

# The qDE/IM Correspondence

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with E. Frenkel and A. Zeitlin, to appear

**Talk at University of Melbourne**

# Motivation

Quantum Geometry and Integrable Systems

[Okounkov et al]

[Pushkar, Zeitlin, Smirnov]

[PK, Pushkar, Smirnov, Zeitlin]

BPS/CFT Correspondence

[Nekrasov Shatashvili]

Geometric q-Langlands Correspondence

[Frenkel] [Aganagic, Frenkel, Okounkov]

CFT as Integrable System

[Bazhanov, Lukyanov, Zamolodchikov]

[Dorey, Dunning, Tateo]

# I. The ODE/IM Correspondence

[Bazhanov, Lukyanov,  
Zamolodchikov]  
[Dorey, Dunning,  
Tateo]

Consider Schrödinger equation

$$\Psi_1'' + \left( E - \frac{l(l+1)}{x^2} - x^{2M} \right) \Psi_1 = 0.$$

Can be presented in the vector form

$$\begin{aligned} \Psi_1' + \frac{l}{x} \Psi_1 + \Psi_2 &= 0 \\ \Psi_2' - \frac{l}{x} \Psi_2 + p(x, E) \Psi_1 &= 0 \end{aligned}$$

Symmetries

$$\Lambda : x \mapsto x, E \mapsto E, l \mapsto -l - 1$$

$$\Omega : x \mapsto qx, E \mapsto q^{-2}E, l \mapsto l$$

$$q = e^{\frac{i\pi}{1+M}}$$

# Asymptotics at Infinity

Solution of the Schrödinger equation at  $x = \infty$

$$\chi(x, E, l) \sim x^{-\frac{M}{2}} \exp\left(-\frac{x^{1+M}}{1+M} + \dots\right)$$

Let  $\chi^+(x, E, l) = \chi(x, E, l)$

Act with  $\Omega$  to get another solution

$$\chi^-(x, E, l) = iq^{-\frac{1}{2}} \chi(qx, q^{-2}E, l)$$

Wronskian

$$W[\chi^+, \chi^-] = 2$$

$$\Omega : x \mapsto qx, E \mapsto q^{-2}E, l \mapsto l$$

$$q = e^{\frac{i\pi}{1+M}}$$

# Asymptotics cont'd

Solutions around  $x = 0$

$$\psi(x, E, l) \sim ax^{l+1} + O(x^{l+3})$$

$$\Lambda : x \mapsto x, E \mapsto E, l \mapsto -l - 1$$

Act with  $\Lambda$  to get a different basis

$$\psi^+(x, E) = \psi(x, E, l)$$

$$\psi^-(x, E) = \psi(x, E, -l - 1)$$

Wronskian

$$W[\psi^+, \psi^-] = 2i(q^{l+\frac{1}{2}} - q^{-l-\frac{1}{2}})$$

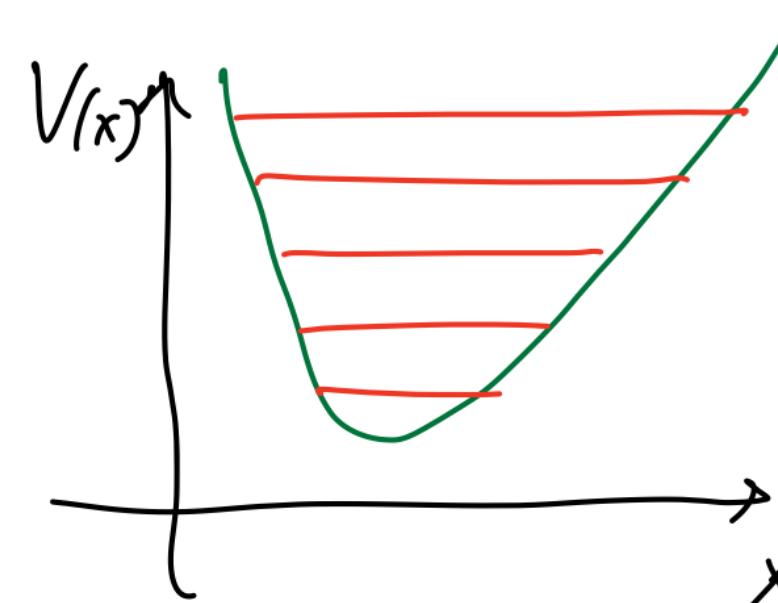
Discrete values of energy arise when  $\psi^\pm \rightarrow 0$  as  $x \rightarrow \infty$   $\{E_n^\pm\}$

Spectral determinants

$$D^\pm(E, l) = \prod_{n=1}^{\infty} \left(1 - \frac{E}{E_n^\pm}\right)$$

WKB approximation

$$E_n^\pm \sim n \frac{2M}{1+M}$$



# Expansion in $\chi$ -basis

Expand  $\psi$  in  $\chi$  basis

$$\psi^+ = C(E, l)\chi^+ + D(E, l)\chi^-$$

Act with  $\Omega, \Lambda$

$$\Lambda\psi^\pm = \psi^\mp$$

$$\Lambda\chi^\pm = \chi^\pm$$

$$\Omega\psi^\pm = q^{\frac{1}{2}\pm l\pm\frac{1}{2}}\psi^\pm$$

$$\Omega\chi^+ = -iq^{\frac{1}{2}}\chi^-$$

$$\Omega\chi^- = -iq^{\frac{1}{2}}\chi^+ + u\chi^-$$

Thus

$$C(E, l) = -iq^{-l-\frac{1}{2}}D(q^{-2}E, l)$$

$$\psi^- = D(E, -l-1)\chi^- - iq^{l+\frac{1}{2}}D(q^{-2}E, -l-1)\chi^+$$

# The QQ-System

Wronskian  $W[\psi^+, \psi^-] = 2i(q^{l+\frac{1}{2}} - q^{-l-\frac{1}{2}})$

$$q^{l+\frac{1}{2}} D(q^E, l) D(E, -l-1) - q^{-l-\frac{1}{2}} D(E, l) D(q^2 E, -l-1) = q^{l+\frac{1}{2}} - q^{-l-\frac{1}{2}}$$

yields the **QQ-system**  $\zeta Q^+(q^2 E) Q^-(E) - \zeta^{-1} Q^+(E) Q^-(q^2 E) = \Delta$

$D(E, l)$  are entire functions — eigenvalues of Baxter Q-operators which appear in the eight-vertex model.

