

# FROM ELIPTIC HYPERGEOMETRIC INTEGRALS TO COMPLEX HYPERGEOMETRIC FUNCTIONS

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**Abstract.** Elliptic hypergeometric integrals are top transcendental special functions of hypergeometric type. They have found applications in  $4d$  supersymmetric quantum fields theories (superconformal indices), in integrable systems (wave functions in quantum  $N$ -body problems) and  $2d$  statistical mechanics (partition functions). I will describe how these integrals can be degenerated in a chain of limits to complex hypergeometric functions related to the representation theory of  $SL(2, \mathbb{C})$ .

## An elliptic analogue of the Euler-Gauss ${}_2F_1$ -function

V.S., 2003

$$V(t_1, \dots, t_8; p, q) = \frac{(p; p)_\infty (q; q)_\infty}{4\pi i} \int_{\mathbb{T}} \frac{\prod_{j=1}^8 \Gamma(t_j x^{\pm 1}; p, q) dx}{\Gamma(x^{\pm 2}; p, q) x},$$

where  $\prod_{j=1}^8 t_j = (pq)^2$ ,  $|t_j| < 1$  and  $(z; p)_\infty = \prod_{j=0}^{\infty} (1 - zp^j)$ ,

$$\Gamma(z; p, q) = \prod_{j,k=0}^{\infty} \frac{1 - z^{-1} p^{j+1} q^{k+1}}{1 - zp^j q^k}, \quad |q|, |p| < 1,$$

$$\Gamma(qz; p, q) = \theta(z; p) \Gamma(z; p, q), \quad \theta(z; p) := (z; p)_\infty (pz^{-1}; p)_\infty.$$

Inversion relation:  $\Gamma\left(\frac{pq}{z}; p, q\right) \Gamma(z; p, q) = 1.$

### The elliptic beta integral

$$V(t_1, \dots, t_7, \frac{pq}{t_7}; p, q) = \prod_{1 \leq j < k \leq 6} \Gamma(t_j t_k; p, q).$$

V.S., 2000

## Applications

- $V$  = a special eigenfunction of  $N = 1$  van Diejen model Hamiltonian (a  $BC_N$  version of the Ruijsenaars model). Similar picture for  $BC_N$  extension of the  $V$ -function. V.S., 2004
- Superconformal indices (counting BPS states) of  $4d$  SUSY theories with  $G = SU(2)$ ,  $F = SU(8)$  or  $SU(6)$  and chiral superfields in the fundamental representation. The elliptic beta integral evaluation proves the Seiberg duality conjecture in the BPS-sector. Dolan, Osborn, 2008
- In statistical mechanics: the star-triangle, star-star relations, top solutions of the Yang-Baxter equation guaranteeing integrability of the corresponding  $2d$  Ising type models and  $1d$  Heisenberg type spin chains.  
**NB.** Seiberg duality = integrability.  
V.S., 2003; Bazhanov, Sergeev, 2010; V.S., 2010; Derkachov, V.S., 2012.

**The trigonometric limit:  $p \rightarrow 0$  for fixed  $q$**

For fixed  $z$  and  $q$ :

$$\Gamma(z; 0, q) = \frac{1}{(z; q)_\infty}, \quad \lim_{p \rightarrow 0} \Gamma(p^\alpha z; p, q) = 1, \quad 0 < \alpha < 1,$$

$$\lim_{p \rightarrow 0} \Gamma(pz; p, q) = (qz^{-1}; p, q)_\infty.$$

Cleaning the  $V$ -balancing condition:

$$t_j = p^{\alpha_j} g_j, \quad \sum_{j=1}^8 \alpha_j = 2, \quad \prod_{j=1}^8 g_j = q^2,$$

for fixed  $q$  and  $g_j$  the limit  $p \rightarrow 0$  becomes well defined. E.g.,

$$\alpha_j = 0, \quad j = 1, \dots, 6, \quad \alpha_7 = \alpha_8 = 1, \quad \Rightarrow$$

$$\lim_{p \rightarrow 0} V(\underline{t}; p, q) = \frac{(q; q)_\infty}{4\pi i} \int_{\mathbb{T}} \frac{(z^{\pm 2}; q)_\infty \prod_{k=7,8} (qg_k^{-1} z^{\pm 1}; q)_\infty dz}{\prod_{k=1}^6 (g_k z^{\pm 1}; q)_\infty z}.$$

Special choice of  $g_j$ :  $6j$ -symbols for  $sl_q(2, \mathbb{R})$ , or Askey-Wilson functions.

Goes down to Euler integral representation for  ${}_2F_1$ -function.

