

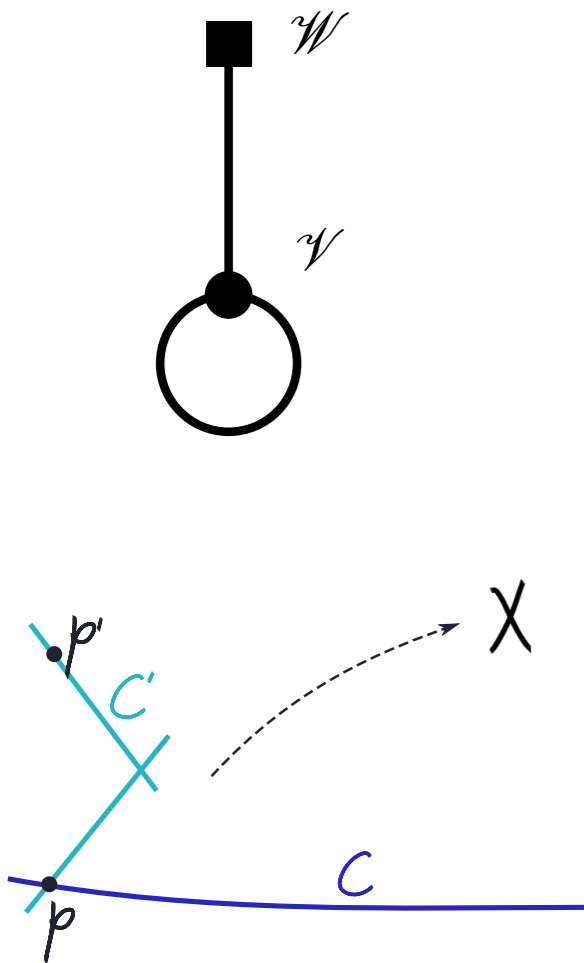
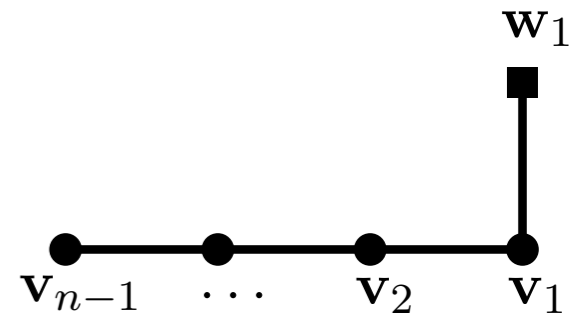
DAHA, Elliptic Hall Algebra

ADHM & $1/2$ ADHM

Peter Koroteev



Talk at workshop 'Higher Symmetries'
Aspen, CO March 19th 2019



Higher Symmetries

In this talk we shall discuss some **higher symmetries** which arise in *physics* and *geometry*.

$$\mathcal{E} \simeq U_{q_1, q_2} \left(\widehat{\widehat{\mathfrak{gl}_1}} \right) \simeq \mathcal{E}_{q_1, q_2} \simeq \mathfrak{gl}_\infty \text{ DAHA}_{q_1, q_2}^S \simeq \text{DIM}_{q_1, q_2} \simeq D(\mathcal{A}_{\text{shuffle}})$$

Quantum Toroidal \mathfrak{gl}_1 Elliptic Hall algebra spherical \mathfrak{gl}_∞ DAHA Ding-Iohara-Miki algebra Drinfeld double of shuffle algebra

This algebra often appears in the **BPS/CFT** correspondence:

Connects BPS observables of $\mathcal{N}=2$ supersymmetric gauge theories with CFT correlators (*Mathematically*: relates structures arising on moduli spaces of sheaves (instantons) with vertex operator algebras or qVOAs)

BPS/CFT and Geometry

Mathematicians have now several *proofs* of BPS/CFT (AGT) in limiting cases (no fundamental matter), those proofs do not use the original class-S construction [Schiffmann Vasserot] [Negut]

Physics **proof*** by Kimura and Pestun uses direct localization computations

One of our goals is to understand BPS/CFT **geometrically**

Namely we want describe instanton counting and vertex operator algebras in terms of **quantum geometry** (quantum cohomology or quantum K-theory) of some family of spaces