[BLZ] description of  $c < 1$  CFT as completely integrable theory

QQ-system is ubiquitous to quantum integrable systems (XXX, XXZ, XYZ)

# Examples

Consider  $M = 1$

$c = -2$  CFT, i.e. the “free fermion” theory

$$E_n^+ = 4n + 2l - 1, \quad n = 1, 2, \dots,$$

$$D^+(E, l) = \frac{\Gamma(\frac{3}{4} + \frac{l}{2}) e^{cE}}{\Gamma(\frac{3}{4} + \frac{l}{2} - \frac{E}{4})}$$

Consider  $M \rightarrow \infty$

The Schrödinger potential becomes spherically symmetric rigid well

$$\begin{cases} 0, & \text{if } 0 < x < 1 \\ +\infty, & \text{if } x > 1 \end{cases}$$

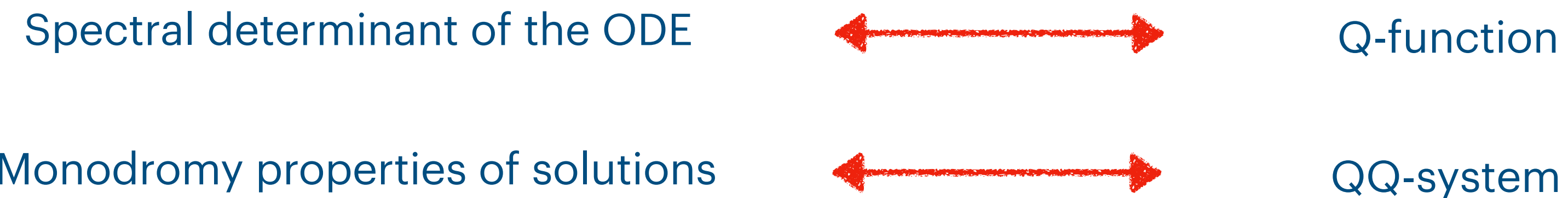
$$D^+(E, l) = \Gamma(l + 3/2) (\sqrt{E}/2)^{-l - \frac{1}{2}} J_{l + \frac{1}{2}}(\sqrt{E})$$



# The ODE/IM Correspondence

$$\psi^- = D(E, -l - 1)\chi^- - iq^{l+\frac{1}{2}}D(q^{-2}E, -l - 1)\chi^+$$

$$\zeta Q^+(q^2 E)Q^-(E) - \zeta^{-1}Q^+(E)Q^-(q^2 E) = \Delta$$



How can we understand this geometrically?

# II. Affine Opers

Study analytic solutions of the following linear problem

$$\begin{aligned}\Psi_1' + \frac{l}{x}\Psi_1 + \Psi_2 &= 0 \\ \Psi_2' - \frac{l}{x}\Psi_2 + p(x, E)\Psi_1 &= 0\end{aligned}$$

$$\Psi_1'' + \left( E - \frac{l(l+1)}{x^2} - x^{2M} \right) \Psi_1 = 0.$$

Can rewrite it using the following **affine** connection

$$\mathcal{L}(x, E) = \partial_x + \frac{\ell}{x} + e + p(x, E)e_0$$

where  $p(x, E) = x^{Mh^\vee} - E$ , with  $M > 0$  and  $E \in \mathbf{C}$ .

# Differential Operator

Consider connection  $\mathcal{L} = \partial_x + \frac{\ell}{x} + A(x, E)$   $A(x, E) \in \mathfrak{g}[x^M, E][[x]]$ .

**Proposition 2.1.** *There exist an element  $U \in G[[x, x^M, E]]$ , such that*

$$U^{-1}\mathcal{L}U = \partial_x + \frac{\ell}{x}$$

*and for any finite-dimensional representation  $V$  of  $\mathfrak{g}$ ,  $U(x, x^M, E)v = v + x\tilde{v}(x)$ , where  $\tilde{v}(x) \in V[E, x^M][[x]]$ .*

Allows to find a formal solution

**Theorem 2.2.** *There following expression  $W_g(x, E) = x^{-\ell}U(x)g$ , where  $g \in G$  and  $U \in G[[x, x^M, E]]$  constructed as in Proposition 2.1 gives a formal group-valued solution to the problem  $\mathcal{L}W = 0$ .*

# Canonical Solutions

**Corollary 2.3.** *In any highest weight representation  $V$ , choosing a standard basis  $\{v_i\}_{i=1}^{\dim(V)}$  according to the weight decomposition, so that  $\lambda_i = \text{wt}(v_i)$ , there is a family of  $V$ -valued solutions of equation  $\mathcal{L}\rho = 0$ , namely  $\{\varphi^{\lambda_i, v_i}(x, E)\}_{i=1}^{\dim(V)}$ , so that*

$$\varphi^{\lambda_i, v_i}(x, E) = W_e(x, E)v_i = x^{-\langle \ell, \lambda_i \rangle} (v_i + x\tilde{v}_i(x)), \quad i = 1, \dots, \dim(V),$$

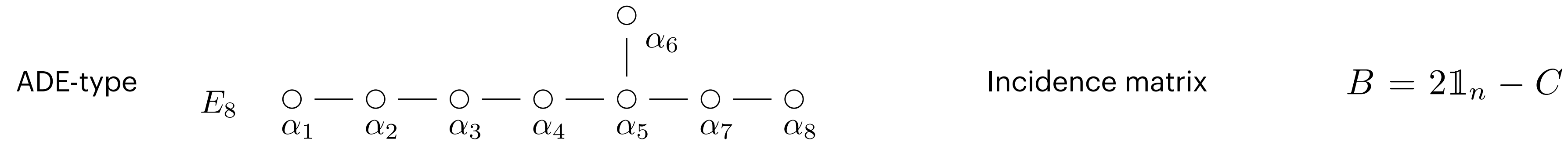
where  $\tilde{v}_i(x) \in V[E, x^M][[x]]$ .