**Rational degeneration:**  $q \rightarrow 1$  (the limit  $q = 0$  is trivial).

Jackson's  $q$ -gamma function

$$\Gamma_q(u) := \frac{(q; q)_\infty}{(q^u; q)_\infty} (1 - q)^{1-u}, \quad |q| < 1,$$

$$\Gamma_q(u + 1) = \frac{1 - q^u}{1 - q} \Gamma_q(u), \quad \lim_{q \rightarrow 1} \Gamma_q(u) = \Gamma(u).$$

Koornwinder (1990): the limit  $q \rightarrow 1^-$  is uniform on compacta

Rains (2006): OK for  $|q| \rightarrow 1$  from inside  $\mathbb{T}$  under fixed angle

Substitute  $z = q^u$ ,  $g_j = q^{a_j}$ ,  $\sum_{j=1}^8 a_j = 2$ , infinite products  $\rightarrow 1/\Gamma_q(u)$   
 $\Rightarrow$  Mellin-Barnes type integral

$$\lim_{q \rightarrow 1} \lim_{p \rightarrow 0} V(\underline{t}; p, q) = \frac{\kappa(q)}{4\pi i} \int_{-i\infty}^{i\infty} \frac{\prod_{j=1}^6 \Gamma(a_j \pm u)}{\Gamma(\pm 2u) \prod_{j=7,8} \Gamma(1 - a_j \pm u)} du,$$

where  $\kappa(q)$  is a diverging factor

$$\kappa(q) = (q; q)^5 (1 - q)^{-6} \log q^{-1}.$$

Goes down to the Barnes representation for  ${}_2F_1$ -function.

## Hyperbolic degeneration

Parametrize

$$t_j = e^{-2\pi v g_j}, \quad z = e^{-2\pi v u}, \quad p = e^{-2\pi v \omega_1}, \quad q = e^{-2\pi v \omega_2}.$$

In the limit  $v \rightarrow 0^+$ ,

Ruijsenaars, 1997

$$\Gamma(e^{-2\pi v u}; e^{-2\pi v \omega_1}, e^{-2\pi v \omega_2}) \underset{v \rightarrow 0^+}{=} e^{-\pi \frac{2u - \omega_1 - \omega_2}{12v\omega_1\omega_2}} \gamma^{(2)}(u; \omega_1, \omega_2).$$

Faddeev's (1994) modular dilogarithm, or hyperbolic gamma function

$$\gamma^{(2)}(u; \omega) = \gamma^{(2)}(u; \omega_1, \omega_2) := e^{-\frac{\pi i}{2} B_{2,2}(u; \omega)} \gamma(u; \omega),$$

second order multiple Bernoulli polynomial

$$B_{2,2}(u; \omega) = \frac{1}{\omega_1 \omega_2} \left( \left( u - \frac{\omega_1 + \omega_2}{2} \right)^2 - \frac{\omega_1^2 + \omega_2^2}{12} \right),$$

$$\gamma(u; \omega) := \frac{(\tilde{\mathbf{q}} e^{2\pi i \frac{u}{\omega_1}}; \tilde{\mathbf{q}})_\infty}{(e^{2\pi i \frac{u}{\omega_2}}; \mathbf{q})_\infty} = \exp \left( - \int_{\mathbb{R}+i0} \frac{e^{ux}}{(1 - e^{\omega_1 x})(1 - e^{\omega_2 x})} \frac{dx}{x} \right),$$

$$\mathbf{q} = e^{2\pi i \frac{\omega_1}{\omega_2}}, \quad \tilde{\mathbf{q}} = e^{-2\pi i \frac{\omega_2}{\omega_1}}.$$

It is well defined for  $\omega_1, \omega_2 > 0$  (i.e.,  $|\mathbf{q}| = 1$ ) and  $0 < \operatorname{Re}(u) < \omega_1 + \omega_2$ .

Rains (2006): this limit is uniform on compacta!

Dedekind function:  $\eta(\tau) = q^{1/24}(q; q)_\infty$ ,  $q = e^{2\pi i\tau}$ ,

$$\eta(-1/\tau) = \sqrt{-i\tau} \eta(\tau).$$

It yields the asymptotics of  $(q; q)_\infty$  for  $\tau \rightarrow 0$  (or  $q \rightarrow 1$ ).