In other words we want (q)VOAs to **emerge** from quantum geometry

Recent Developments

Vertex Algebras at the Corner [Gaiotto Rapcak]

VOAs at junctions of supersymmetric intersections in N=4 SYM

COHA and VOAs [Rapcak Soibelman Yang Zhao]

Action of COHA on the moduli space of *spiked* instantons

Quiver W-algebras [Kimura Pestun]

4,5,6d quiver gauge theories on $R^4 \times S$ in Omega background

The Magnificent Four [Nekrasov]

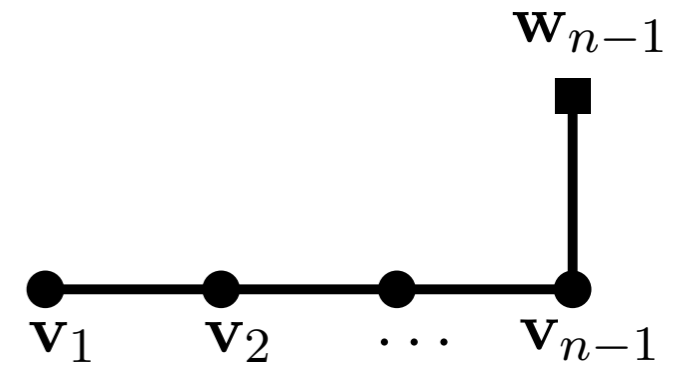
D8 brane probed by D0 branes in B field

$U(1)^4 \subset \text{Spin}(8)$ + additional *nongeometric* $U(1)$ symmetry
 q_1, q_2, q_3, q_4

Nakajima Quiver Varieties

$\text{Rep}(\mathbf{v}, \mathbf{w})$ — linear space of quiver reps

$\mu : T^*\text{Rep}(\mathbf{v}, \mathbf{w}) \rightarrow \text{Lie}(G)^*$ moment map



Nakajima quiver variety

$$X = \mu^{-1}(0) // G$$

$$G = \prod GL(V_i)$$

Automorphism group

$$\text{Aut}(X) = \prod GL(Q_{ij}) \times \prod GL(W_i) \times \mathbb{C}_{\hbar}^{\times}$$

Maximal torus

$$T = \mathbb{T}(\text{Aut}(X))$$

Tensorial polynomials of tautological bundles V_i, W_i and their duals generate *classical T-equivariant K-theory* ring of X

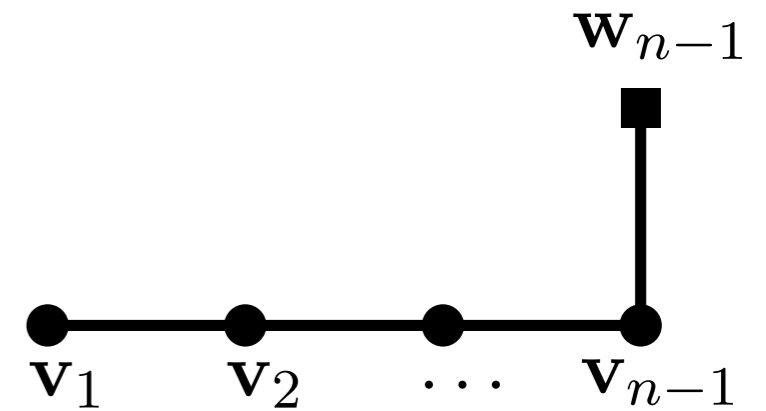
value of a quasimap defines a map to a quotient stack which contains stable locus as an open subset

Quasimaps

$\mathcal{C} \rightarrow X$ is described by collection of vector bundles \mathcal{V}_i on \mathcal{C} of ranks \mathbf{v}_i with section $f \in H^0(\mathcal{C}, \mathcal{M} \oplus \mathcal{M}^* \otimes \mathfrak{h})$ satisfying $\mu = 0$

where $\mathcal{M} = \sum_{i \in I} \text{Hom}(\mathcal{W}_i, \mathcal{V}_i) \oplus \sum_{i, j \in I} Q_{ij} \otimes \text{Hom}(\mathcal{V}_i, \mathcal{V}_j)$