**Proposition 2.4.** *Any analytic solution  $\Psi(x, E)$  in  $x, E$  of  $\mathcal{L}\Psi(x, E) = 0$  on  $x \in D \in \mathbb{C} \setminus \mathbb{R}$  and  $E \in \mathbb{C}$  can be decomposed in terms of formal solutions  $\varphi^{\lambda_i, v_i}$  in the following way:*

$$\Psi(x, E) = \sum_{i=1}^{\dim(V)} Q_{v_i}(E) \varphi^{\lambda_i, v_i}(x, E),$$

where  $Q_{v_i}(E)$  are analytic functions of  $E$ .

# Lie-Theoretic Data



$$[h_i, h_j] = 0, \quad [h_i, e_j] = C_{ij}e_j, \quad [h_i, f_j] = -C_{ij}f_j, \quad [e_i, f_j] = \delta_{ij}h_i$$

Fundamental weights  $\omega_i(h_j) = \delta_{ij}$       Simple roots  $\omega_i = \sum_{j \in I} (C^{-1})_{ji} \alpha_j, \quad i \in I$

Fundamental representations  $L(\omega_i), i \in I$

Affine Kac-Moody algebra  $\widehat{\mathfrak{g}} = \mathfrak{g} \otimes \mathbb{C}[t, t^{-1}] \oplus \mathbb{C}c$

$$[a \otimes f(t), b \otimes g(t)] = [a, b] \otimes f(t)g(t) + (a|b) \operatorname{Res}_{t=0} (f'(t)g(t)dt) c$$

$$[c, \widehat{\mathfrak{g}}] = 0.$$

Evaluation representation: as a vector space we take  $V(\zeta) = V$ .

$$(a \otimes f(t))v = f(\zeta)(av), \quad cv = 0, \quad \text{for } a \in \mathfrak{g}, f(t) \in \mathbb{C}[t, t^{-1}], v \in V$$

# Twisted Connection

Consider the evaluation representation  $V^{(i)}$  corresponding to the highest weight representation  $V_{\omega_i}$  of  $\mathfrak{g}$  with evaluation parameter  $t = e^{\pi i p(i)}$

Here  $p(i)$  is a homomorphism from the ordered set of vertices of the Dynkin diagram of  $\mathfrak{g}$  to  $\mathbb{Z}$

Denote  $\mathcal{L}_k$  the twisted differential operator by automorphism of

$$\widehat{\mathfrak{g}}: t \longrightarrow te^{2\pi i k}$$

$$\omega = e^{\frac{2\pi i}{h^\vee(M+1)}}, \quad \Omega = e^{\frac{2\pi i M}{M+1}} = \omega^{h^\vee M}$$

**Proposition 3.2.** *If  $\phi(x, E)$  is the solution of the linear problem  $\mathcal{L}(x, E)\phi(x, E) = 0$ , then  $\phi_k(x, E) = \omega^{-k\rho^\vee} \phi(\omega^k x, \Omega^k E)$  is a solution of linear problem  $\mathcal{L}_k(x, E)\phi_k(x, E) = 0$ .*

# WKB Analysis

**Proposition 3.8.** *Let  $\mathcal{L}(x, E)$  be the differential operator defined above. Then, we have the following gauge transformation of  $\mathcal{L}$ :*

$$(3.10) \quad q(x, E)^{\text{ad}\rho^\vee} \mathcal{L}(x, E) = \partial_x + q(x, E)\Lambda + \frac{\ell - M\rho^\vee}{x} + O(x^{-1-\delta})$$

$$p(x, E)^{\frac{1}{h^\vee}} = q(x, E) + O(x^{-1-\delta})$$

$$\Lambda = e_0 + e$$

$$\rho^\vee = \sum_{i=1}^r \omega_i^\vee$$

$$[\rho^\vee, e] = e, \quad [\rho^\vee, e_0] = -(h^\vee - 1)e_0$$

Let

$$S(x, E) = \int_0^x q(y, E) dy$$

**Theorem:** There exist a unique solution

$$\Psi(x, E) = e^{-\lambda S(x, E)} q(x, E)^{-h} (\psi + o(1)), \quad \text{in the sector } |\arg x| < \frac{\pi}{2(M+1)} - \delta$$

$$\Lambda\psi = \lambda\psi$$

# Family of Solutions

Consider  $\Psi_k(x, E) = \omega^{-k\rho^\vee} \Psi(\omega^k x, \Omega^k E)$

**Proposition 3.10.** *For any  $k \in \mathbb{R}$  such that  $|k| < \frac{h^\vee(M+1)}{2}$ , on the positive real semi-axis the function  $\Psi_k$  has the asymptotic behavior*

$$\Psi_k(x, E) = e^{-\lambda\gamma^k S(x, E)} q(x, E)^{-h} \gamma^{-kh} (\psi + o(1)), \quad x \gg 0,$$

where  $\Lambda\psi = \lambda\psi$   
 $[\frac{1}{2} - \frac{h^\vee}{2}, \frac{h^\vee}{2} + \frac{1}{2}]$ , are solutions

and  $\gamma = e^{\frac{2\pi i}{h^\vee}}$ . Then, the functions  $\Psi_s$ ,  $s \in \mathbb{Z}$ ,  $s \in$   
in the representation  $V_{\frac{h^\vee}{2} + \frac{1}{2}}$ .