Then,

Rains, 2006

$$V(e^{-2\pi v g_k}; e^{-2\pi v \omega_1}, e^{-2\pi v \omega_2}) \underset{v \rightarrow 0^+}{=} e^{\frac{\pi}{4v} \left( \frac{1}{\omega_1} + \frac{1}{\omega_2} \right)} I_h(\underline{g}),$$

$$I_h(\underline{g}) = \int_{-i\infty}^{i\infty} \frac{\prod_{j=1}^8 \gamma^{(2)}(g_j \pm z; \omega_1, \omega_2)}{\gamma^{(2)}(\pm 2z; \omega_1, \omega_2)} \frac{dz}{2i\sqrt{\omega_1\omega_2}},$$

$$\operatorname{Re}(g_j) > 0, \quad \sum_{j=1}^8 g_j = 2(\omega_1 + \omega_2).$$

Special choice of parameters  $\Rightarrow$  eigenfunction of  $N = 1$  Hamiltonian of hyperbolic Ruijsenaars model (Ruijsenaars, 1994)  $\propto 6j$  symbols for the Faddeev modular double  $sl_q(2, \mathbb{R}) \times sl_{\tilde{q}}(2, \mathbb{R})$  (Ponsot, Tschner, 2001)

Remind

$$q = e^{2\pi i \frac{\omega_1}{\omega_2}}, \quad \tilde{q} = e^{-2\pi i \frac{\omega_2}{\omega_1}}.$$

Then,

$$\begin{aligned} \gamma^{(2)}(\omega_1 x; \omega) &= \Gamma_q(x) \frac{e^{-\frac{\pi i}{2} B_{2,2}(\omega_1 x; \omega)}}{(q; q)_\infty (1-q)^{1-x}} (\tilde{q} e^{2\pi i x}; \tilde{q})_\infty \\ &\underset{\omega_1 \rightarrow 0}{\approx} \frac{\Gamma(x)}{\sqrt{2\pi}} \left( \frac{\omega_2}{2\pi\omega_1} \right)^{\frac{1}{2}-x}, \quad q \rightarrow 1^-, \quad \tilde{q} \rightarrow 0. \end{aligned} \quad \text{Ruijsenaars, 1997}$$

Replace  $g_{7,8} \rightarrow g_{7,8} + \omega_1$  and apply the inversion formula

$$\gamma^{(2)}(x; \omega) \gamma^{(2)}(\omega_1 + \omega_2 - x; \omega) = 1.$$

Then the limit  $\omega_1 \rightarrow 0 \Rightarrow$  the previous plain hypergeometric integral.



## Complex hypergeometric functions

Euler's beta integral:

$$\int_0^1 x^{\alpha-1}(1-x)^{\beta-1}dx = \frac{\Gamma(\alpha)\Gamma(\beta)}{\Gamma(\alpha+\beta)}, \quad \operatorname{Re}(\alpha), \operatorname{Re}(\beta) > 0.$$

Take  $\alpha, \alpha' \in \mathbb{C}$ ,  $\alpha - \alpha' = n_\alpha \in \mathbb{Z}$  and

$$[z]^\alpha := z^\alpha \bar{z}^{\alpha'} = |z|^{2\alpha'} z^{n_\alpha}, \quad \int_{\mathbb{C}} d^2z := \int_{\mathbb{R}^2} d(\operatorname{Re} z) d(\operatorname{Im} z),$$

$\bar{z}$  is a complex conjugate of  $z$ .

Then, the complex beta integral is

Gelfand, Graev, Vilenkin, 1962

$$\int_{\mathbb{C}} [w - z_1]^{\alpha-1} [z_2 - w]^{\beta-1} \frac{d^2w}{\pi} = \frac{\Gamma(\alpha)\Gamma(\beta)\Gamma(1 - \alpha' - \beta')}{\Gamma(1 - \alpha')\Gamma(1 - \beta')\Gamma(\alpha + \beta)} [z_2 - z_1]^{\alpha+\beta-1}$$

Complex gamma function

$$\mathbf{\Gamma}(x, n) = \mathbf{\Gamma}(\alpha|\alpha') := \frac{\Gamma(\alpha)}{\Gamma(1 - \alpha')} = \frac{\Gamma(\frac{n+ix}{2})}{\Gamma(1 + \frac{n-ix}{2})},$$

$$\alpha = \frac{n + ix}{2}, \quad \alpha' = \frac{-n + ix}{2}, \quad x \in \mathbb{C}, \quad n \in \mathbb{Z}.$$

From the reflection relation  $\Gamma(x)\Gamma(1-x) = \frac{\pi}{\sin \pi x} \Rightarrow$

$$\mathbf{\Gamma}(\alpha|\alpha') = (-1)^{\alpha-\alpha'} \mathbf{\Gamma}(\alpha'|\alpha), \quad \mathbf{\Gamma}(x, -n) = (-1)^n \mathbf{\Gamma}(x, n),$$

$$\mathbf{\Gamma}(\alpha|\alpha') \mathbf{\Gamma}(1-\alpha|1-\alpha') = (-1)^{\alpha-\alpha'}, \quad \mathbf{\Gamma}(x, n) \mathbf{\Gamma}(-x-2i, n) = 1.$$

$$\mathbf{\Gamma}(\alpha+1|\alpha') = \mathbf{\Gamma}(x-i, n+1) = \alpha \mathbf{\Gamma}(\alpha|\alpha'), \quad \mathbf{\Gamma}(\alpha|\alpha'+1) = \mathbf{\Gamma}(x-i, n-1) = -\alpha' \mathbf{\Gamma}(\alpha|\alpha').$$

