

Non-Abelian Vortices and 4d/2d Correspondence

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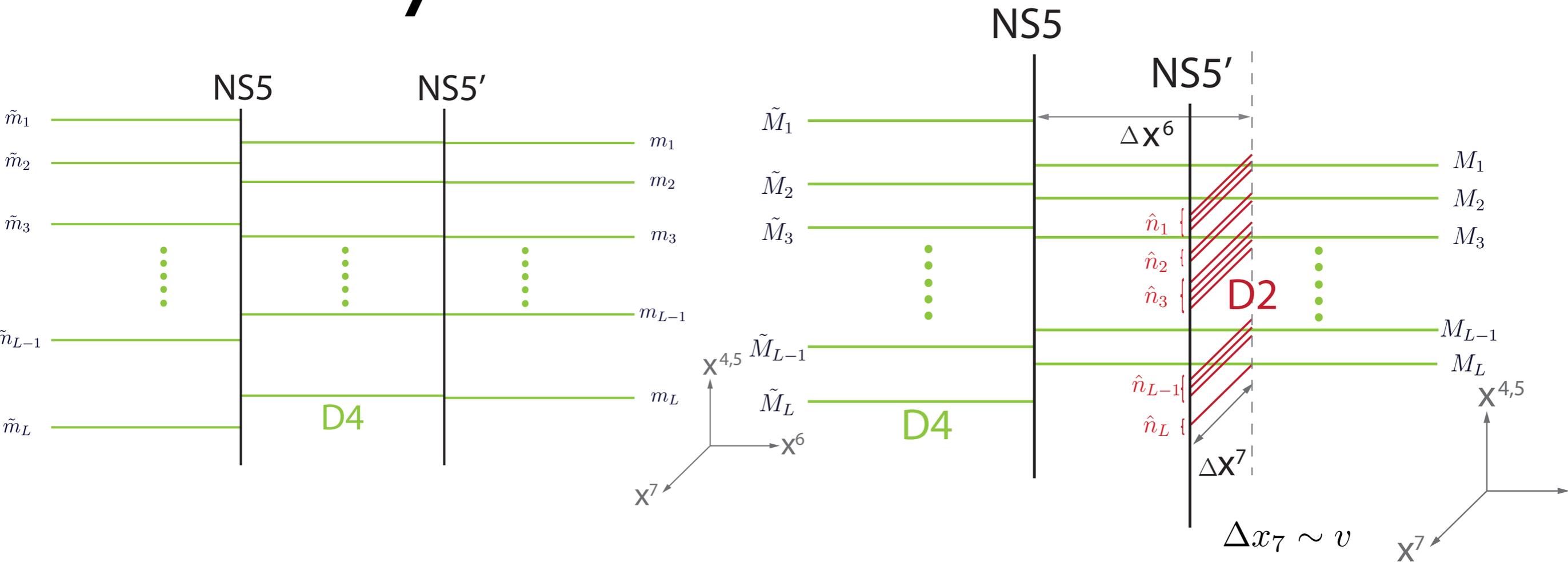


work in progress with A.Gorsky, H. Chen

Outline

- 4d/2d correspondence in a nutshell
 - ★ *From brane construction*
 - ★ *The Dictionary of 4d/2d*
 - ★ *Derivation from vortices*
- 4d/2d in Omega background
 - ★ *N=2 theory in Nekrasov-Shatashvili limit*
 - ★ *Monopoles vortices and strings in Omega background*

Hanany-Witten construction



SQCD $N_f = 2N_c$

2d FI parameter $r = \frac{\Delta x^6}{2\pi g_s l_s} = \frac{4\pi}{e^2}$

Higgs branch root

$$\sigma = X^4 + iX^5, \quad Z = X^1 + iX^2$$

Color-flavor locked phase of SQCD

$$V = \frac{1}{g^2} \text{Tr} |[\sigma, \sigma^\dagger]|^2 + \text{Tr} |[\sigma, Z]|^2 + \text{Tr} |[\sigma, Z^\dagger]|^2 + \sum_{a=1}^N \psi_a^\dagger \sigma^\dagger \sigma \psi_a + \frac{g^2}{2} \text{Tr} \left(\sum_a \psi_a \psi_a^\dagger + [Z, Z^\dagger] - r \mathbf{1}_k \right)^2$$

In the simplest case $\mathbb{C}\mathbb{P}^{N-1}$ model

4d / 2d duality

[Dorey Hollowood, Tong]

$\mathcal{N} = 2$ $SU(N)$ SQCD

$N_f = N + \tilde{N}$ fund hypers

w/ masses

m_1, \dots, m_N $\mu_1, \dots, \mu_{\tilde{N}}$

$$\tau = \frac{4\pi i}{g^2} + \frac{\theta}{2\pi}$$

on baryonic Higgs branch

(2,2) $U(1)$ GLSM e

N chiral +1 \tilde{N} chiral -1

w/ *twisted* masses

m_1, \dots, m_N $\mu_1, \dots, \mu_{\tilde{N}}$

$$\tau = ir + \frac{\theta}{2\pi}$$

vortex moduli space

BPS dyons
(Seiberg-Witten)

kinks interpolating
between different vacua

BPS spectra (as functions of masses, Λ) are the same

4d / 2d duality

[Dorey Hollowood, Tong]

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BPS spectra (as functions of masses, Λ) are the same
Nonabelian vortices help to understand it from
pure field theory constructions

