin{aligned} &\Downarrow \\ \pi_0 \Sigma^{n-2} F^{\vee} &\rightarrow \pi_0 \Sigma^{n-1} X^{\vee} \text{ ONTO} \\ \pi_{n-1} \Sigma^{n-2} F^{\vee} &\rightarrow \pi_{n-1} \Sigma^{n-1} X^{\vee} \text{ ZERO} \end{aligned}$$

Ex explain π_{n-1} condition.

Symmetries

Def A (generalized) SYMMETRY of a QFT Z is:



This should be viewed as an action of $\text{End}_Q(C)$ on Z

Extended topological operators in Z are the projections of topological defects in (C, Q) (not touching Z boundary)

Those not touching C are "central" in a sense (below)

Remarks (1) "(-1)-form symmetries" = self-domain walls of Q .

They change C . "Half-space" gauging can be done that way.

(2) "0-form symmetries" = self-defects of C .

(3) All interior operators may be pushed to the ∂ .
(but may acquire singular support)

(4) D_p, D_q may "braid" interestingly if they link in the ∂ .
But if one comes from the interior, braiding = trivial.

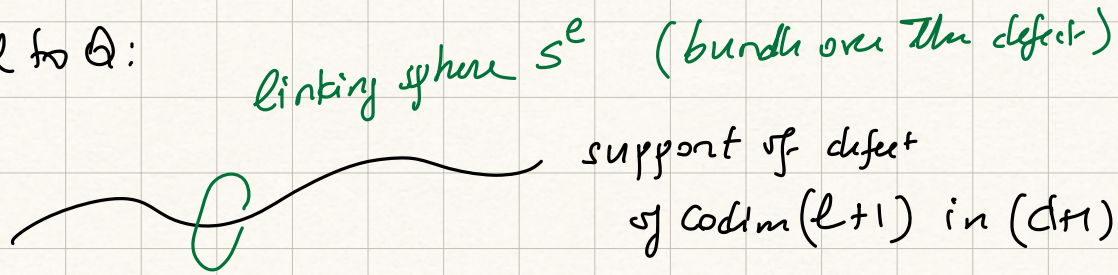
(5) More generally, this is true (in bulk or ∂)

if one of the defects is topologically condensed

$Q_1(x)$ case: Braiding is ~~nontrivial~~ on defects with full support.
nontrivial

Calculus of Defects.

Internal to \mathcal{Q} :



(Quantum) defect local label:

Boundary theory for $\mathcal{Q}(S^e)$ ($d+1-l$ dim TQFT)

If \mathcal{Q} comes from X , this is a representation

on $\Sigma^{d-e} \mathbb{C}^n$ of $\mathcal{Q}_{d-e}(\text{Map}(S^e; X))$

Caution - The spheres may rotate along the defect (w/ normal both)
- Even on oriented theories (eg X with $\tau \in H^{d+1}(X; \mathbb{C}^n)$)
the $SO(l+1)$ action on \dots may be nontrivial.
(Source: K -invariants of X)

We need a "flat" or locally constant ∂ theory
for a choice of quantum label all along defect

- Constraints on tangential structures are tricky to work out.