$$\begin{aligned} \Rightarrow \int_{\mathbb{C}} [w-z_1]^{\alpha-1} [z_2-w]^{\beta-1} \frac{d^2 w}{\pi} \\ = \frac{\mathbf{\Gamma}(\alpha|\alpha') \mathbf{\Gamma}(\beta|\beta')}{\mathbf{\Gamma}(\alpha+\beta|\alpha'+\beta')} [z_2-z_1]^{\alpha+\beta-1} = \frac{\mathbf{\Gamma}(\alpha, \beta, \gamma)}{[z_1-z_2]^\gamma}, \end{aligned}$$

where  $\alpha + \beta + \gamma = \alpha' + \beta' + \gamma' = 1$  and

$$\mathbf{\Gamma}(\alpha_1, \dots, \alpha_k) := \mathbf{\Gamma}(\alpha_1|\alpha'_1) \cdots \mathbf{\Gamma}(\alpha_k|\alpha'_k).$$

Inversion  $w \rightarrow w^{-1}$ ,  $z_1 \rightarrow z_1^{-1}$ ,  $z_2 \rightarrow z_2^{-1}$  and shifts  $w \rightarrow w - z_3$ ,  $z_1 \rightarrow z_1 - z_3$ ,  $z_2 \rightarrow z_2 - z_3 \Rightarrow$  the star-triangle relation:

$$\int_{\mathbb{C}} [z_1 - w]^{\alpha-1} [z_2 - w]^{\beta-1} [z_3 - w]^{\gamma-1} \frac{d^2 w}{\pi} \\ = \frac{\mathbf{\Gamma}(\alpha, \beta, \gamma)}{[z_3 - z_2]^\alpha [z_1 - z_3]^\beta [z_2 - z_1]^\gamma}, \quad \alpha + \beta + \gamma = 1.$$

Nice applications to Feynman diagrams and non-compact spin chains:  
Vasil'ev, Pismak, Khonkonen, Derkachov, Manashov, Valinevich, 1981 - ...

Complex hypergeometric functions:

Emergence in 2d conformal field theory: Dotsenko, Fateev, 1985

Complex Selberg integral: Aomoto, 1987

Mellin-Barnes representations of complex hypergeometric functions:

$3j$ -symbols of  $SL(2, \mathbb{C})$  for unitary principal series (real  $x$ ): Naimark, 1959

$6j$ -symbols of  $SL(2, \mathbb{C})$ : Ismagilov, 2006; Derkachov, V.S., 2017

Rigorous spectral analysis: Molchanov, Neretin, 2018

## From hyperbolic integrals to complex hypergeometric functions

Formally: Bazhanov, Mangazeev, Sergeev, 2008; Kels, 2014;

Complete rigorous consideration: Sarkissian, V.S, 2019

Take

$$\gamma(u; \omega_1, \omega_2) = \frac{(e^{2\pi i \frac{u}{\omega_1}} e^{-2\pi i \frac{\omega_2}{\omega_1}}; e^{-2\pi i \frac{\omega_2}{\omega_1}})_\infty}{(e^{2\pi i \frac{u}{\omega_2}}; e^{2\pi i \frac{\omega_1}{\omega_2}})_\infty}$$

and set

$$“b” = \sqrt{\frac{\omega_1}{\omega_2}} = i + \delta, \quad \delta \rightarrow 0^+ \quad (c_{CFT} \rightarrow 1).$$

Then  $\omega_1 + \omega_2 = 2\delta\sqrt{\omega_1\omega_2} + O(\delta^2) \rightarrow 0$  and

$$\sqrt{\frac{\omega_2}{\omega_1}} = -i + \delta + O(\delta^2), \quad \frac{\omega_1}{\omega_2} = -1 + 2i\delta + \delta^2, \quad \frac{\omega_2}{\omega_1} = -1 - 2i\delta + O(\delta^2).$$

Special choice of the argument  $u$ :

$$u = i\sqrt{\omega_1\omega_2}(n + x\delta), \quad n \in \mathbb{Z}, \quad x \in \mathbb{C}.$$

Then, for  $q = e^{2\pi i \frac{\omega_1}{\omega_2}} \rightarrow 1$

$$(e^{2\pi i \frac{u}{\omega_2}}; e^{2\pi i \frac{\omega_1}{\omega_2}})_\infty = (e^{-2\pi\delta(n+ix+\delta x)}; q)_\infty \underset{\delta \rightarrow 0^+}{=} \frac{(q; q)_\infty (1-q)^{1-\frac{n+ix}{2}+O(\delta)}}{\Gamma_q\left(\frac{n+ix}{2} + O(\delta)\right)}.$$