$U(N_c)$ $\mathcal{N} = 2$ $d = 4$ SQCD w/ N_f quarks

$$\{Q_\alpha^I, \bar{Q}_{\dot{\beta}}^J\} = 2\delta^{IJ} P_{\alpha\dot{\beta}} + 2\delta^{IJ} Z_{\alpha\dot{\beta}}$$

$$\{Q_\alpha^I, Q_\beta^J\} = 2Z_{\alpha\beta}^{IJ}$$

$$\mathcal{L} = \text{Im} \left[\tau \int d^4\theta \text{Tr} \left(Q^{i\dagger} e^V Q_i + \tilde{Q}^{i\dagger} e^V \tilde{Q}_i + \Phi^\dagger e^V \Phi \right) \right] \\ + \text{Im} \left[\tau \int d^2\theta \left(\text{Tr} W^{\alpha 2} + m_j^i \tilde{Q}_i Q^j + Q_i \Phi \tilde{Q}^i \right) \right]$$

bosonic part

FI term

$$S = \int d^4x \text{Tr} \left\{ \frac{1}{2g^2} F_{\mu\nu}^2 + \frac{1}{g^2} |D_\mu \Phi|^2 + |\nabla_\mu Q|^2 + \frac{g^2}{4} (Q\bar{Q} - \xi)^2 + |\Phi Q + QM|^2 \right\}$$

BPS conditions

$$B_3 - g^2 (Q\bar{Q} - \xi^2) = 0 \\ \nabla_3 Q = 0$$

String tension

$$T = \xi \int d^2x \text{Tr} F_{12} = 2\pi\xi n$$

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strings

monopoles domain walls

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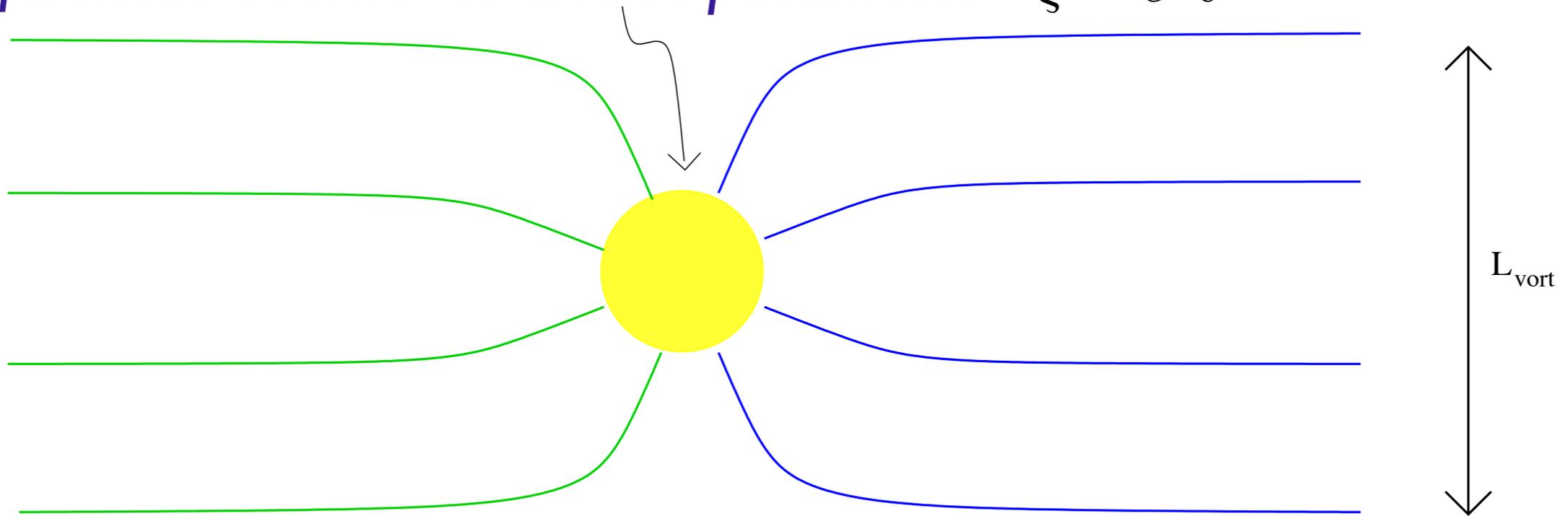
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Monopoles on Higgs Phase [Shifman, Yung] [Tong]

Higgs branch condition $\phi = \text{diag}(m_i)$, $q_i^a = v\delta_i^a$, $\tilde{q}_i^a = 0$

Pattern of symmetry breaking depends on the relationship between the differences of masses and FI parameter $\xi = e^2 v^2$



$$ev \gg \Delta m$$

$$\longleftrightarrow L_{\text{mon}}$$

$$U(N)_G \times SU(N)_F \xrightarrow{v} SU(N)_{\text{diag}} \xrightarrow{m} U(1)_{\text{diag}}^{N-1}$$

$$ev \ll \Delta m$$

$$U(N)_G \times SU(N)_F \xrightarrow{m} U(1)_G^N \times U(1)_F^{N-1} \xrightarrow{v} U(1)_{\text{diag}}^{N-1}$$

(2,2) 2d GLSM

[Witten]

Consider U(1) gauge theory

$$\mathcal{L}_{\text{vortex}} = \frac{1}{2g^2} (F_{01}^2 + |\partial\sigma|^2) + \sum_{i=1}^{N_c} (|\mathcal{D}\psi_i|^2 + |\sigma - m_i|^2 |\psi_i|^2) + \frac{g^2}{2} \left(\sum_{i=1}^{N_c} |\psi_i|^2 - r \right)^2$$

Vacuum i : $\sigma = m_i$, $|\psi_j|^2 = r\delta_{ij}$

for vortex embedded into
i's U(1) subgroup

FI term runs $r(\mu) = r_0 - \frac{N_c}{2\pi} \log \left(\frac{M_{UV}}{\mu} \right) \Rightarrow \Lambda = \mu \exp \left(-\frac{2\pi r(\mu)}{N_c} \right)$