Degree $(\mathbf{v}_1, \dots, \mathbf{v}_{n-1})$



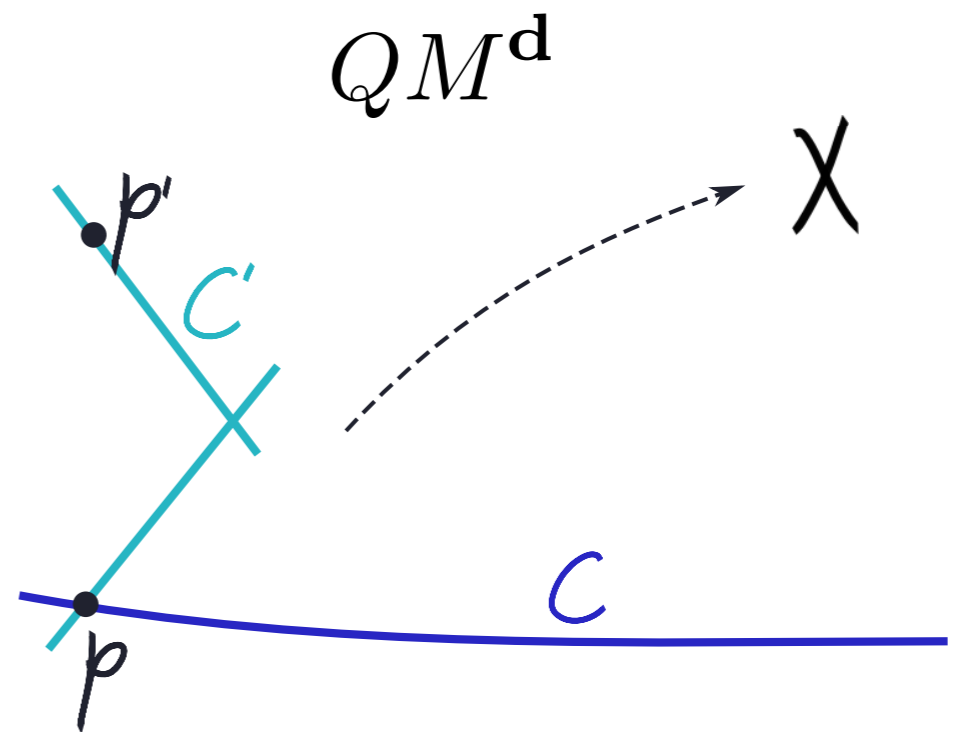
Evaluation map

$$\text{ev}_p(f) = f(p) \in [\mu^{-1}(0)/G] \supset X$$

Stable if $f(p) \in X$

for all but finitely many singular points

Resolve to make proper ev map



Vertex Function (g)

Say this in words: equivariant pushforward, etc.

Spaces of quasimaps admit an action of an extra torus \mathbb{C}_q base \mathbb{P}^1 keeping two fixed points (0, infinity)

Define **vertex function** with quantum (Novikov) parameters $z^{\mathbf{d}} = \prod_{i \in I} z_i^{d_i}$

$$V^{(\tau)}(z) = \sum_{\mathbf{d}=\vec{0}}^{\infty} z^{\mathbf{d}} \text{ev}_{p_2,*} \left(\text{QM}_{\text{nonsing } p_2}^{\mathbf{d}}, \hat{\mathcal{O}}_{\text{vir}} \tau(\mathcal{Y}_i|_{p_1}) \right) \in K_{\mathbb{T}_q}(X)_{\text{loc}}[[z]]$$

[Okounkov]

[PK Pushkar Smirnov Zeitlin]