Similarly for  $\tilde{q} = e^{-2\pi i \frac{\omega_2}{\omega_1}} \rightarrow 1$

$$(e^{2\pi i \frac{u}{\omega_1}} e^{-2\pi i \frac{\omega_2}{\omega_1}}; e^{-2\pi i \frac{\omega_2}{\omega_1}})_\infty \underset{\delta \rightarrow 0^+}{=} \frac{(\tilde{q}; \tilde{q})_\infty (1 - \tilde{q})^{-\frac{n+ix}{2} + O(\delta)}}{\Gamma_{\tilde{q}} \left(1 + \frac{n-ix}{2} + O(\delta)\right)}.$$

As a result,

$$\begin{aligned} \gamma(u; \omega_1, \omega_2) &\underset{\delta \rightarrow 0^+}{=} \frac{\Gamma_q \left(\frac{n+ix}{2} + O(\log q)\right)}{\Gamma_{\tilde{q}} \left(1 + \frac{n-ix}{2} + O(\log \tilde{q})\right)} \frac{(\tilde{q}; \tilde{q})_\infty (1 - \tilde{q})^{-\frac{n+ix}{2} + O(\log \tilde{q})}}{(q; q)_\infty (1 - q)^{1 - \frac{n+ix}{2} + O(\log q)}}. \\ \frac{(\tilde{q}; \tilde{q})_\infty}{(q; q)_\infty} &= e^{\frac{\pi i}{12} \left(\frac{\omega_2}{\omega_1} + \frac{\omega_1}{\omega_2}\right)} \left(-i \frac{\omega_1}{\omega_2}\right)^{\frac{1}{2}} \underset{\delta \rightarrow 0^+}{=} e^{\frac{\pi i}{12}}. \end{aligned}$$

Finally, for  $\sqrt{\frac{\omega_1}{\omega_2}} = i + \delta$ , uniformly on the compacta

$$\gamma^{(2)}(i\sqrt{\omega_1 \omega_2}(n + x\delta); \omega_1, \omega_2) \underset{\delta \rightarrow 0^+}{\approx} (4\pi\delta)^{ix-1} e^{\frac{\pi i}{2} n^2} \mathbf{\Gamma}(x, n).$$

## Degeneration of integrals

Consider

$$\begin{aligned}
& \int_{-i\infty}^{i\infty} \Delta(z) \frac{dz}{i\sqrt{\omega_1\omega_2}} = \int_{-\infty}^{\infty} \Delta(i\sqrt{\omega_1\omega_2}x) dx, \quad x = \frac{z}{i\sqrt{\omega_1\omega_2}}, \\
& = \sum_{N \in \mathbb{Z}} \int_{N-1/2}^{N+1/2} \Delta(i\sqrt{\omega_1\omega_2}x) dx = \sum_{N \in \mathbb{Z}} \int_{-1/2}^{1/2} \Delta(i\sqrt{\omega_1\omega_2}(N+x)) dx \\
& = \sum_{N \in \mathbb{Z}} \int_{-1/2\delta}^{1/2\delta} \delta \Delta(i\sqrt{\omega_1\omega_2}(N+y\delta)) dy \underset{\delta \rightarrow 0^+}{=} \sum_{N \in \mathbb{Z}} \int_{-\infty}^{\infty} [\delta \Delta(i\sqrt{\omega_1\omega_2}(N+y\delta))] dy, \quad x = y\delta.
\end{aligned}$$

Apply this limit to the general univariate hyperbolic beta integral.

$$I_h(g_1, \dots, g_7, \omega_1 + \omega_2 - g_7) = \prod_{1 \leq j < k \leq 6} \gamma^{(2)}(g_j + g_k; \omega).$$

Parametrization of the integral kernel for  $I_h(\underline{g})$ :

$$z = i\sqrt{\omega_1\omega_2}(N + \delta y), \quad y \in \mathbb{C}, \quad N \in \mathbb{Z} + \nu, \quad \nu = 0, \frac{1}{2},$$

$$g_k = i\sqrt{\omega_1\omega_2}(N_k + \delta\alpha_k), \quad \alpha_k \in \mathbb{C}, \quad N_k \in \mathbb{Z} + \nu, \quad \nu = 0, \frac{1}{2},$$