Effective twisted superpotential

$$\mathcal{W}(\Sigma) = \frac{i}{2} \tau \Sigma - \frac{1}{4\pi} \sum_{i=1}^{N_c} (\Sigma - m_i) \log \left(\frac{2}{\mu} (\Sigma - m_i) \right) \Rightarrow \text{Vacua } \exp \frac{\partial \widetilde{\mathcal{W}}}{\partial \sigma} = 1$$

Central charge

$$Z = -i \sum_{i=1}^{N_c} (m_i S_i + m_{D i} T_i)$$

$$m_{D i} = -2i\mathcal{W}(e_i) = \frac{1}{2\pi i} N_c e_i + \frac{1}{2\pi i} \sum_{j=1}^{N_c} m_j \log \left(\frac{e_i - m_j}{\Lambda} \right)$$

'ANO' String

$U(N)$ gauge theory with fundamental matter $q \rightarrow UqV$ $U \in U(N)_G, V \in SU(N)_F$

$$N_f = N_c$$

$$S = \int d^4x \operatorname{Tr} \left(\frac{1}{2e^2} F^{\mu\nu} F_{\mu\nu} + \frac{1}{e^2} (\mathcal{D}_\mu \phi)^2 \right) + \sum_{i=1}^{N_f} |\mathcal{D}_\mu q_i|^2$$

$$- \sum_{i=1}^{N_f} q_i^\dagger \phi^2 q_i - \frac{e^2}{4} \operatorname{Tr} \left(\sum_{i=1}^{N_f} q_i q_i^\dagger - v^2 1_N \right)^2$$

Vacuum

$$\phi = 0, \quad q_i^a = v \delta_i^a$$

breaks symmetry

(color-flavor locking)

$$U(N)_G \times SU(N)_F \rightarrow SU(N)_{\text{diag}}$$

**Induces nontrivial topology
on moduli space**

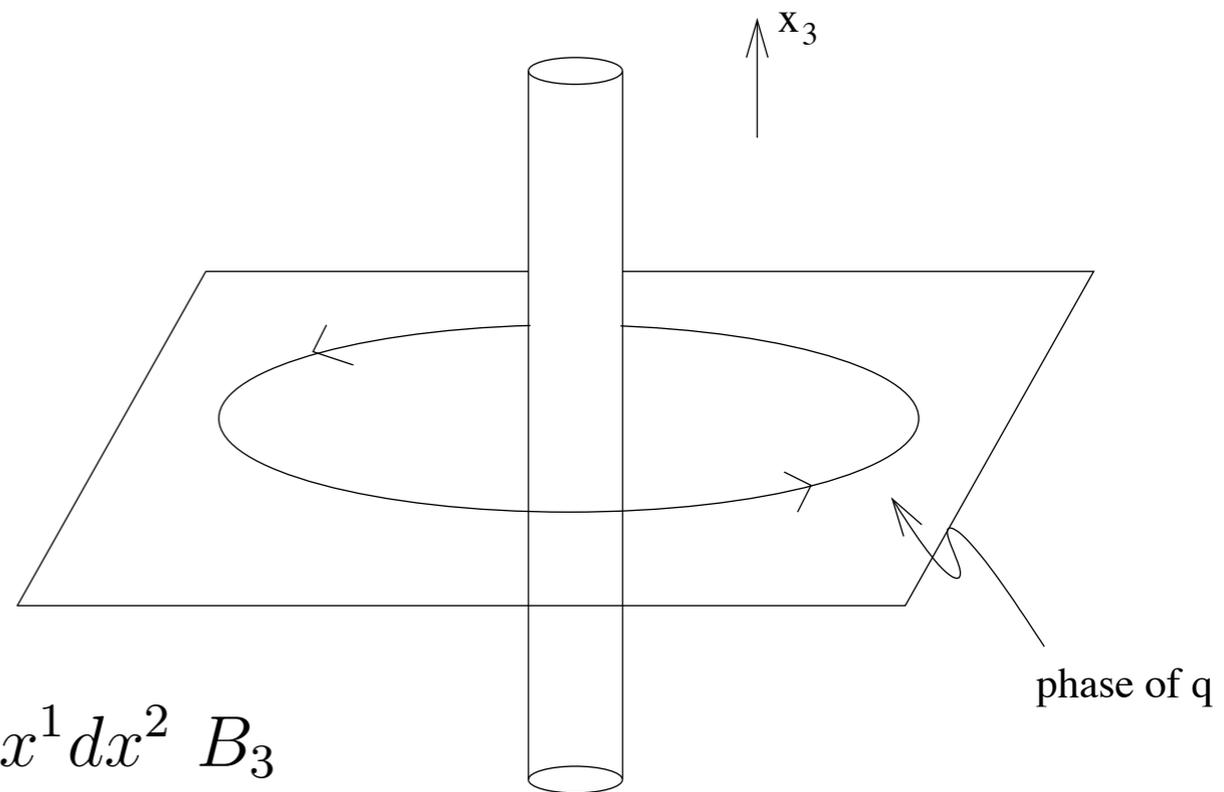
$$\Pi_1 (U(N) \times SU(N) / SU(N)_{\text{diag}}) \cong \mathbf{Z}$$

**To find a string need
winding at infinity**

$$q_N \sim q e^{ik\theta}$$

$$A_\theta \sim \frac{k}{\rho}$$

$$2\pi k = \operatorname{Tr} \oint_{S_\infty^1} i \partial_\theta q q^{-1} = \operatorname{Tr} \oint_{S_\infty^1} A_\theta = \operatorname{Tr} \int dx^1 dx^2 B_3$$