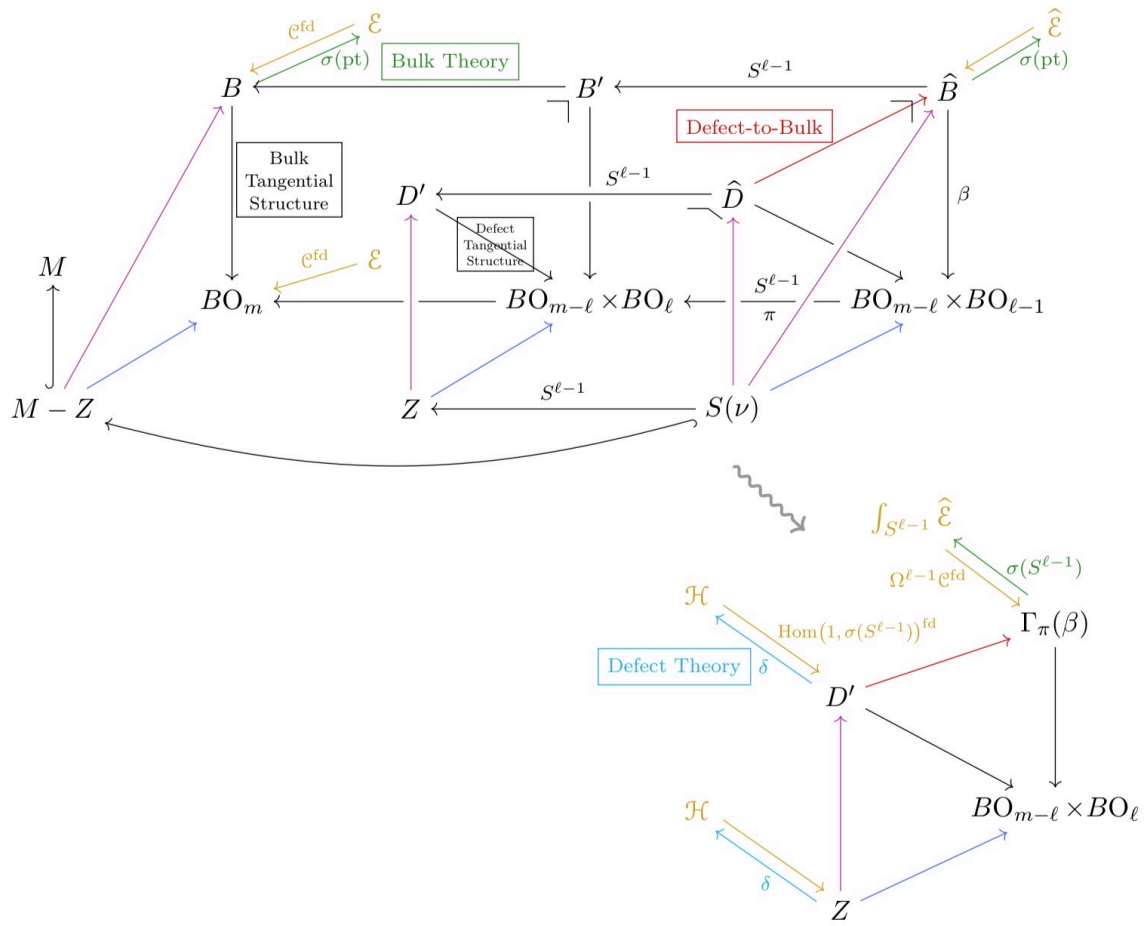


FIGURE 8. Local defect data, including tangential structures

Defects in an n -dimensional TQFT from a space X

$\mathcal{Q}_m(Y) :=$ quantization as algebra in an m -category

Top row = constant X = "Wilson defects"

Bottom row = ΩX = 't Hooft defects

Codimension of defects	Local quantum labels are modules over...	Contributing homotopy groups, placed in dim $[d]$	
		all/new defects	Condensed
n (=points)	$\mathcal{Q}_0(X^{S^{n-1}})$	π_0 $\pi_{n-1}[0]$	Wilson 't Hooft None
$n-1$ (lines)	$\mathcal{Q}_1(X^{S^{n-2}})$	π_0, π_1 $\pi_{n-2}[0], \pi_{n-1}[1]$	π_0 $\pi_{n-1}[1]$
$n-2$ (surfaces)	$\mathcal{Q}_2(X^{S^{n-3}})$	π_0, π_1, π_2 $\pi_{n-3}[0], \pi_{n-2}[1], \pi_{n-1}[2]$	π_0, π_1 $\pi_{n-2}[1], \pi_{n-1}[2]$
$k = n - \lfloor \frac{n-1}{2} \rfloor$	$\mathcal{Q}_{n-k}(X^{S^{k-1}})$	$\pi_0, \pi_1, \dots, \pi_{\lfloor \frac{n-1}{2} \rfloor}$ $\pi_{\lfloor \frac{n-1}{2} \rfloor}[0], \pi_k[1], \pi_{k+1}[2], \dots$	$\pi_0, \pi_1, \dots, \pi_{\lfloor \frac{n-1}{2} \rfloor - 1}$ $\pi_k[1], \pi_{k+1}[2], \dots$
$k = n - \lfloor \frac{n-1}{2} \rfloor - 1$		$\pi_0, \pi_1, \dots, \pi_{\lfloor \frac{n-1}{2} \rfloor + 1}$ $\pi_{k-1}[0], \pi_k[1], \pi_{k+1}[2], \dots$	$\pi_0, \pi_1, \dots,$ $\pi_k[1], \pi_{k+1}[2],$
2	$\mathcal{Q}_{n-2}(X^{S^1})$	$\pi_0, \pi_1, \dots, \pi_{n-2}$ $\pi_1[0], \pi_2[1], \dots, \pi_{n-1}[n-2]$	$\pi_0, \pi_1, \dots,$ $\pi_2[1], \dots, \pi_{n-1}[n-2]$
1 (interface)	$\mathcal{Q}_{n-1}(X^{S^0})$	$[\pi_0, \pi_1, \dots, \pi_{n-1}]$ $[\pi_0[0], \pi_1[1], \dots, \pi_{n-1}[n-1]]$	$\Delta \pi_0$ $\pi_1, \dots, \pi_{n-2},$ $\pi_1[1], \pi_2[2], \dots, \Delta \pi_{n-1}$
0 (space-filling)	$\mathcal{Q}_n(X^{\mathbb{S}^n})$	$\pi_0, \pi_1, \dots, \pi_n$ $\pi_0[0], \pi_1[1], \pi_2[2], \dots$	$\pi_0, \pi_1, \dots, \pi_{n-1}$ $\pi_0[1], \pi_1[2], \dots$