# $\Psi$ -System

**Theorem:** Let  $\mathfrak{g}$  be a simple Lie algebra of ADE type, and let the solution  $\Psi(x, E)$

Let us denote the solution  $\Psi$  corresponding to representation  $V = V^{(i)}$  as  $\Psi^{(i)}$

$$m_i(\Psi_{-1/2}^{(i)}(x, E) \wedge \Psi_{1/2}^{(i)}(x, E)) = \bigotimes_{j>i \in I} \Psi^{(j)}(x, E)^{\otimes(-a_{ij})} \bigotimes_{j<i \in I} \Psi^{(j)}(x, E)^{\otimes(-a_{ij})}, \quad \forall i \in I.$$

Here unique morphism of representations

$$m_i : \bigwedge^2 L(\omega_i) \rightarrow \bigotimes_{j \in I} L(\omega_j)^{\otimes B_{ij}}$$

Such that

$$m_i(f_i v_i \wedge v_i) = w_i$$

$$w_i = \bigotimes_{j \in I} v_j^{\otimes B_{ij}}$$

# III. q-Operas

G-simple simply-connected complex Lie group

Principal G-bundle  $\mathcal{F}_G$  over  $D$

$$M_q : D \rightarrow D \\ z \mapsto qz$$

A meromorphic **(G,q)-connection** on  $\mathcal{F}_G$  is a section  $A$  of  $\text{Hom}_{\mathcal{O}_U}(\mathcal{F}_G, \mathcal{F}_G^q)$

U-open dense set

Choose U so that the restriction  $\mathcal{F}_G|_U$  of  $\mathcal{F}_G$  to  $U$  is isomorphic to a trivial G-bundle

$$A(z) \in G(\mathbb{C}(z)) \quad \text{on} \quad U \cap M_q^{-1}(U)$$

Change of trivialization  $A(z) \mapsto g(qz)A(z)g(z)^{-1}$

# (G,q)-Oper

A meromorphic (G,q)-oper on disc  $D$  is a triple  $(\mathcal{F}_G, A, \mathcal{F}_{B_-})$

$A$  is a meromorphic  $(G, q)$ -connection

$\mathcal{F}_{B_-}$  is a reduction of  $\mathcal{F}_G$  to  $B_-$

**Oper condition:** Restriction of the connection on some open dense set  $U$

$$A : \mathcal{F}_G \longrightarrow \mathcal{F}_G^q \text{ to } U \cap M_q^{-1}(U)$$

takes values in the double Bruhat cell

$$B_-(\mathbb{C}[U \cap M_q^{-1}(U)])cB_-(\mathbb{C}[U \cap M_q^{-1}(U)])$$

Coxeter element:  $c = \prod_i s_i$

Locally

$$A(z) = n'(z) \prod_i (\phi_i(z)^{\check{\alpha}_i} s_i) n(z)$$

$\phi_i(z) \in \mathbb{C}(z)$  and  $n(z), n'(z) \in N_-(z)$

# Miura $(G, q)$ -Operers

**Definition:** A *Miura  $(G, q)$ -oper* on  $D$  is a quadruple  $(\mathcal{F}_G, A, \mathcal{F}_{B_+}, \mathcal{F}_{B_-})$ , where  $(\mathcal{F}_G, A, \mathcal{F}_{B_+})$  is a meromorphic  $(G, q)$ -oper on  $D$  and  $\mathcal{F}_{B_-}$  is a reduction of the  $G$ -bundle  $\mathcal{F}_G$  to  $B_-$  that is preserved by the  $q$ -connection  $A$ .

Choose a trivialization  $\mathcal{F}_{G,x} \simeq G$  under this isomorphism

$$\begin{aligned}\mathcal{F}_{B_-,x} &\simeq aB_- \subset G \\ \mathcal{F}_{B_+,x} &\simeq bB_+\end{aligned}$$

Then  $a^{-1}b$  is a well defined element of the double quotient of  $B_- \backslash B / B_+ \simeq W_G$

Flags  $\mathcal{F}_{B_-}$  and  $\mathcal{F}_{B_+}$  are in *generic relative position* at  $x \in X$  if the corresponding element of the Weyl group assigned to them at  $x$  is equal to 1 or  $a^{-1}b \in B_- \cdot B_+$

# Structure Theorem

- Theorem:**
- i) Each regular Miura  $(G, q)$ -oper on a disk can be represented as an element of  $N_+((z)) \prod_i s_i N_+((z)) \cap B_-((z))$*
  - ii) Elements from the above intersection may be written in the form*

$$(4.3) \quad A(z) = \prod_i g_i(z)^{-\check{\alpha}_i} e^{\frac{t_i f_i}{g_i(z)}}, \quad g_i(z) \in \mathbb{C}((z))^\times,$$

where  $t_i \in \mathbb{C}^\times$  are complex parameters corresponding to the lift of  $c_i$  to  $s_i$ .