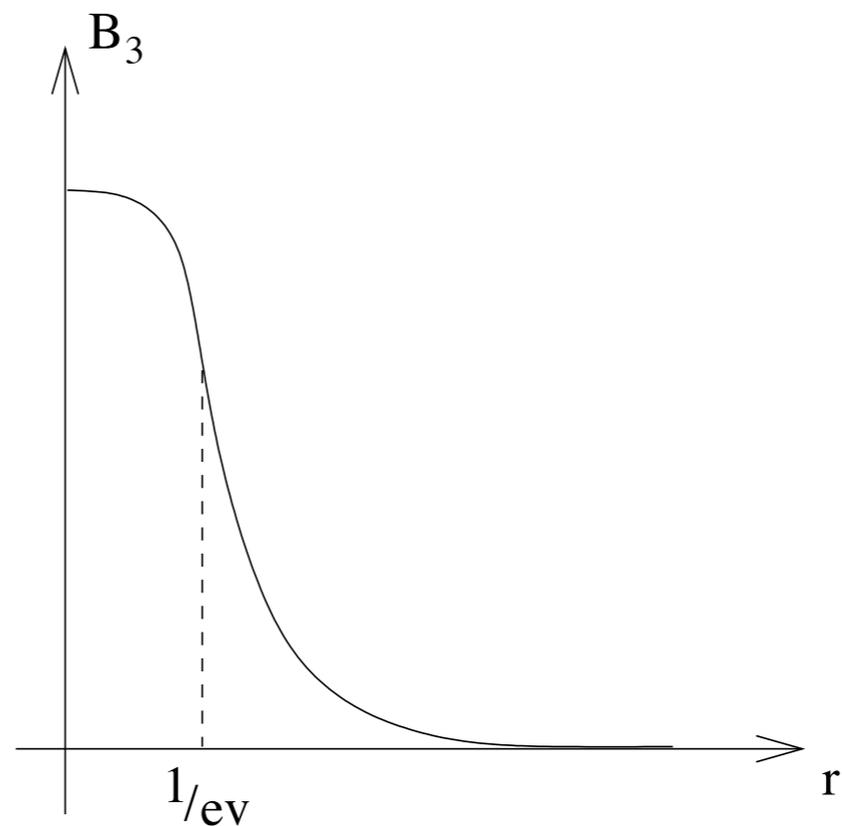
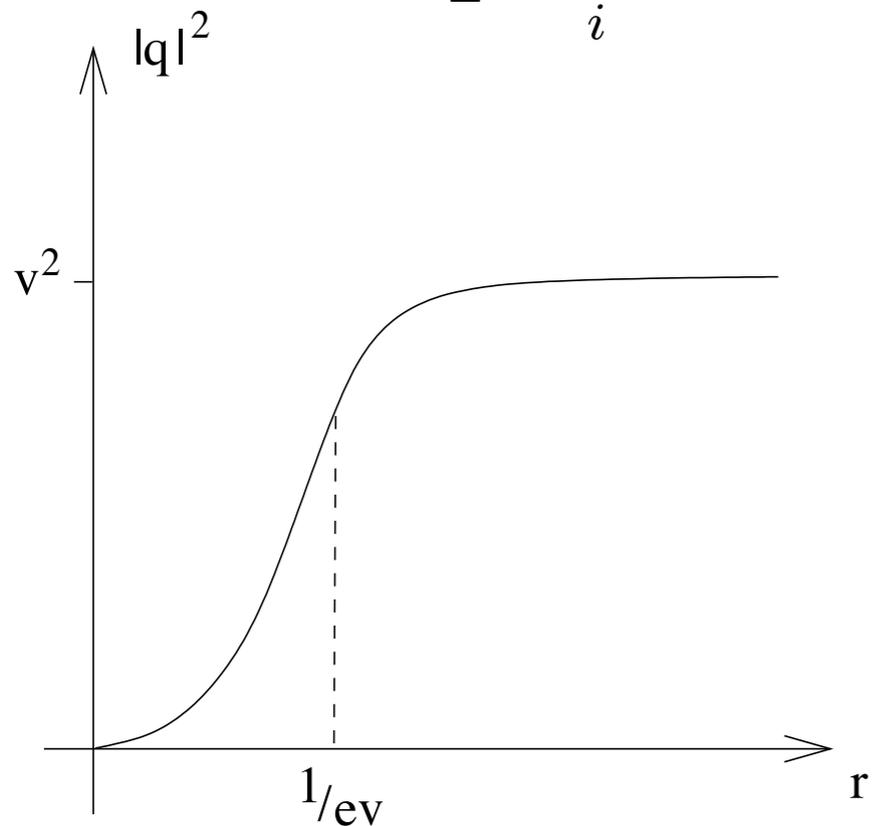


BPS equations for vortex

$$\begin{aligned}
 T_{\text{vortex}} &= \int dx^1 dx^2 \text{Tr} \left(\frac{1}{e^2} B_3^2 + \frac{e^2}{4} \left(\sum_{i=1}^N q_i q_i^\dagger - v^2 1_N \right)^2 \right) + \sum_{i=1}^N |\mathcal{D}_1 q_i|^2 + |\mathcal{D}_2 q_i|^2 \\
 &= \int dx^1 dx^2 \frac{1}{e^2} \text{Tr} \left(B_3 \mp \frac{e^2}{2} \left(\sum_{i=1}^N q_i q_i^\dagger - v^2 1_N \right) \right)^2 + \sum_{i=1}^N |\mathcal{D}_1 q_i \mp i \mathcal{D}_2 q_i|^2 \\
 &\mp v^2 \int dx^1 dx^2 \text{Tr} B_3 \quad \geq \mp v^2 \int d^2 x \text{Tr} B_3 = 2\pi v^2 |k| \quad (
 \end{aligned}$$

gives

$$B_3 = \frac{e^2}{2} \left(\sum_i q_i q_i^\dagger - v^2 1_N \right) \quad (\mathcal{D}_x - i \mathcal{D}_y) q_i = 0$$



Vortices

Simple vortex w/ $N=1$, $k=1$ (ANO) has two collective coordinates-translations in x,y directions