- DW cocycle $\tau \in \mathbb{I}_{\mathbb{C}^*}^n(X)$ can be integrated over S^k

π_1 -gauging of 't Hooft defects needed if π_1 acts nontrivially

1-fold Condensation

Work in progress w Freed, Hopkins;
 (multifold condensation, singular.. in progress)

Key ingredients

- ① notion of Dirichlet boundary condition for defects
 [Higher generalization with handles]
- ② Correspondence between internal algebra objects
 in a tensor category & module categories
 [Oshik for fusion cats; awaits full generalization]

Remark It is clear how this works if we work
 in the category of iterated algebras or categories.
 But extra wishes on the list need more care.

Example (3d, connected theory, fusion cat \mathcal{F})

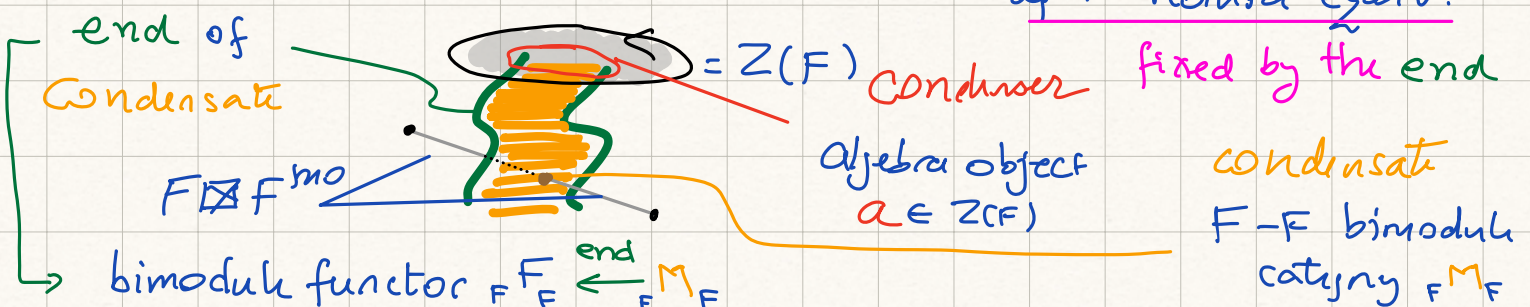
- (1) All condensed line operators are scalars
- (2) All surface operators are condensed

PS (1) Only scalar point ops (connectedness)

(2) surface ops $\leftrightarrow \mathcal{F} \boxtimes \mathcal{F}^{mo}$ -module cats

$\hookleftarrow \mathcal{Z}(\mathcal{F})$ -module cats \leftrightarrow algebra objects in $\mathcal{Z}(\mathcal{F})$

up to Morita equiv.



Precise Condition Observe the adjunction (TFTS \mathcal{J})

$$\text{Hom}_{\Omega^{e+1}C} (\mathbb{1}_{e+1}, \mathcal{J}(S^{e+1})) \cong \text{End}_{\text{Hom}(\mathbb{1}_e, \mathcal{J}(S^e))} (\mathcal{J}(D^{e+1}))$$

$$\mathbb{1}_e \xrightarrow{\mathcal{J}(D^{e+1})} \mathcal{J}(S^e)$$

$$\xleftarrow{\mathcal{J}(D^{e+1})^*} \mathbb{1}_e$$

$$\mathcal{J}(D^{e+1})^* \circ \mathcal{J}(D^{e+1})$$

$$\mathcal{J}(S^e)$$

$$\Omega_C^e := \underbrace{\text{End}(\text{End}(\dots(\text{End}(\mathbb{1}_0)\dots))}_{l \text{ times}}$$

target N -rat.

