

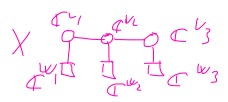
g-Opers, QQ-System and Betti Ansatz

(BPS/CFT correspondence) w/ D Sogo, A Zertlin (SP/n), g-opers
 conformal CIRH 03/13/19 F. Frenkel, D Sogo, A Zertlin (G, g)-opers

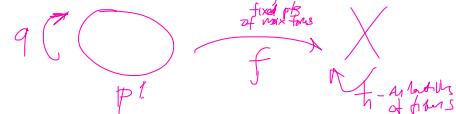
Motivation 1) BPS/CFT correspondence, dualities between integrable systems and gauge/string theory
 In particular, quantum/classical duality [DK, Gaiotto 2013]

XXZ/ERS or sd w=4 quiver gauge theory / Ad w=2 theory w RPS boundary conditions

2) Recent progress in enumerative AG (Okawa et al [AF0], [AO]) [PSZ], [KPS2]
 Quantum enumerative K-theory of toric/quiver varieties



In physics terms $V = \left(\begin{array}{l} \text{equiv. left push forward of} \\ \text{boundary space of} \\ \text{quasimaps from } \mathbb{P}^1 \text{ to } X \end{array} \right) = \sum_{\text{vert}} \mathbb{Z} d$
 k-th vertex factor of X \rightarrow $\sum_{\text{vert}} \mathbb{Z} d$ \rightarrow $\sum_{\text{vert}} \mathbb{Z} d$ \rightarrow $\sum_{\text{vert}} \mathbb{Z} d$

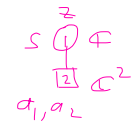


$V \sim e^{\frac{\widehat{W}(\vec{z})}{\log q}}$ $\xrightarrow{q \rightarrow 1}$ $K_*(X)$
 key - lang function for $\widehat{W}(\vec{z})$ \rightarrow $K_*(X)$
 $\widehat{W}(\vec{z}) = \sum_{i=1}^n \log z_i$

Quantum version of the duality

$X = T^* \mathbb{P}^1$ in type A

$H_{i(2)}^{\text{ERS}} V_{\vec{p}} = e_i(a_1, a_2) V_{\vec{p}}$
 $H_1 = \frac{1}{2} \frac{z_1 - z_2}{z_1 - z_2} p_1 + \frac{1}{2} \frac{z_2 - z_1}{z_2 - z_1} p_2$



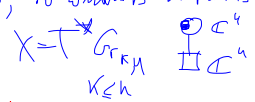
$X = T^* \mathbb{P}^1, z = \frac{z_1}{z_2}$
 $\hbar = p_1 = p_2$

\widehat{V}_1^*
 $p_1 f(z_1, z_2) = f(q z_1, z_2)$
 $p_2 f(z_1, z_2) = f(z_1, q z_2)$

Classically ($q \rightarrow 1$) $V_{\vec{p}} \rightarrow e^{\frac{\widehat{W}(\vec{z})}{\log q}} \Rightarrow \left\{ \frac{\partial \widehat{W}(\vec{z})}{\partial s} = 0 \right\} \Leftrightarrow \{ H_i = e_i \}$

Note that this only works for type-A (albeit one can use 'arbitrarily' direct to get some insight about types B and C)

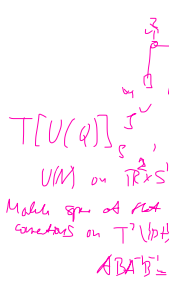
Also spinors exist for all possible weights, no constraints on ranks of bundles V_i, W_i
 \Rightarrow Need to extend the duality

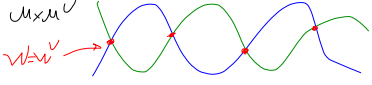
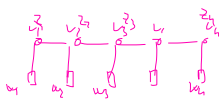


It turns out that the proper language to formulate such duality is geometric q-langlands correspondence
 3) Langlands. Categorical equivalence $D_x(\text{Bun } G) \cong D_{\frac{1}{\text{max}}}(G, q)$ in Riemann surface X

(Automorphic) Betti representations \Leftarrow $D_{\text{mod}}(\text{Bun } G) = \text{Loc } \mathcal{L}_G$ \Rightarrow Opers (Galois) B-side

Consider different version of the above relationship
 geometric q-langlands correspondence $\mathcal{A} \text{ (QQ-system } \mathcal{L}_g) \mid \mathcal{B} \text{ (G, g)-opers}$
 Categorical conjecture made by [Elliot-Lestoun]