# Adding Singularities and Twists

Consider family of polynomials  $\{\Lambda_i(z)\}_{i=1,\dots,r}$

**(G,q)-oper with regular singularities** can be written as

$$A(z) = n'(z) \prod_i (\Lambda_i(z)^{\check{\alpha}_i} s_i) n(z), \quad n(z), n'(z) \in N_-(z)$$

Using structure theorem every Miura (G,q)-oper with singularities reads

$$A(z) = \prod_i g_i(z)^{\check{\alpha}_i} e^{\frac{\Lambda_i(z)}{g_i(z)} e_i}, \quad g_i(z) \in \mathbb{C}(z)^\times$$

**(G,q)-oper is Z-twisted** if it is equivalent to a constant element of  $G$   $Z \in H \subset H(z)$   $Z$  is regular semisimple. There are  $W_G$

$$A(z) = g(qz) Z g(z)^{-1}$$

Miura (G,q)-opers for each (G,q)-opers

**Z-twisted Miura (G,q)-oper** if gauge transform is from Borel

$$A(z) = v(qz) Z v(z)^{-1}, \quad v(z) \in B_+(z)$$

# Plücker Relations

$V_i^+$  irrep of  $G$  with highest weight  $\omega_i$       Line  $L_i \subset V_i$  stable under  $B_+$

Plucker relations: for two integral dominant weights  $L_{\lambda+\mu} \subset V_{\lambda+\mu}$  is the image of  $L_\lambda \otimes L_\mu \subset V_\lambda \otimes V_\mu$   
 under canonical projection  $V_\lambda \otimes V_\mu \longrightarrow V_{\lambda+\mu}$

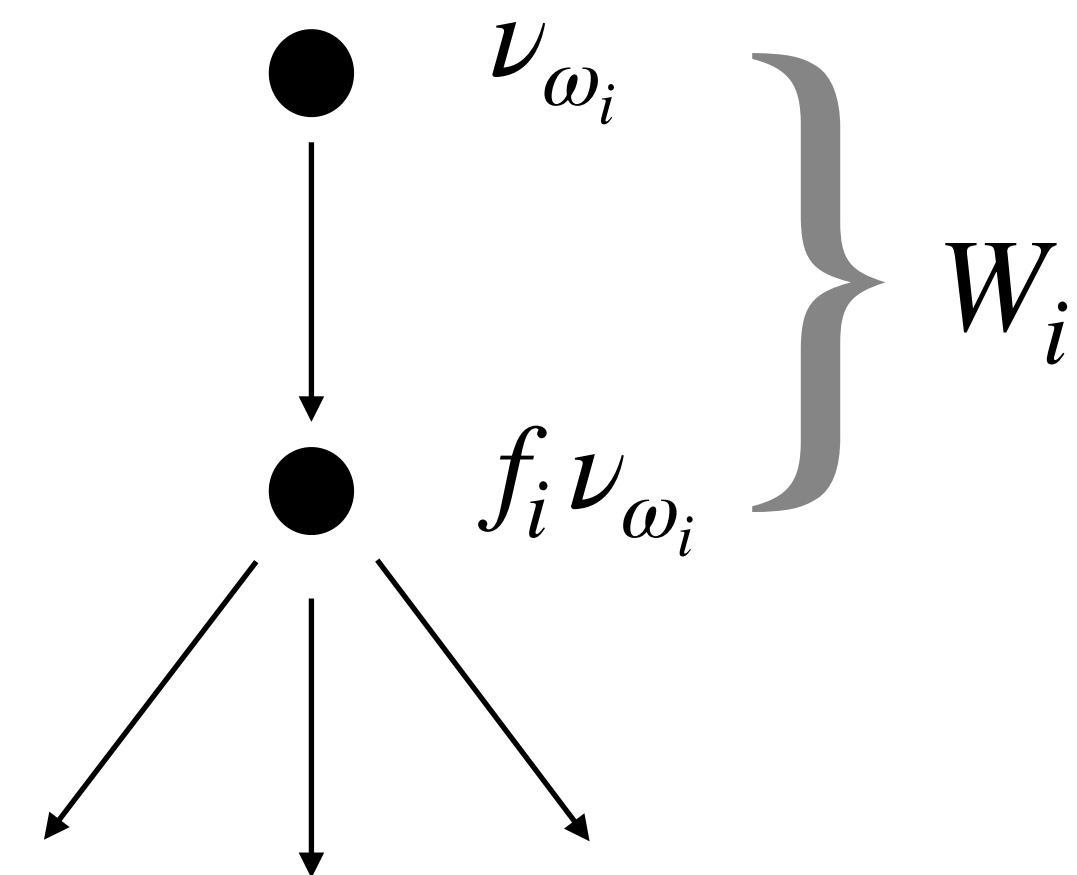
Conversely, for a collection of lines  $L_\lambda \subset V_\lambda$  satisfying Plucker relations  $\exists B \subset G$  such that  $L_\lambda$  is stabilized by  $B$  for all  $\lambda$

A choice of  $B$  is equivalent to a choice of  $B_+$ -torsor in  $G$

Let  $\nu_{\omega_i}$  be a generator of the line  $L_i \subset V_i$ . This is a vector of weight  $\omega_i$  wrt  $H \subset B_+$

The subspace of  $V_i$  of weight  $\omega_i - \alpha_i$  is one-dimensional and spanned  $f_i \cdot \nu_{\omega_i}$

Thus the 2d subspace spanned by  $\{\nu_{\omega_i}, f_i \cdot \nu_{\omega_i}\}$  is a  $B_+$ -invariant subspace of  $V_i$



# Miura-Plücker (G,q)-Operators

let  $(\mathcal{F}_G, A, \mathcal{F}_{B_-}, \mathcal{F}_{B_+})$  be a Miura  $(G, q)$ -oper with regular singularities  $\{\Lambda_i(z)\}_{i=1, \dots, r}$

Associated vector bundle  $\mathcal{V}_i = \mathcal{F}_{B_+} \times_{B_+} V_i = \mathcal{F}_G \times_G V_i$  contains rank-two subbundle  $\mathcal{W}_i = \mathcal{F}_{B_+} \times_{B_+} W_i$

associated to  $W_i \subset V_i$ , and  $\mathcal{W}_i$  in turn contains a line subbundle  $\mathcal{L}_i = \mathcal{F}_{B_+} \times_{B_+} L_i$

Using structure theorems we obtain  $r$  Miura  $(GL(2), q)$ -operators

$$A_i(z) = \begin{pmatrix} g_i(z) & \Lambda_i(z) \prod_{j>i} g_j(z)^{-a_{ji}} \\ 0 & g_i^{-1}(z) \prod_{j \neq i} g_j(z)^{-a_{ji}} \end{pmatrix}$$