U(N) vortex
has more moduli

$$A_z = \begin{pmatrix} A_z^* & & & \\ & 0 & & \\ & & \ddots & \\ & & & 0 \end{pmatrix}, \quad q = \begin{pmatrix} q^* & & & \\ & v & & \\ & & \ddots & \\ & & & v \end{pmatrix}$$

Moduli space
($k=1$)

$$SU(N)_{\text{diag}} / S[U(N-1) \times U(1)] \cong \mathbb{C}P^{N-1}$$

For higher k

$$\dim(\mathcal{V}_{k,N}) = 2kN \quad \mathcal{V}_{1,N} \cong \mathbf{C} \times \mathbb{C}P^{N-1}$$

Non-Abelian String

[Auzzi, Bolognesi, Evslin, Konishi, Yung]

[Shifman Yung]

$$\varphi = U \begin{pmatrix} \phi_2(r) & 0 & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & \dots & \phi_2(r) & 0 \\ 0 & 0 & \dots & \phi_1(r) \end{pmatrix} U^{-1},$$

Take Abelian string solution
Make global rotation

$$A_i^{\text{SU}(N)} = \frac{1}{N} U \begin{pmatrix} 1 & \dots & 0 & 0 \\ \dots & \dots & \dots & \dots \\ 0 & \dots & 1 & 0 \\ 0 & 0 & \dots & -(N-1) \end{pmatrix} U^{-1} (\partial_i \alpha) f_{NA}(r)$$

Matrix U parameterizes
orientational modes

$$A_i^{\text{U}(1)} = -\frac{1}{N} (\partial_i \alpha) f(r), \quad A_0^{\text{U}(1)} = A_0^{\text{SU}(N)} = 0,$$

Gauge group is broken to \mathbb{Z}_N

$$\frac{SU(N)}{SU(N-1) \times U(1)} = \mathbb{CP}^{N-1}$$

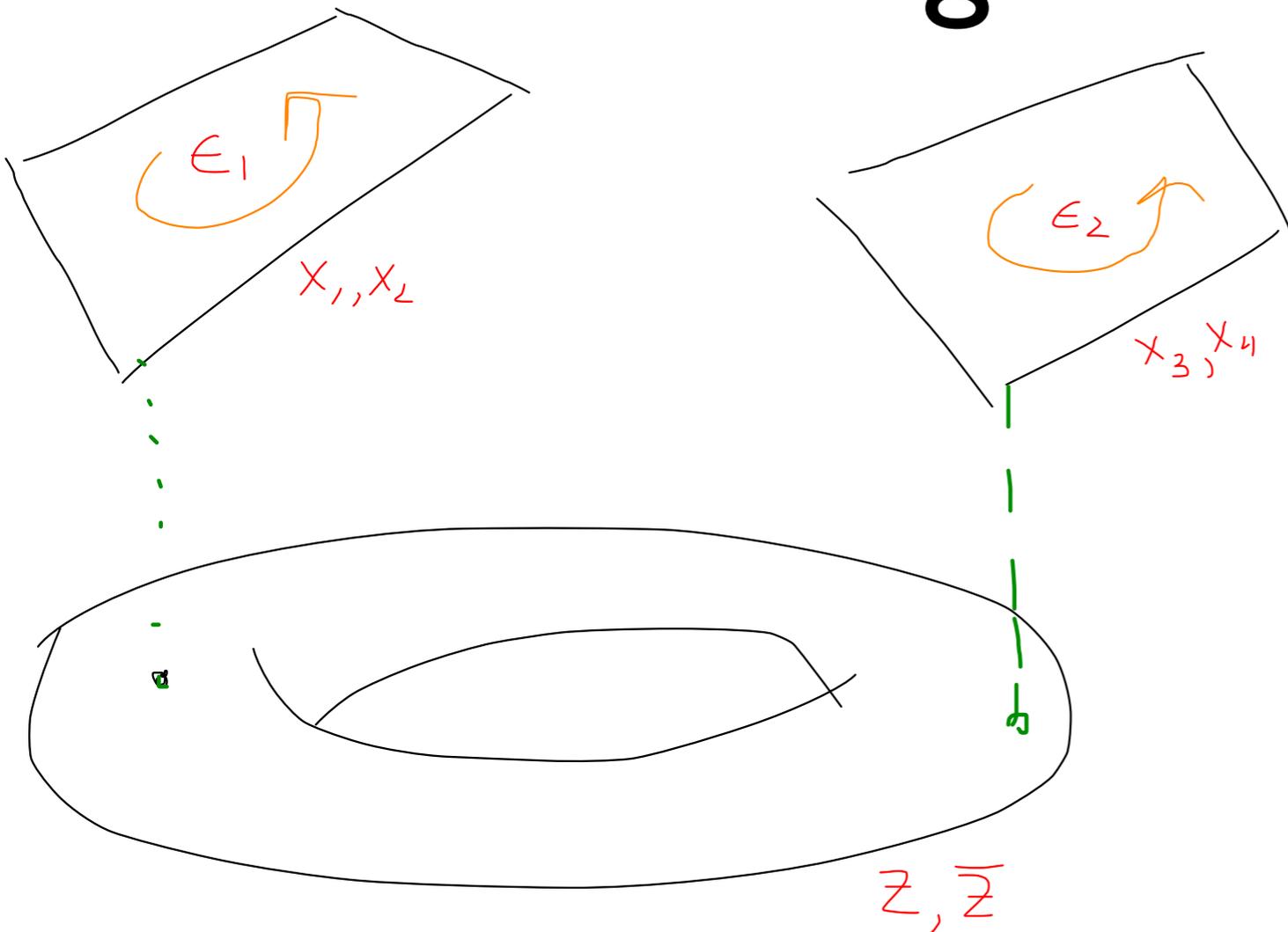
All bulk degrees of freedom massive

$$M^2 = e^2 v^2$$

Theory is fully Higgsed

Omega background

[Nekrasov et al]