(Oshik principle)

(nice) algebra object in LHS = modules over RHS
(with a generator)

(Localization assumption)

Objects of $\text{Hom}(\mathbb{1}_e, \mathcal{J}(S^e))$ are "determined by"
their "localization at the unit" $\mathcal{J}(D^{e+1})$
"Localization at $\mathcal{J}(D^{e+1})$ is an equivalence"

$$M \in \text{Hom}_{\Omega_C^e} (\mathbb{1}_e, \mathcal{J}(S^e))$$

$$+ \text{end} \in \text{Hom}(\mathcal{J}(D^{e+1}), M)$$

localize



$$\text{Hom}_{\text{Hom}_{\Omega_C^e}(\mathbb{1}_e, \mathcal{J}(S^e))} (\mathcal{J}(D^{e+1}), M)$$

a module over

$$\text{End}_{\text{Hom}(\mathbb{1}_e, \mathcal{J}(S^e))} (\mathcal{J}(D^{e+1}))$$

+ generator

$$\text{algebra object}$$

$$a \in \text{Hom}(\mathbb{1}, \mathcal{J}(S^{e+1}))$$

$$+ \text{regular module}$$

≅



Oshik principle

Dirichlet condition on end says that you can "condense back" to the original (defect, end).

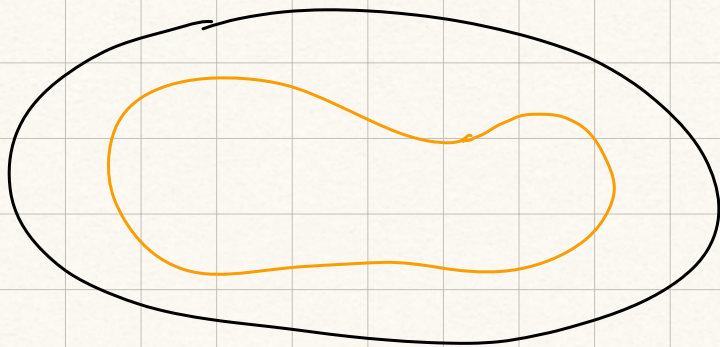
Overall

\mathcal{A} = algebra in $\text{codim}-(l+2)$ defect labels $\text{Hom}(\mathbb{1}_{2+l}, \mathcal{G}_d(x)[S^{2+l}])$
 + regular module (Condenser)

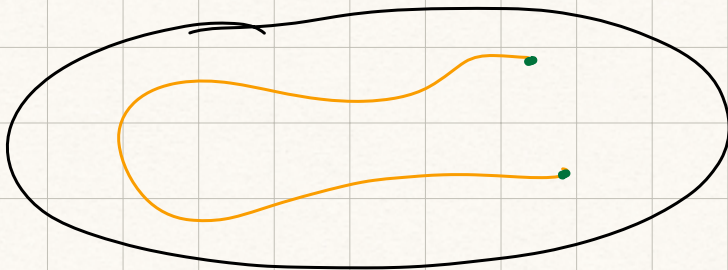
→ $\text{End} \dots [\mathcal{J}(D^{2+l})]$ module (+ generator)

→ $\text{codim}(l+1)$ defect label in $\text{Hom}(\mathbb{1}_e, \mathcal{G}_d(x)[S^e])$
 (+ end) (condensate)

"Ripping open theorem" [Weak assumption on Id. completion]



||



||



A top-dimensionally integrated 1-fold condensed defect may be punctured and retracted to one codim lower to produce the same state vector in a given enclosing hypersurface.

at the cost of embedding additional defects in the boundary.

$\mathbb{Q}_d(x) [s^{e+1}] \in (d-e-1)$ algebras //

δ label $\in \text{Hom}_{\Sigma^{d-e}} (\mathbb{1}, \mathbb{Q}_d(x) [s^{e+1}])$ in part a
 $(d-e-1)$ alj

if algebra object then $a(d-e)$ algebra

$$\text{Hom}_{\Sigma^{d-e}} (\mathbb{1}, \mathbb{Q}_d(x) [s^e])$$

$$\mathbb{Q}_d(x) [s^{e+1}] = \bigcup_{\mathbb{Q}_d(x) [s^e]} \otimes U^r$$

$$\text{Hom} (\mathbb{1}, \mathbb{Q}_d(x) [s^{e+1}]) = \Omega \text{Hom} (\mathbb{1}, \mathbb{Q}_d(x) [s^e])$$

\cup
 δ algebra object \rightarrow module m \uparrow
 object in
 $\text{Hom} (\mathbb{1}, \mathbb{Q}_d(x) [s^e])$
 with end

Adjunction:

$$\text{Hom} (\mathbb{1}, Z(s^{e+1})) = \text{End}_{\text{Hom} (\mathbb{1}, Z(s^e))} (Z(D^{e+1}))$$

\cup
 δ algebra \rightarrow module m \uparrow