**Z-twisted Miura-Plücker (G,q)-oper** is meromorphic Miura  $(G, q)$ -oper on  $P^1$  such that for each Miura  $(GL(2), q)$ -oper

$$A_i(z) = v(zq)Zv(z)^{-1}|_{\mathcal{W}_i} = v_i(zq)Z_i v_i(z)^{-1}$$

where  $v_i(z) = v(z)|_{\mathcal{W}_i}$  and  $Z_i = Z|_{\mathcal{W}_i}$



# QQ-System

**Theorem:** *There is a one-to-one correspondence between the set of nondegenerate  $Z$ -twisted Miura-Plücker  $(G, q)$ -opers and the set of nondegenerate polynomial solutions of the QQ-system*

$$\tilde{\xi}_i Q_-^i(z) Q_+^i(qz) - \xi_i Q_-^i(qz) Q_+^i(z) = \Lambda_i(z) \prod_{j>i} [Q_+^j(qz)]^{-a_{ji}} \prod_{j<i} [Q_+^j(z)]^{-a_{ji}}, \quad i = 1, \dots, r,$$

$$\tilde{\xi}_i = \zeta_i \prod_{j>i} \zeta_j^{a_{ji}}, \quad \xi_i = \zeta_i^{-1} \prod_{j<i} \zeta_j^{-a_{ji}}$$

Proof uses

$$v(z) = \prod_{i=1}^r y_i(z)^{\check{\alpha}_i} \prod_{i=1}^r e^{-\frac{Q_-^i(z)}{Q_+^i(z)} e_i} \dots, \quad g_i(z) = \zeta_i \frac{Q_+^i(qz)}{Q_+^i(z)}.$$

# SL(2) Example

$$A(z) = \begin{pmatrix} g(z) & \Lambda(z) \\ 0 & g(z)^{-1} \end{pmatrix}$$

Z-twisted q-oper condition

$$A(z) = v(zq)Zv(z)^{-1}, \quad Z = \begin{pmatrix} \zeta & 0 \\ 0 & \zeta^{-1} \end{pmatrix}$$

Gauge transformation reads

$$v(z) = \begin{pmatrix} y(z) & 0 \\ 0 & y(z)^{-1} \end{pmatrix} \begin{pmatrix} 1 & -\frac{Q_-(z)}{Q_+(z)} \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} y(z) & -y(z)\frac{Q_-(z)}{Q_+(z)} \\ 0 & y(z)^{-1} \end{pmatrix}$$

We find

$$g(z) = \zeta y(zq)y(z)^{-1}$$

$$\Lambda(z) = y(z)y(zq) \left( \zeta \frac{Q_-(z)}{Q_+(z)} - \zeta^{-1} \frac{Q_-(zq)}{Q_+(zq)} \right)$$

The q-oper condition becomes the **SL(2) QQ-system**

$$\zeta Q_-(z)Q_+(zq) - \zeta^{-1} Q_-(zq)Q_+(z) = \Lambda(z)$$

To get Bethe equations

$$Q_+(z) = \prod_{k=1}^m (z - w_k)$$

evaluate at roots of Q

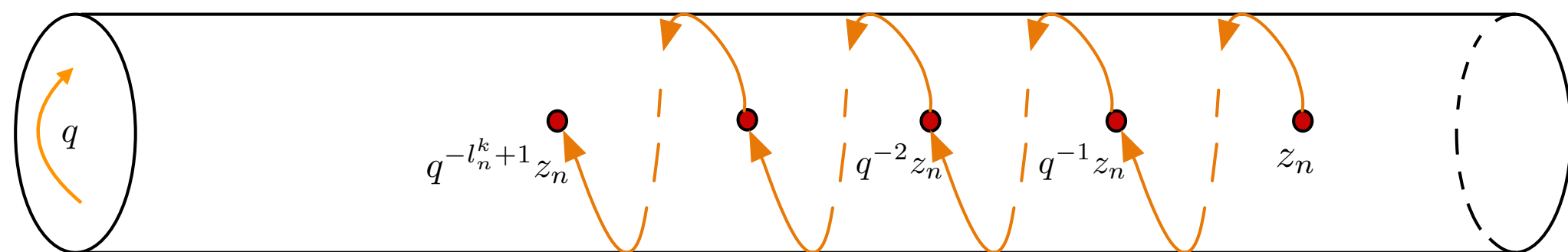
$$\frac{\Lambda(w_k)}{\Lambda(q^{-1}w_k)} = -\zeta^2 \frac{Q_+(qw_k)}{Q_+(q^{-1}w_k)}, \quad k = 1, \dots, m.$$

Singularities

$$\Lambda(z) = \prod_{p=1}^L \prod_{j_p=0}^{r_p-1} (z - q^{-j_p} z_p)$$

**XXZ Bethe equations**

$$q^r \prod_{p=1}^L \frac{w_k - q^{1-r_p} z_p}{w_k - q z_p} = -\zeta^2 q^m \prod_{j=1}^m \frac{qw_k - w_j}{w_k - qw_j}, \quad k = 1, \dots, m.$$



# XXZ Bethe Ansatz Equations for G

roots of Q+

$$\frac{Q_+^i(qw_i^k)}{Q_+^i(q^{-1}w_i^k)} \prod_j \zeta_j^{a_{ji}} = - \frac{\Lambda_i(w_k^i) \prod_{j>i} [Q_+^j(qw_k^i)]^{-a_{ji}} \prod_{j<i} [Q_+^j(w_k^i)]^{-a_{ji}}}{\Lambda_i(q^{-1}w_k^i) \prod_{j>i} [Q_+^j(w_k^i)]^{-a_{ji}} \prod_{j<i} [Q_+^j(q^{-1}w_k^i)]^{-a_{ji}}}$$

Space of nondegenerate solutions of  
QQ-system for G

Nondegenerate **Z-twisted Miura-Plucker** (G,q)-opers  
with regular singularities



Space of nondegenerate solutions of  
XXZ for G

?