Rotational symmetry
broken to maximal torus

$$SO(4) \rightarrow SO(2) \times SO(2)$$

6d Metric

$$G_{AB}dx^A dx^B = Adz d\bar{z} + (dx^m + \Omega^m dz + \bar{\Omega}^m d\bar{z})^2$$

We will be interested in Nekrasov-Shatashvili limit

$$\Omega^m = (-i\epsilon x^2, i\epsilon x^1, 0, 0)$$

$$\epsilon_2 \rightarrow 0$$

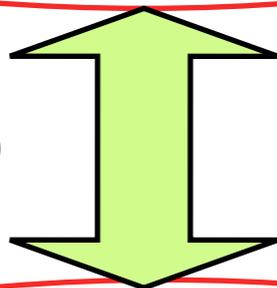
4d/2d in Omega background

[Dorey
Hollowood Lee]

N=2 SQCD in Omega background
in NS limit with $N_f=2N_c$

$$\vec{a} = \vec{m}_F - \vec{n}\epsilon \quad \vec{n} = (n_1, \dots, n_L) \in \mathbb{Z}^L$$

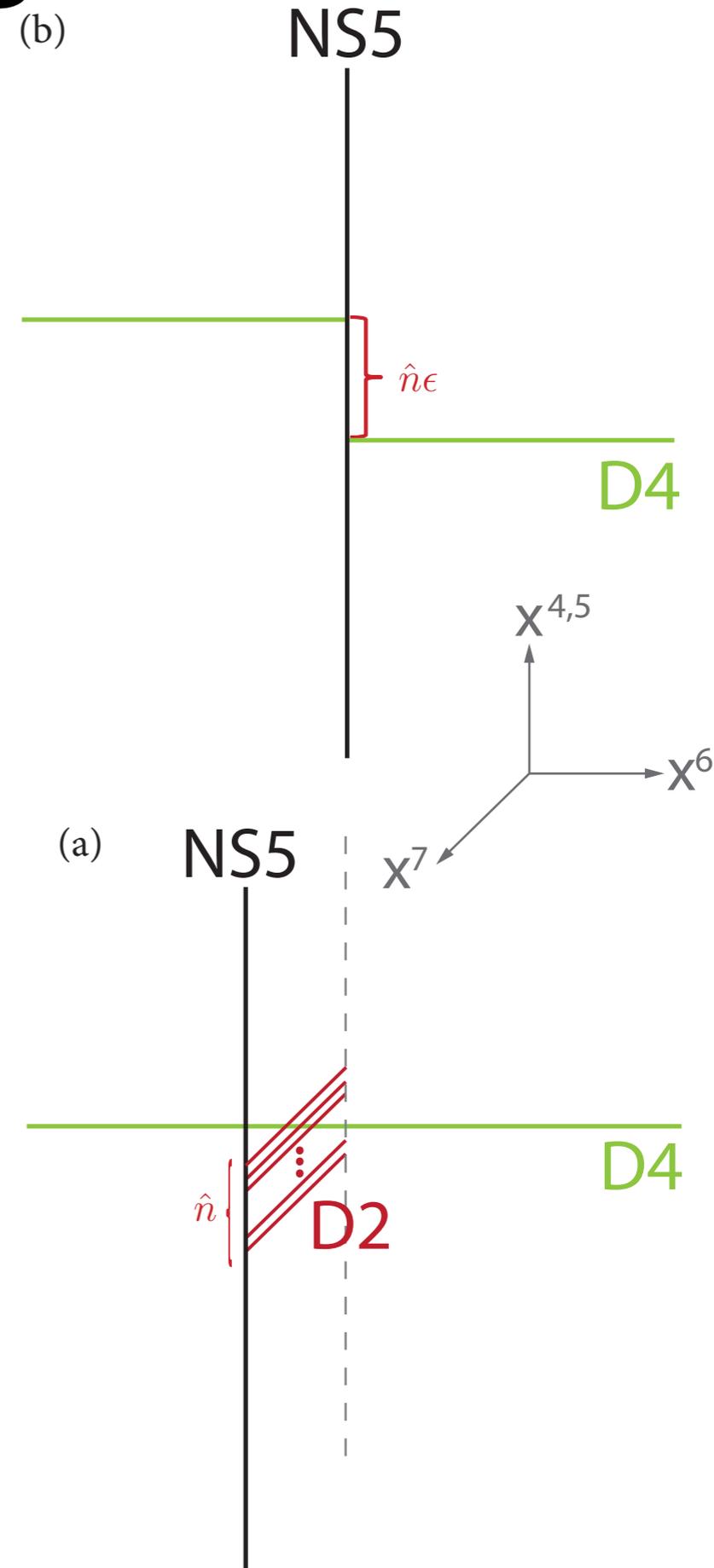
$$\mathcal{W}^{(I)} \stackrel{\text{on-shell}}{\equiv} \mathcal{W}^{(II)}$$



(2,2) GLSM w/ gauge group $U(K)$
massive adjoint and twisted masses

$$\vec{M}_F = \vec{m}_F - \frac{3}{2}\vec{\epsilon}, \quad \vec{M}_{AF} = \vec{m}_{AF} + \frac{1}{2}\vec{\epsilon}.$$