Define **quantum K-theory** as a ring with multiplication

$$A \circledast B = A \otimes B + \sum_{d=1}^{\infty} A \circledast_d B z^d$$

$$\mathcal{F} \circledast = \sum_{\mathbf{d}=\vec{0}}^{\infty} z^{\mathbf{d}} \text{ev}_{p_1,p_3,*} \left(\text{QM}_{p_1,p_2,p_3}^{\mathbf{d}}, \text{ev}_{p_2}^* (\mathbf{G}^{-1} \mathcal{F}) \hat{\mathcal{O}}_{\text{vir}} \right) \mathbf{G}^{-1} \quad \left(\overbrace{\quad}^{\mathbf{G}^{-1} \mathcal{F}} \right) \mathbf{G}^{-1}$$

gluing

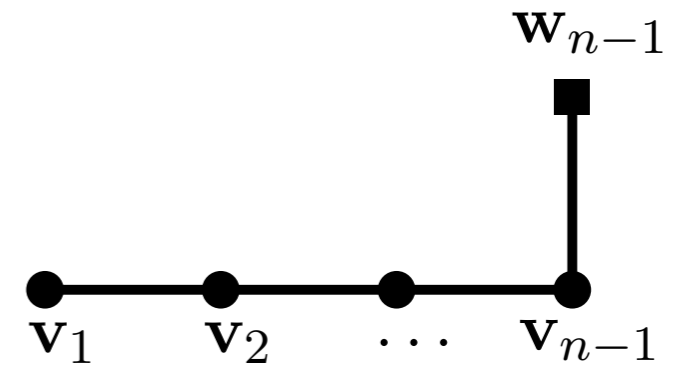
$$\mathcal{C}_0 = \mathcal{C}_{0,1} \cup_p \mathcal{C}_{0,2} \quad \text{---} = \text{X} = \text{---} \mathbf{G}^{-1} \left(\text{---} \right)$$

Vertex Functions

After classifying fixed points of space of nonsingular quasimaps we can compute the vertex

$$V_p^{(\tau)}(z) = \sum_{d_{i,j} \in C} z^{\mathbf{d}} q^{N(\mathbf{d})/2} EHG \tau(x_{i,j} q^{-d_{i,j}})$$

$$E = \prod_{i=1}^{n-1} \prod_{j,k=1}^{v_i} \{x_{i,j}/x_{i,k}\}_{d_{i,j}-d_{i,k}}^{-1} \quad x_{i,j} \in \{a_1, \dots, a_{w_n}\}$$



Vertex (trivial class)

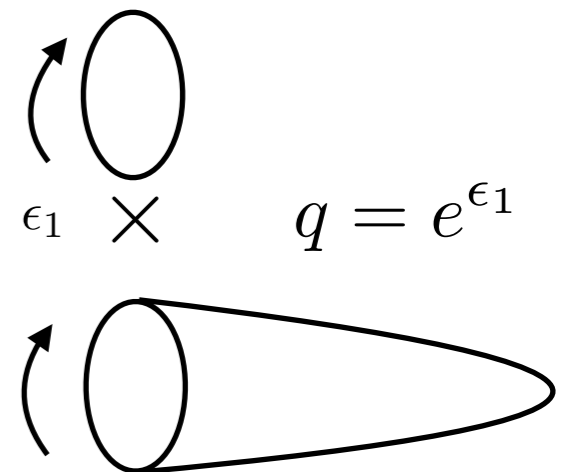
$$V = {}_2\phi_1 \left(\hbar, \hbar \frac{a_1}{a_2}, q \frac{a_1}{a_2}; q; z \right) \quad v_1 = 1, w_1 = 2$$

Vortex

$\mathcal{N} = 2^*$ quiver gauge theory on $X_3 = \mathbb{C}_{\epsilon_1} \times S^1_\gamma$

Lagrangian depends on twisted masses a_1, a_2

FI parameter z and $U(1)$ R-symmetry $\log \hbar$

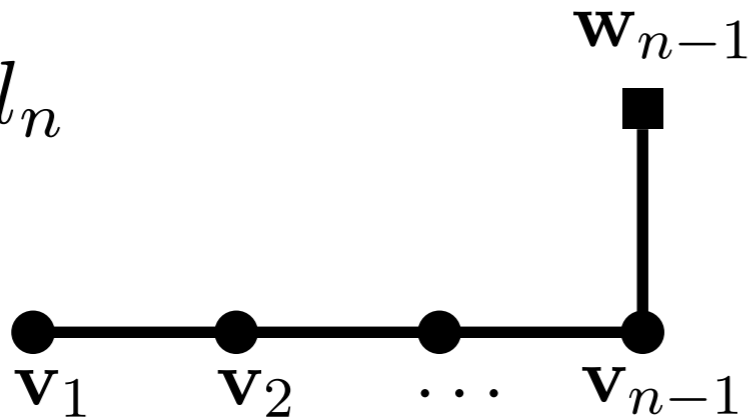


Difference Equations

[PK Pushkar Smirnov Zeitlin]