Red dots
represent
in explicit form 5 nodes

Vectors of $S_d \rightarrow$ solution of XX^2
given by $L = L_{i, a}^S \wedge L_{j, b}^{S^L}$

$L_{L, q}^S$
 $L_{R, z}^{S^L}$

fixes a_j of $A + S$ and j in B
fixes a_j of $B + S$ and j in A

looks like generic
Langmuir + S duality

$$T[U(q)] = \sum_{i=1}^n u_i, N = \sum_{i=1}^n u_i$$

UM on $\mathbb{R}^2 \times S^1 \times \mathbb{R}^2$ S -Nambu pair B, C, L
generalized Nambu B, C, R

Mobius sym. of flat Riemannian
connections on $T^2 \times U(1)$

$$ABA^{-1}B^{-1} = C$$

function
associated with
of fundamental
over X

(X)

formalism
on X

(Galois)

te

$$\sum_{i=1}^n \left[\frac{B}{g_i(G, q)} \right]^{-1}$$

q-gaps & Quantum/Classical duality

Talks @ UC Davis
QMAP seminar 11/9/8

almost complete
A catch, to Singe
(2/2)

① Motivation of gaps • (quantum) (q-)geometric Langlands correspondence
Classical (rational) anal
 $K \rightarrow 0$

$D_k(\text{Bun}_G) \leftrightarrow D_{\text{inv}}(\text{Bun}_G)$
 $D_{\text{inv}}(\text{Bun}_G) \leftrightarrow \mathcal{L} \in \mathcal{L}_G$
In this file is the following about difference version of this

Bun_G on X compact Riemann surface

q-Geometric Langlands

- Recent progress in enumerative algebraic geometry and integrable systems (I guess of the latter)
- observed at my recent work w/ Zorin
- Quantum K-theory of symplectic resolutions, linkages between integrable systems which arise in gauge theory
- Ed Frenkel pointed out importance of gaps to langlands (07)
- Drinfeld & Sokolov (PS) studied them locally as gauge equivalence classes of certain left quantum bundles in 2d
- Baldwin, Brinkfield gave a constant-independent construction
- Timmer (07) Mura structure

Brief history

① $SL(2)$ gaps

1.1

Def A $G(\mathbb{R})$ -gap on P^1 is a triple (E, ∇, \mathcal{L}) , where E - rank-2 v-bundle on P^1
connection $\nabla: E \rightarrow E \otimes K$, \mathcal{L} - line subbundle such that residual map $\nabla|_{\mathcal{L}}: \mathcal{L} \rightarrow E/\mathcal{L} \otimes K$ is an iso
If the structure group can be reduced to $SL(2)$ we can get an $SL(2)$ gap.

P^1, \mathcal{L} is a line subbundle on E . In the vicinity of any point z we have

$$S(z) \wedge \nabla_z S(z) \neq 0, \quad \nabla_z = \frac{d}{dz} \nabla$$

Talk about $SL(2)$ -gaps w/ regular singularities and trivial monodromies around singular points



An $SL(2)$ -gap w/ sing at z_1, z_2 , z_3, z_4 of weights k_1, k_2, k_3, k_4 is the same as above, where ∇ is also a connection except at z_3, z_4 where it has order of vanishing k_3, k_4

$$S(z) \wedge \nabla_z S(z) \sim (z-z_3)^{k_3} \quad (1)$$

Trivial monodromy condition \Rightarrow after an appropriate gauge transformation A is trivial

in $P^1 \setminus \{z_i\}$ subbundle \mathcal{L} has section $S(z) = \begin{pmatrix} q_1(z) \\ q_2(z) \end{pmatrix}$, q_1, q_2 have no common roots

then (1) reads $q_2 \partial_z q_1 - \partial_z q_1 q_2 = \rho(z)$, $\rho(z) = \prod_{i=1}^4 (z-z_i)^{k_i}$ (2)

By a gauge transformation we can make $\deg(q_1) < \deg(q_2)$ and q_2 is monic (always possible)

Nondegenerate \Rightarrow none of z_i are roots of $q_2 \Rightarrow$ each root of q_2 has multiplicity 1

Let $k = \sum_{i=1}^4 k_i = \deg(\rho)$ $\Rightarrow \deg(q_1) \deg(q_2) = k+1 \Rightarrow \deg(q_1) = \ell \leq \frac{k}{2}$

Brankle (2) as $\partial_z \begin{pmatrix} q_1(z) \\ q_2(z) \end{pmatrix} = - \frac{\rho(z)}{q_2(z)^2}$

Compute residues at each z_i we get

$$\sum_{i=1}^4 \frac{1}{z_i - z_j} \frac{k_i}{z_i - z_j} = \sum_{j=1}^4 \frac{E}{z_j - z_i} \quad i=1, \dots, \ell$$

to study \Rightarrow do analytical transform $z \rightarrow \frac{1}{z}$

Better question for $SL(2)$ Geometric moduli of level $k-2\ell > 0$

Theorem 1 There is a 1-1 correspondence between spectra of the S_2 Ginzburg model and spec of non-degenerate S_2 quiver w/ trivial boundary and regular singularities

1.2 Mirna structures

below to condition of non-degenerate dual $K-2\mathbb{P}_-$.