$$M_{adj} = \epsilon \quad K = \sum_{i=1}^N n_i - N$$



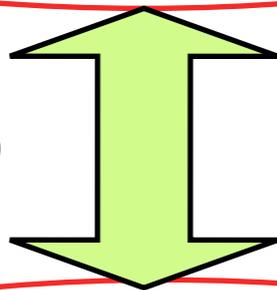
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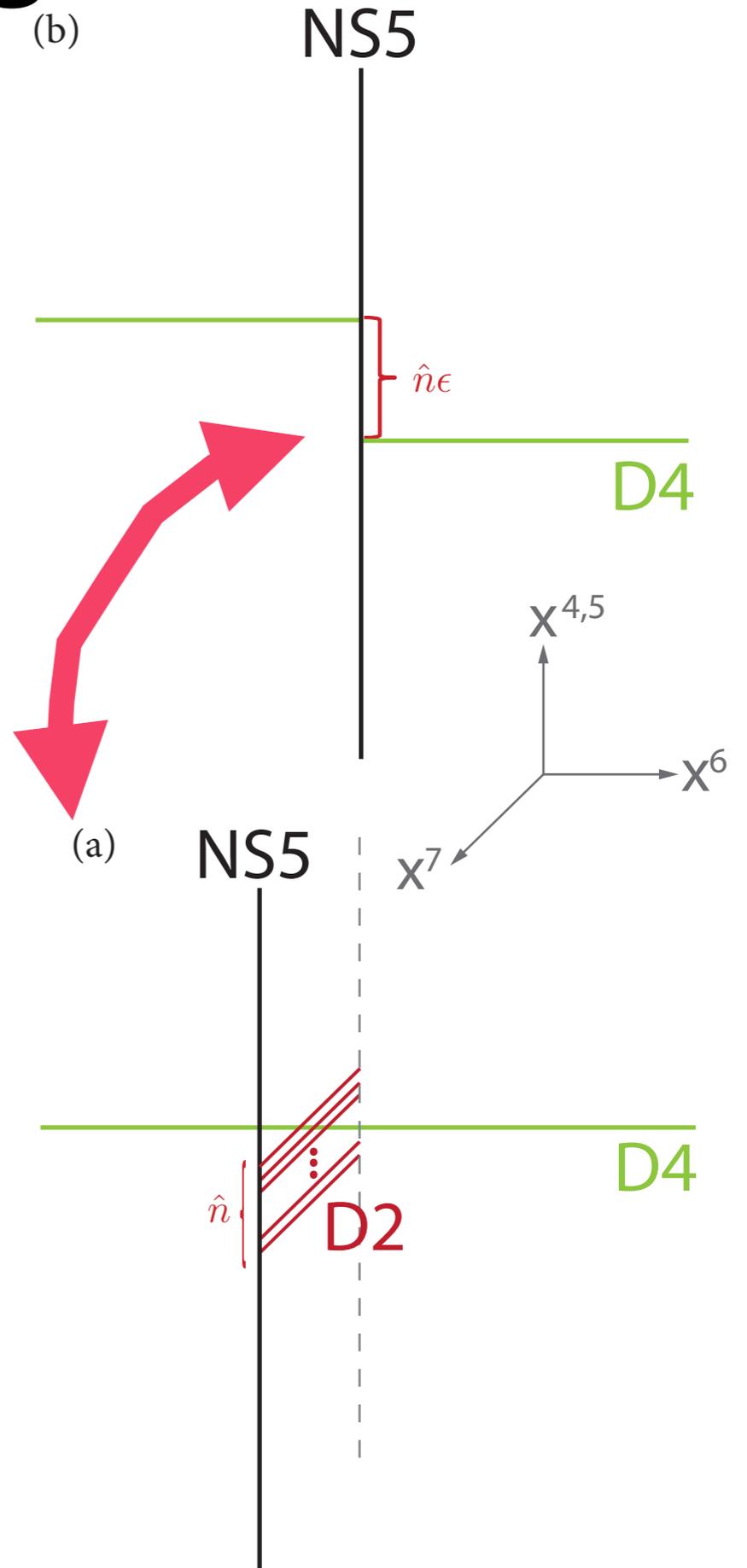
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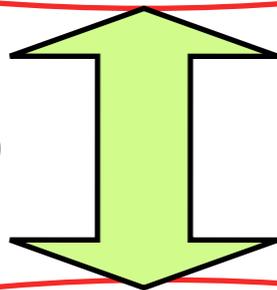
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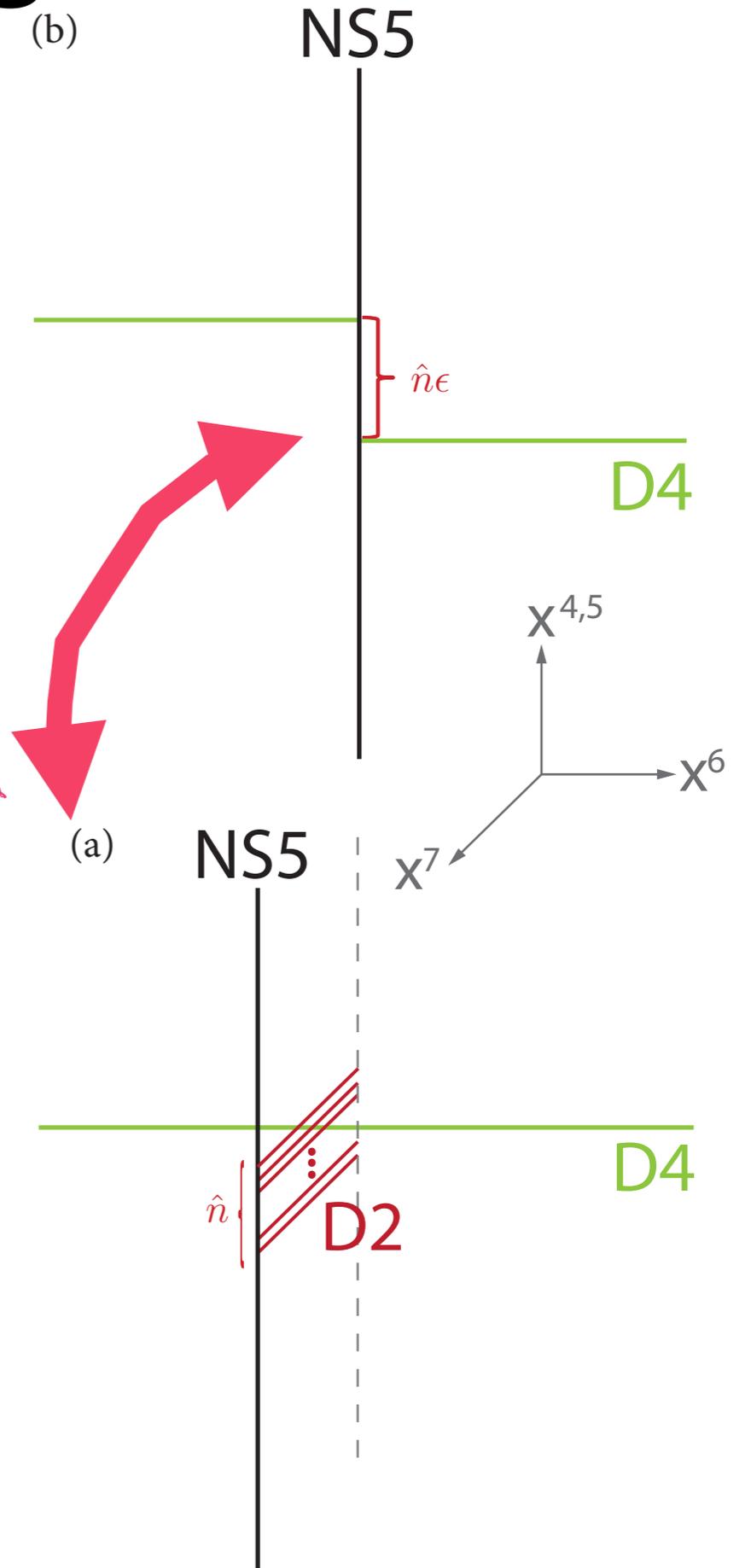