[PK, PK Zeitlin]

$$X = T^* \mathbb{F}l_n$$



Quantum K-theory Ring $q \rightarrow 0$

$$QK_T(T^* \mathbb{F}l_n) = \frac{\mathbb{C}[z_i^{\pm 1}, a_i^{\pm 1}, T_q^{\pm 1} \hbar, q]}{\mathcal{I}_{\text{tRS}}}$$

The K-theory vertex function satisfies equation of motion of *trigonometric Ruijsenaars-Schneider model*

$$\hat{H}_d V = e_d(z_1, \dots, z_{n-1}) V$$

$$\hat{H}_d = \sum_{I \subset \{1, \dots, n\}, |I|=d} \left(\prod_{i \in I, j \notin I} \frac{a_i \hbar^{\frac{1}{2}} - a_j \hbar^{-\frac{1}{2}}}{a_i - a_j} \right) \prod_{i \in I} T_i^q$$

3d Mirror version (a.k.a. bispectral dual)

$$\hat{H}_d^! V = e_d(a_1, \dots, a_{n-1}) V$$

$$\hat{H}_d^!(a_i, \hbar, T_a^q) = \hat{H}_d(z_i/z_{i+1}, \hbar^{-1}, T_z^q)$$

If time will be tight then say in words that A and B are holonomies of electric and magnetic operators

Should be around 50% of time here!!!!

Spherical DAHA

Deift-Schneider Hamiltonians form a maximal commuting subalgebra inside **spherical double affine Hecke algebra for $\mathfrak{gl}(n)$**

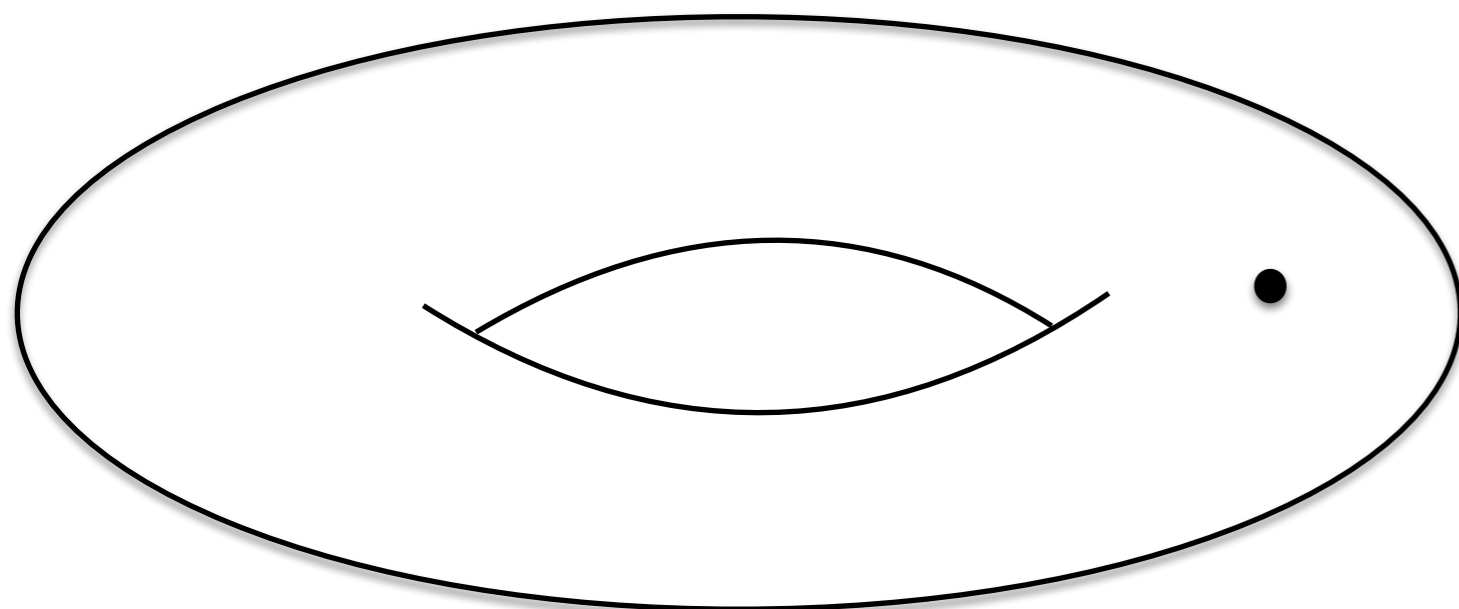
$$\{\hat{H}_1, \dots, \hat{H}_n\} \subset \text{DAHA}_{q, \hbar}^{\mathfrak{S}_n}(\mathfrak{gl}_n) =: \mathcal{A}_n$$