The limit  $\delta \rightarrow 0^+ \Rightarrow$  the balancing condition

$$\sum_{k=1}^6 \alpha_k = -2i, \quad \sum_{k=1}^6 N_k = 0.$$

The parameter  $\nu = 0, 1/2$  emerges because only the sums  $N_k \pm N$  should be integers. Now,  $\delta \rightarrow 0^+$  limiting relations:

$$\prod_{k=1}^6 \gamma^{(2)}(g_k \pm z; \omega) \rightarrow \frac{(-1)^{2\nu}}{(4\pi\delta)^8} \prod_{k=1}^6 \Gamma(\alpha_k + y, N_k + N) \Gamma(\alpha_k - y, N_k - N),$$

$$\prod_{1 \leq j < k \leq 6} \gamma^{(2)}(g_j + g_k; \omega) \rightarrow \frac{(-1)^{2\nu}}{(4\pi\delta)^5} \prod_{1 \leq j < k \leq 6} \mathbf{\Gamma}(\alpha_j + \alpha_k, N_j + N_k),$$

$$\gamma^{(2)}(\pm 2z; \omega) \rightarrow \frac{(-1)^{2\nu}}{(4\pi\delta)^2} \frac{\Gamma(N + iy)}{\Gamma(1 + N - iy)} \frac{\Gamma(-N - iy)}{\Gamma(1 - N + iy)} = \frac{(4\pi\delta)^{-2}}{y^2 + N^2}.$$

The diverging factors  $(4\pi\delta)^{-5}$  on both sides cancel.

The final complex beta integral evaluation formula:

$$\frac{1}{8\pi} \sum_{N \in \mathbb{Z} + \nu} \int_{-\infty}^{\infty} (y^2 + N^2) \prod_{k=1}^6 \mathbf{\Gamma}(\alpha_k \pm y, N_k \pm N) dy$$

$$= \prod_{1 \leq j < k \leq 6} \mathbf{\Gamma}(\alpha_j + \alpha_k, N_j + N_k),$$

where  $\sum_{k=1}^6 \alpha_k = -2i$ ,  $\sum_{k=1}^6 N_k = 0$ ,  $N_k \in \mathbb{Z} + \nu$ ,  $\nu = 0, \frac{1}{2}$ , and

$$\mathbf{\Gamma}(x_1 \pm x_2, n_1 \pm n_2) := \mathbf{\Gamma}(x_1 + x_2, n_1 + n_2) \mathbf{\Gamma}(x_1 - x_2, n_1 - n_2).$$

Special cases (non-rigorously): Bazhanov, Mangazeev, Sergeev, 2008; Kels, 2014

General case (rigorously): Sarkissian, V.S., 2019; Derkachov, Manashov, 2019.



Degeneration of the  $W(E_7)$ -group transformation law for  $V$ -function: Seiberg duality for  $4d$  SCI on  $S^3 \times S^1 \rightarrow$  mirror symmetry of  $3d$  PF on  $S_b^3 \rightarrow 2d$  duality of what on  $\mathbb{C}$  ? vortices ?)

$$\begin{aligned}
& \sum_{N \in \mathbb{Z} + \nu} \int_{-\infty}^{\infty} (y^2 + N^2) \prod_{k=1}^8 \Gamma(\alpha_k \pm y, N_k \pm N) dy \\
&= (-1)^L \prod_{1 \leq j < k \leq 4} \Gamma(\alpha_j + \alpha_k, N_j + N_k) \prod_{5 \leq j < k \leq 8} \Gamma(\alpha_j + \alpha_k, N_j + N_k) \\
&\quad \times \sum_{N \in \mathbb{Z} + \mu} \int_{-\infty}^{\infty} (y^2 + N^2) \prod_{k=1}^4 \Gamma(\alpha_k \pm y - \frac{1}{2}X - i, N_k \pm N - \frac{1}{2}L) \\
&\quad \quad \quad \times \prod_{k=5}^8 \Gamma(\alpha_k \pm y + \frac{1}{2}X + i, N_k \pm N + \frac{1}{2}L) dy,
\end{aligned}$$

with  $X := \sum_{j=1}^4 \alpha_j$ ,  $L := \sum_{j=1}^4 N_j$  and the balancing

$$\sum_{k=1}^8 \alpha_k = -4i, \quad \alpha_k \in \mathbb{C}, \quad \sum_{k=1}^8 N_k = 0, \quad N_k \in \mathbb{Z} + \nu, \quad \nu = 0, \frac{1}{2}.$$

### Limits of parameters $g_k \rightarrow \infty$ .

Degenerate form of the hyperbolic beta integral

$$\int_{-i\infty}^{i\infty} \prod_{j=1}^3 \gamma^{(2)}(f_j + z; \omega) \gamma^{(2)}(h_j - z; \omega) \frac{dz}{i\sqrt{\omega_1 \omega_2}} = \prod_{j,k=1}^3 \gamma^{(2)}(f_j + h_k; \omega),$$

$\sum_{j=1}^3 (f_j + h_j) = \omega_1 + \omega_2$ . The limit  $\omega_1 + \omega_2 \rightarrow 0$  ( $b \rightarrow i$ ) yields

$$\begin{aligned} & \frac{i^{2\nu}}{4\pi} \sum_{N \in \mathbb{Z} + \nu} \int_{-\infty}^{\infty} \prod_{j=1}^3 \Gamma(s_j + y, N + N_j) \Gamma(t_j - y, N - M_j) dy \\ &= \prod_{j,k=1}^3 \Gamma(s_j + t_k, N_j + M_k), \quad \sum_{j=1}^3 (N_j + M_j) = 0, \quad \sum_{j=1}^3 (s_j + t_j) = -2i. \end{aligned}$$

For  $\nu = 0$  this is the Mellin-Barnes representation of the complex star-triangle relation.  
Derkachov, Manashov, Valinevich, 2017.