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conifold
transition



Vortices in Omega background [PK Gorsky Chen]

in progress

Symmetry breaking pattern

$$SU(2)_c \times SU(2)_R \times SU(2)_{\mathcal{R}} \rightarrow U(1)_c \times SU(2)_{R+\mathcal{R}}$$

SUSY transform pure SYM

$$\begin{aligned} \delta\Lambda_\alpha^I &= \zeta_\beta^I ((\sigma^{mn})_\alpha^\beta F_{mn} + i[\phi, \bar{\phi}] \delta_\alpha^\beta + \nabla_m (\bar{\Omega}^m \phi - \Omega^m \bar{\phi}) \delta_\alpha^\beta) \\ &\quad + \bar{\zeta}_{\dot{\beta}}^I (\sigma^m)_{\alpha}^{\dot{\beta}} (\nabla_m \phi - F_{mn} \Omega^n) \end{aligned}$$

String central charge current

$$\zeta_3 = \frac{1}{2} \partial_m ((\phi^a \bar{\Omega}^m - \bar{\phi}^a \Omega^m) B_3^a) \sigma_{\alpha\dot{\alpha}}^3 \delta^{IJ} = \frac{i}{2} B_3^a \partial_\varphi (\phi^a \bar{\epsilon} - \bar{\phi}^a \epsilon) \sigma_{\alpha\dot{\alpha}}^3 \delta^{IJ}$$

yields for a string of tension \sim epsilon

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} |B_z^a + \phi \tau^a \bar{\phi} - i \nabla_m (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)|^2 + \frac{1}{2} |\mathcal{D}_1 \phi^a + i \mathcal{D}_2 \phi^a - (\Omega_2 - i \Omega_1) B_z^a|^2 \\ &\quad + \partial_m (B_z^a (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)) \geq \partial_m (B_z^a (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)). \end{aligned}$$

Vortices in Omega background [PK Gorsky Chen] in progress

Symmetry breaking pattern

$$SU(2)_c \times SU(2)_R \times SU(2)_{\mathcal{R}} \rightarrow U(1)_c \times SU(2)_{R+\mathcal{R}}$$

SUSY transform
pure SYM

$$\begin{aligned} \delta\Lambda_\alpha^I &= \zeta_\beta^I ((\sigma^{mn})_\alpha^\beta F_{mn} + i[\phi, \bar{\phi}] \delta_\alpha^\beta + \nabla_m (\bar{\Omega}^m \phi - \Omega^m \bar{\phi}) \delta_\alpha^\beta) \\ &\quad + \bar{\zeta}_{\dot{\beta}}^I (\sigma^m)_\alpha^{\dot{\beta}} (\nabla_m \phi - F_{mn} \Omega^n) \end{aligned}$$

String central charge
current

$$\zeta_3 = \frac{1}{2} \partial_m ((\phi^a \bar{\Omega}^m - \bar{\phi}^a \Omega^m) B_3^a) \sigma_{\alpha\dot{\alpha}}^3 \delta^{IJ} = \frac{i}{2} B_3^a \partial_\varphi (\phi^a \bar{\epsilon} - \bar{\phi}^a \epsilon) \sigma_{\alpha\dot{\alpha}}^3 \delta^{IJ}$$

yields for a string of tension \sim epsilon

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} |B_z^a + \phi \tau^a \bar{\phi} - i \nabla_m (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)|^2 + \frac{1}{2} |\mathcal{D}_1 \phi^a + i \mathcal{D}_2 \phi^a - (\Omega_2 - i \Omega_1) B_z^a|^2 \\ &\quad + \partial_m (B_z^a (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)) \geq \partial_m (B_z^a (\Omega^m \bar{\phi}^a - \bar{\Omega}^m \phi^a)). \end{aligned}$$

*Searching for the field theoretical explanation of the
new duality*

Conclusions and open questions

- Nonabelian vortices to study BPS spectrum of SQCD
- Generalization of the 4d/2d duality to theories in Omega background
- Connections to integrable systems in 2d...
- Relationship w/ another 4d/2d duality [Vafa et al]
- Holography for Non-Abelian vortices