?

Nondegenerate **Z-twisted Miura** (G,q)-opers  
with regular singularities

# Quantum Bäcklund Transformation

**Theorem:** Consider the following q-gauge transformation

$$A \mapsto A^{(i)} = e^{\mu_i(qz)f_i} A(z) e^{-\mu_i(z)f_i}, \quad \text{where} \quad \mu_i(z) = \frac{\prod_{j \neq i} [Q_+^j(z)]^{-a_{ji}}}{Q_+^i(z) Q_-^i(z)}$$

changes the set of Q-functions

$$\begin{aligned} Q_+^j(z) &\mapsto Q_+^j(z), & j \neq i, & & \{\tilde{Q}_+^j\}_{j=1, \dots, r} &= \{Q_+^1, \dots, Q_+^{i-1}, Q_-^i, Q_+^{i+1}, \dots, Q_+^r\} \\ Q_+^i(z) &\mapsto Q_-^i(z), & Z &\mapsto s_i(Z) & \{\tilde{z}_j\}_{j=1, \dots, r} &= \{z_1, \dots, z_{i-1}, z_i^{-1} \prod_{j \neq i} z_j^{-a_{ji}}, \dots, z_r\} \end{aligned}$$

Now the strategy is to successively apply Bäcklund transformations according to the reduced decomposition of the element of the Weyl group

Consider longest element  $w_0 = s_{i_1} \dots s_{i_\ell}$

**Theorem:** Every Z-twisted Miura-Plucker (G,q)-oper is Z-twisted Miura (G,q)-oper

The proof based on properties of double Bruhat cells addresses the existence of the diagonalizing element  $v(z)$  (to be constructed later)

# Generalized Wronskians

Consider big cell in Bruhat decomposition

$$G_0 = N_- H N_+ \\ g = n_- h n_+$$

$$V_i^+ \text{ irrep of } G \text{ with highest weight } \omega_i \\ h\nu_{\omega_i}^+ = [h]^{\omega_i} \nu_{\omega_i}^+$$

Define **principal minors** for group element  $g$

$$\Delta^{\omega_i}(g) = [h]^{\omega_i}, \quad i = 1, \dots, r$$

For  $SL(N)$  they are standard minors of matrices

Then **generalized minors** are regular functions on  $G$

$$\Delta_{u\omega_i, v\omega_i}(g) = \Delta^{\omega_i}(\tilde{u}^{-1} g \tilde{v}) \quad u, v \in W_G.$$

**Proposition 4.5.** *For a  $W$ -generic  $Z$ -twisted Miura-Plücker  $(G, q)$ -oper with  $q$ -connection  $A(z) = v(qz)Zv(z)^{-1}$ , where  $v(z) \in B_-(z)$  we have the following relation:*

$$(4.5) \quad \Delta_{w \cdot \omega_i, \omega_i}(v^{-1}(z)) = Q_+^{w, i}(z)$$

for any  $w \in W$ .

# Generalized Minors and QQ-system

The set of generalized minors  $\{\Delta_{w \cdot \omega_i, \omega_i}\}_{w \in W; i=1, \dots, r}$  creates a set of coordinates on  $G/B^+$ , known as *generalized Plücker coordinates*. In particular, the set of zeroes of each of  $\Delta_{w \cdot \omega_i, \omega_i}$  is a uniquely and unambiguously defined hypersurface in  $G/B$ .

**Proposition** For a  $W$ -generic  $Z$ -twisted Miura-Plücker  $(G, q)$ -oper with  $q$ -connection  $A(z) = v(qz)Zv(z)^{-1}$ , where  $v(z) \in B_-(z)$  we have the following relation:

$$\Delta_{w \cdot \omega_i, \omega_i}(v^{-1}(z)) = Q_+^{w, i}(z)$$

for any  $w \in W$ .

Proof: Since  $\Delta^{\omega_i}(v^{-1}(z)) = Q_+^i(z)$  Diagonalizing gauge transformation  $v^{-1}(z) = \prod_{i=1}^r e^{\frac{Q_-^i(z)}{Q_+^i(z)} f_i} \prod_{i=1}^r [Q_+^i(z)]^{\check{\alpha}_i} \dots$

$$v^{-1}(z)\nu_{\omega_i}^+ = Q_+^i(z)\nu_{\omega_i}^+ + Q_-^i(z)f_i\nu_{\omega_i}^+ + \dots$$

# Generalized Wronskians

The approach is similar to Miura-Plucker q-Operators

Let  $\nu_{\omega_i}^+$  be a generator of the line  $L_i^+ \subset V_i^+$   $V_i^+$  irrep of  $G$  with highest weight  $\omega_i$

The subspace  $L_{c,i}^+$  of  $V_i$  of weight  $c^{-1} \cdot \omega_i$  is one-dimensional and is spanned by  $s^{-1}\nu_{\omega_i}^+$

Associated vector bundle  $\mathcal{V}_i^+ = \mathcal{F}_{B_+} \times_{B_+} V_i^+ = \mathcal{F}_G \times_G V_i^+$

Contains line subbundles  $\mathcal{L}_i^+ = \mathcal{F}_H \times_H L_i^+$ ,  $\mathcal{L}_{c,i}^+ = \mathcal{F}_H \times_H L_{c,i}^+$

Define **generalized Wronskian** as quadruple  $(\mathcal{F}_G, \mathcal{F}_{B_+}, \mathcal{G}, Z)$