**6j-symbols** for unitary principal series of  $SL(2, \mathbb{C})$  = triple integral over complex plane (i.e. 6 real integrations) = infinite bilateral sum of Mellin-Barnes integrals.

Ismagilov, 2006; Derkachov, V.S., 2017

$$R_\ell(c, c') = \int d^2z \Psi_2(a_1, a_2, a_3 | \ell, c, z) \overline{\Psi_1(a_1, a_2, a_3 | \ell, c', z)},$$

$$\overline{\Psi_1(a_1, a_2, a_3 | \ell, c', z)} = \int \frac{d^2y}{[1-y]^{\frac{1-\ell-c'-a_3}{2}} [y]^{\frac{1+\ell-c'+a_3}{2}} [z-y]^{\frac{1-a_1+a_2+c'}{2}}},$$

$$\Psi_2(a_1, a_2, a_3 | \ell, c, z) = \frac{1}{[z]^{\frac{1+a_1-\ell+c}{2}}} \int \frac{d^2z_0}{[1-z_0]^{\frac{1+a_1+\ell+c}{2}} [z_0]^{\frac{1+a_2-a_3-c}{2}} [z-z_0]^{\frac{1-a_2+a_3-c}{2}}}.$$

Mellin-Barnes representation:

$$\begin{array}{llll} a_1 = N_1/2 + i\sigma_1 & a_3 = N_3/2 + i\sigma_3 & c = M_1/2 + i\rho_1 & s = N/2 + iu/2 \\ a_2 = N_2/2 + i\sigma_2 & l = N_4/2 + i\sigma_4 & c' = M_2/2 + i\rho_2 & \end{array}$$

$$\left\{ \begin{array}{l} \sigma_1, N_1 \quad \sigma_2, N_2 \quad | \quad \rho_1, M_1 \\ \sigma_3, N_3 \quad \sigma_4, N_4 \quad | \quad \rho_2, M_2 \end{array} \right\} = (-1)^{M_2 - N_2 + N_4} \frac{\pi^2 \Gamma(\sigma_1 - \sigma_2 + \rho_2 - i, A_1) \Gamma(\sigma_2 - \sigma_3 + \rho_1 - i, A_2)}{4 \Gamma(-\sigma_3 - \sigma_4 + \rho_2 - i, A_3)}$$

$$\times \Gamma(\sigma_1 + \sigma_4 + \rho_1 - i, A_4) \sum_{N \in \mathbb{Z}} \int \prod_{k=1}^4 \Gamma(R_k - u, S_k - N) \Gamma(U_k + u, T_k + N) du$$

where

$$\begin{aligned} R_1 &= -\sigma_1 + \sigma_2 - \rho_2 - i & U_1 &= -\rho_1 - \sigma_2 + \sigma_4 + \rho_2 & S_1 &= (-N_1 + N_2 - M_2)/2 \\ R_2 &= \sigma_1 + \sigma_2 - \rho_2 - i & U_2 &= \rho_1 - \sigma_2 + \sigma_4 + \rho_2 & S_2 &= (N_1 + N_2 - M_2)/2 \\ R_3 &= -\sigma_3 - \sigma_4 - \rho_2 - i & U_3 &= 0 & S_3 &= -(N_3 + N_4 + M_2)/2 \\ R_4 &= \sigma_3 - \sigma_4 - \rho_2 - i & U_4 &= 2\rho_2 & S_4 &= (N_3 - N_4 - M_2)/2 \\ T_1 &= (-M_1 - N_2 + N_4 + M_2)/2, & T_2 &= (M_1 - N_2 + N_4 + M_2)/2, & T_3 &= 0, & T_4 &= M_2, \end{aligned}$$

$$A_1 = \frac{N_1 - N_2 + M_2}{2} \quad A_2 = \frac{N_2 - N_3 + M_1}{2} \quad A_3 = \frac{-N_3 - N_4 + M_2}{2} \quad A_4 = \frac{N_1 + N_4 + M_1}{2}$$

$$\text{Constraints } \sum_{k=1}^4 (R_k + U_k) = -4i, \quad \sum_{k=1}^4 (S_k + T_k) = 0.$$