Mirna oper $(E, \mathcal{D}, \mathcal{L})$ together with additional subbundle $\hat{\mathcal{L}}$ generated by ∇

\mathcal{L} span E everywhere except maybe finitely many points.

Thus for any oper we have a family of Mirna oper's parameterized by G/B flag (Frankel)

Pick $\hat{\mathcal{L}}$ being generated by $\mathbb{3} = (1)$

Theorem 2 There is a 1-1 relation between solutions of Bethe Ansatz equations of the Ginzburg model and S_2 Mirna oper's w/ reg sing & trivial monodromy with weights and their points z_1, z_2

1.3 Casimirs and differential operators

Eigenvalues of Casimir Hamiltonian study from oper

Any $SU(2)$ gauge transformation $g(z) = \begin{pmatrix} q & -q_+ \\ 0 & q_-^{-1} \end{pmatrix}$ so that $g(z)S(z) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$A = -(\partial_z g)g^{-1} = \begin{pmatrix} -\frac{\partial_z q}{q} & -\frac{q_+}{q} \\ 0 & \frac{\partial_z q_-}{q_-} \end{pmatrix} \xrightarrow[\text{gauge transform}]{\text{diagonal}} A(z) = \begin{pmatrix} -\frac{q_+}{q} & -1 \\ 0 & u(z) \end{pmatrix}$$

($g^{-1} \begin{pmatrix} 1 & 0 \\ 0 & g_-^{-1} \end{pmatrix}$)

$$\text{where } u(z) = -\frac{\partial_z q}{q} + \frac{\partial_z q_-}{q_-} = -\sum_m \frac{k_m}{z - z_m} + \sum_i \frac{1}{z - w_i}$$

Apply Mirna basis from $g(z) = \begin{pmatrix} u & 1 \\ 1 & 0 \end{pmatrix}$ so the two conditions read $B(z) = \begin{pmatrix} 0 & -1 \\ -t(z) & 0 \end{pmatrix}$.

$$\text{where } t(z) = z_2 u(z) + u^2(z) = \sum_m \frac{u_m (k_m + 1) A}{(z - z_m)^2} + \sum_i \frac{C_i}{z - w_i}$$

↑ eigenvalues of Casimir Hamiltonians

$$C_{w_i} = k_{w_i} \left(\sum_m \frac{k_m}{z_m - w_i} - \sum_{j=1}^L \frac{1}{z_j - w_i} \right)$$

Horizontal section $f = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$ to $B(z)$ is determined by $(z_2 - B(z)) \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} = 0$

$$(\partial_z^2 - t(z))f_1(z) = 0 \quad \leftarrow \text{projector condition}$$

1.4 Irregular singularities [FFTL, FFR]

$A \sim d + \begin{pmatrix} a & \\ & a \end{pmatrix} dz$ at $P^1 \setminus \infty$ section $S(z) = e^{-\int a dz} \begin{pmatrix} q_+ \\ q_- \end{pmatrix}$

then Bethe equations will take the following form $-2a + \sum_{j=1}^L \frac{k_j}{z_j - w_j} = \sum_{j=1}^L \frac{c_j}{w_j - w_j}$

Theorem 3 There is a 1-1 correspondence between solutions of inhomogeneous Bethe equations and non-degenerate S_2 quiver w/ irregular singularities and z_1, z_2 and W -double pole w/ 2-regular $-\begin{pmatrix} a & \\ & a \end{pmatrix}$

① Q-ops

②

Let $q \in \mathbb{C}^X$, consider E^T - pull back of E - vector bundle over \mathbb{P}^1 with map $z \mapsto qz$
 Given map of vector bundles $A: E \rightarrow E^T$ gives linear map of given local $E_2 \rightarrow E_2$

q -quasilinearities $A(z) \mapsto g(qz)A(z)g^{-1}(z)$
 change of transition by $g(z)$
 Let $D_g: E \rightarrow E^q$ then $D_g(z) = AS$
 $S(z) \mapsto S(qz)$