$\mathcal{G}$  is a meromorphic section of a principle bundle  $\mathcal{F}_G$

s.t. for sections  $\{v_i^+, v_{c,i}^+\}_{i=1,\dots,r}$  of line bundles  $\{\mathcal{L}_i^+, \mathcal{L}_{c,i}^+\}_{i=1,\dots,r}$  on  $U \cap M_q^{-1}(U)$

$$\mathcal{G}^q \cdot v_i^+ = Z \cdot \mathcal{G} \cdot v_{c,i}^+$$

# Adding Singularities

Effectively the above definition means that the Wronskian, written as an element of  $G(z)$ , satisfies

$$Z^{-1} \mathcal{G}(qz) \nu_{\omega_i}^+ = \mathcal{G}(z) \cdot s_{\phi}(z)^{-1} \cdot \nu_{\omega_i}^+$$

$$s_{\phi}(z) = \prod_i \phi_i^{-\check{\alpha}_i} s_i$$

Define **generalized Wronskian with regular singularities** if

$$s_{\Lambda}(z)^{-1} = \prod_i^{\text{inv}} s_i \Lambda_i^{\check{\alpha}_i}$$



# q-Operators and q-Wronskians

## Theorem 1:

Nondegenerate generalized q-Wronskians  
with regular singularities  $\{\Lambda_i\}_{i=1,\dots,r}$



Nondegenerate Z-twisted Miura (G,q)-opers  
with regular singularities  $\{\Lambda_i\}_{i=1,\dots,r}$

## Theorem 2:

*For a given Z-twisted (G,q)-Miura oper, there exists a unique generalized q-Wronskian*

$$\mathcal{W}(z) \in B_-(z)w_0B_-(z) \cap B_+(z)w_0B_+(z) \subset G(z),$$

*satisfying the system of equations*

$$\begin{aligned} \mathcal{W}(q^{k+1}z)\nu_{\omega_i}^+ &= Z^k \mathcal{W}(z)s^{-1}(z)s^{-1}(qz) \dots s^{-1}(q^k z)\nu_{\omega_i}^+, \\ i &= 1, \dots, r, \quad k = 0, 1, \dots, h-1, \end{aligned}$$

*where h is the Coxeter number of G.*

# Examples: SL(2)

$$\mathcal{W}(qz)\nu_{\omega}^+ = Z\mathcal{W}(z)s^{-1}(z)\nu_{\omega}^+$$

$$s^{-1}(z) = \tilde{s}^{-1}\Lambda(z)^{\check{\alpha}} = \begin{pmatrix} 0 & \Lambda(z)^{-1} \\ \Lambda(z) & 0 \end{pmatrix}, \quad \nu_{\omega}^+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad Z = \begin{pmatrix} \zeta & 0 \\ 0 & \zeta^{-1} \end{pmatrix}$$

In terms of Q-polynomials

$$\mathcal{W}(z) = \begin{pmatrix} Q^+(z) & \zeta^{-1}\Lambda(z)^{-1}Q_+(qz) \\ Q^-(z) & \zeta\Lambda(z)^{-1}Q^-(qz) \end{pmatrix}$$

$$\zeta Q^+(z)Q^-(qz) - \zeta^{-1}Q^+(qz)Q^-(z) = \Lambda(z)$$

is equivalent to  $\det \mathcal{W}(z) = 1$ .

# The qDE/IM Correspondence

[PK, Frenkel, Zeitlin]

**Proposition 5.1.** *i) In any evaluation representation  $\bar{V}_k^i$  there is a system of formal solutions of the linear problem*

$$\varphi_k^{i,v_s}(x) = x^{-\langle \ell, \lambda_s \rangle} (v_s + x\tilde{v}_s(x)),$$

where  $\tilde{v}_s(x) \in V[E, x^M][[x]]$ . Here  $\lambda_s = \text{wt}(v_s)$ .

*ii) The following transformation properties hold:*

$$\omega^{-k\rho^\vee} \varphi_{k'}^{i,v_s}(\omega^k x, \Omega^k E) = \omega^{-k\langle \ell + \rho^\vee, \lambda_s \rangle} \varphi_{k+k'}^{i,v_s}(x)$$

**Proposition 5.3.** *The solutions  $\Psi_{-k+1/2}^i$  have the following form:*

$$\Psi_{-k+1/2}^{(i)}(x, E) = \sum_{s=1}^{\dim V} Q_{v_s}^i(q^{-k+1/2} E) \varphi_{1/2}^{i, Z^k v_s}(x, E)$$

# Main Theorem

Affine  $\mathfrak{g}$  oper on formal disk



$(G, q)$ -oper on a disc

**Theorem:** 1) In case if  $M \in \mathbb{Z}_+$  ( $q^{M+1} = 1$ ), the monodromy matrix is represented by regular semisimple element  $Z^{(M+1)h^\vee}$  in the basis of  $\varphi_0^{i, v_s}$ :

$$\varphi_k^{i, v_s}(e^{2\pi i} x, E) = (-1)^{2\langle \rho^\vee, \omega_i \rangle} \varphi_k^{i, Z^{(M+1)h^\vee} v_s}(x, E),$$

2) For  $\Psi_{-k+1/2}^{(i)}(x, E)$  solutions, the monodromy operator can be expressed as follows:

$$\Psi_{-k+1/2}^{(i)}(e^{2\pi i} x, E) = (-1)^{2\langle \rho^\vee, \omega_i \rangle} W_e(x, E) Z^{(M+1)h^\vee} W_e^{-1}(x, E) \Psi_{-k+1/2}^{(i)}(x, E)$$

3) The Monodromy matrix is conjugated to the following operator

$$\left[ A(q^M E) A(q^{M-1} E) \dots A(E) \right]^{h^\vee} = v(E) Z^{(M+1)h^\vee} v(E)^{-1},$$

where  $A(E) = v(qE) Z v(E)^{-1}$  is the Miura  $(G, q)$ -oper connection, defined by the QQ-system.