## Elliptic hypergeometric equation

V.S., 2004

$$A(x)(f(qx) - f(x)) + A(x^{-1})(f(q^{-1}x) - f(x)) + \lambda f(x) = 0,$$

$$A(x) = \frac{\prod_{k=1}^8 \theta(\varepsilon_k x; p)}{\theta(x^2, qx^2; p)}, \quad \lambda = \prod_{k=1}^6 \theta\left(\frac{\varepsilon_k \varepsilon_8}{q}; p\right), \quad \prod_{k=1}^8 \varepsilon_k = p^2 q^2, \quad \varepsilon_7 = \frac{\varepsilon_8}{q},$$

$$f(x) \propto \frac{V(q/c\varepsilon_1, \dots, q/c\varepsilon_5, cx, c/x, c/\varepsilon_8; p, q)}{\Gamma(c^2 x^{\pm 1}/\varepsilon_8, x^{\pm 1}\varepsilon_8; p, q)}, \quad c = \frac{\sqrt{\varepsilon_6 \varepsilon_8}}{p^2}.$$

Other solutions: multiply  $\varepsilon_1, \dots, \varepsilon_5, x$  by  $p^\#$  or permute  $\varepsilon_1, \dots, \varepsilon_5$  with  $\varepsilon_6$ , etc.

Degeneration to the complex  $6j$ -symbols level

$$B(s_2, N_2; s_3, N_3)(U(s_2 - i, N_2 - 1) - U) + B(s_3, N_3; s_2, N_2)(U(s_2 + i, N_2 + 1) - U) + U = 0,$$

$$B = \frac{(s_2 - s_4 + i(2 + N_2 - N_4))(s_4 - s_2 + i(N_4 - N_2))}{(s_3 - s_2 + i(2 + N_3 - N_2))(s_2 - s_3 + i(N_2 - N_3))} \prod_{k=1}^4 \frac{s_3 + t_k + i(2 + N_3 + M_k)}{s_4 + t_k + i(N_4 + M_k)}$$

$$U = \frac{\sum_{N \in \mathbb{Z}} \int_{-\infty}^{\infty} \prod_{a=1}^4 \Gamma(s_a - y, N_a - N) \Gamma(t_a + y, M_a + N) dy}{\Gamma(s_2 - s_4, N_2 - N_4) \Gamma(s_3 - s_4, N_3 - N_4)},$$

$$\sum_{a=1}^4 (N_a + M_a) = 0, \quad \sum_{a=1}^4 (s_a + t_a) = -4i,$$

$$U(s_2 - i, N_2 - 1) := U(s_2 - i, N_2 - 1; s_3 + i, N_3 + 1).$$

$\Rightarrow$   $q$ -difference equation becomes mixed recurrence + difference equation

Derkachov, Sarkissian, V.S., in preparation.

New rational degeneration of the hyperbolic Ruijsenaars  $N$ -body problem.

## Rational hypergeometric identities

Sarkissian, V.S., 2020

$b = \sqrt{\frac{\omega_1}{\omega_2}} = 1 + i\delta$ ,  $\delta \rightarrow 0^+$  limit of the Faddeev modular dilogarithm ( $c_{CFT} = 25$ )

$$\gamma^{(2)}\left(\sqrt{\omega_1\omega_2}(n+1+y\delta); \omega\right) \underset{\delta \rightarrow 0^+}{=} (4\pi\delta)^n e^{-\frac{\pi i}{2}n^2} \left(\frac{1-n-iy}{2}\right)_n,$$

where  $n \in \mathbb{Z}$ ,  $y \in \mathbb{C}$  and  $(a)_n$  is the standard Pochhammer symbol

$$(a)_n = \begin{cases} a(a+1)\cdots(a+n-1), & \text{for } n > 0, \\ \frac{1}{(a-1)(a-2)\cdots(a+n)}, & \text{for } n < 0. \end{cases}$$

Degeneration of the hyperbolic beta integral. For  $N_k \in \mathbb{Z}$ ,  $a_k \in \mathbb{C}$  and

$$\sum_{k=1}^6 N_k = 2, \quad \sum_{k=1}^6 a_k = 0,$$

$$\begin{aligned} & \frac{1}{8\pi i} \sum_{N \in \mathbb{Z} + \nu} \int_{C_N} (y^2 - N^2) \prod_{k=1}^6 \left(1 + \frac{a_k - N_k \pm (y - N)}{2}\right)_{N_k - 1 \pm N} dy \\ & = \prod_{1 \leq j < k \leq 6} \left(1 + \frac{a_j + a_k - N_j - N_k}{2}\right)_{N_j + N_k - 1}, \end{aligned}$$

contour  $C_N$  separates the poles emerging from  $\pm$  signs Pochhammer symbols.

New formula for rational hypergeometric identities.