Def A monodromic $(\alpha(z), q)$ -connection over \mathbb{P}^1 is (E, A) , where E is a fractional vector bundle of rank N over \mathbb{P}^1 and A is a meromorphic section of the sheaf $\text{Hom}_{\mathcal{O}_{\mathbb{P}^1}}(E, E^T) \otimes \mathcal{O}_{\mathbb{P}^1}(q)$ for which $A(z)$ is invertible. If there exists a trivialization where $\det A(z) = 1$ then we have an $(SL(N, \mathbb{C}))_q$ monodromy of (E)

Def $(\alpha(z), q)$ -oper is a triple (E, A, \mathcal{L}) , where (E, A) is (q) -monodromic
 \mathcal{L} is a line subbundle such that the induced map $\mathcal{L} \rightarrow (E/\mathcal{L})^T$ is an iso

- $(SL(N, q))$ -oper w/ regular singularities $z_1, z_2, z_3 \neq 0$, can w/ consists k_1, k_2, k_3 is a monodromic $(SL(N, q))$ -oper for which \mathcal{L} is an iso everywhere on $\mathbb{P}^1 \setminus \{z_1, z_2, z_3\}$ except the points $z_m, q^1 z_m, q^2 z_m, \dots, q^{k_m-1} z_m, m=1, 2, 3$ S(qz)A(z)S(z) \neq 0
- An $(SL(N, q))$ -oper (E, A, \mathcal{L}) w/ reg singularities is called a Z -twisted q -oper if A is gauge equivalent to Z^{-1} , where $Z = \begin{pmatrix} z & 0 \\ 0 & 1 \end{pmatrix}$
- Monodromic q -oper (E, A, \mathcal{L}, Z) , where Z is preserved by A



2.2 Q-connection and Beilinson Ansatz

Theorem 4 There is a 1-1 correspondence between the set of non-degenerate solutions of $SL(N, \mathbb{C})^X$ Beilinson equation and the set of non-degenerate Z -twisted $(SL(N, q))$ -opers w/ regular sing at $z_1, z_2, z_3 \neq 0$ w/ weights k_1, k_2 and can be represented by monodromic q -connection

$$\left(D_q^T - T(qz)D_q - \frac{S(qz)}{S(z)} \right) f_1 = 0$$

Proof $S(qz)A(z)S(z) = f(z)$
 $\int^1 Q_1(z)Q_1(qz) - \int^1 Q_2(qz)Q_2(z) = \prod_{m=1}^k \prod_{j=0}^{k_m-1} (z - q^j z_m)$

Can always write $Q_i(z) = \prod (z - w_i)$
 assume that q -conjugates of z_m, w_i do not overlap

Then $\prod_{m=1}^k \frac{w_i - q^{k_m} z_m}{w_i - q z_m} = \sum_{i=1}^N q^{-k_i} \prod_{j=1}^k \frac{w_j - w_i}{w_i - q w_j}, i=1, \dots, N$

Analogously get X^2 Beilinson's (transfer matrix)

$$A(z) = \begin{pmatrix} a(z) & f(z) \\ 0 & a(z) \end{pmatrix}, \quad a(z) = \sum^1 Q_1(qz)Q_1^{-1}(z) \Rightarrow \hat{A}(z) = \begin{pmatrix} 0 & f(z) \\ -f^{-1}(z) & T(qz)P(qz) \end{pmatrix}$$

$$T(z) = \sum^1 P(q^{-1}z) \frac{Q_1(qz)}{Q_1(z)} + \sum^2 P(z) \frac{Q_2(q^{-1}z)}{Q_2(z)} - q$$

- q -signature of the X^2 transfer matrix

23 trigonometric Riesz means - Steiner Model w/ 2 degrees of freedom

Take $q_+(z) = 1$, $p(z) = (z - z_+) / (z - z_-)$

$$Q_- = z - p_-, \quad Q_+ = c(z - p_+), \quad c = \bar{q}^{-1}(\beta - \gamma^{-1})$$

From previous iteration relation

$$z^2 - \frac{z}{q} \left[\frac{\beta - \gamma}{\gamma - \beta} p_+ + \frac{\gamma - \beta}{\gamma - \beta} p_- \right] + \frac{p_+ p_-}{q} = (z - z_+)(z - z_-)$$

from Whor $H_1 = q(z_+ + z_-)$

fRS $H_2 = z_+ z_-$

Newton's

$$q \text{ iter } \Leftrightarrow \det(zI - L_{\text{fRS}}) = p(z)$$

24 SL(N) Generalization

25 QRS theory

26 Eigenvalue theory