\hat{H}_d are also known as Macdonald operators

[Oblomkov]

Spherical $\mathfrak{gl}(n)$ DAHA is a **geometric quantization** of the moduli space of flat $GL(n; \mathbb{C})$ connections on a torus with one simple puncture

$$\mathcal{M}_n = \{A, B, C\} / GL(n; \mathbb{C})$$



$$ABA^{-1}B^{-1} = C$$

$$C = \text{diag}(\hbar, \dots, \hbar, \hbar^{1-n})$$

$$\mathcal{A}_n = \widehat{\mathbb{C}}_J[\mathcal{M}_n]$$

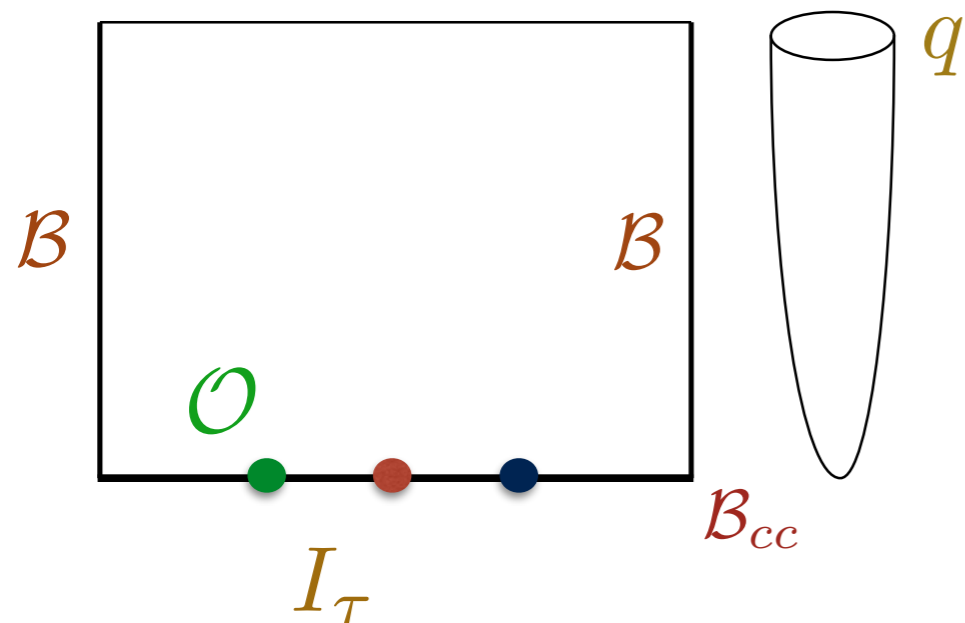
Line Operators and Branes

\mathcal{M}_n is the moduli space of vacua in $\mathcal{N}=2^*$ gauge theory on $\mathbb{R}^3 \times S^1$ with gauge group $U(n)$ and is described by VEVs of line operators wrapping the circle.

A and B are holonomies of *electric* and *magnetic* line operators

Omega background along real 2-plane $\mathbb{R}_q^2 \times \mathbb{R} \times S^1$

Line operators are forced to stay at the tip of the cigar and slide along the remaining line, hence **non-commutativity**



algebra — open strings

$$\mathcal{A}_n = \text{Hom}(\mathcal{B}_{cc}, \mathcal{B}_{cc})$$

representations

(Hilbert space of SUSY QM)

$$\mathcal{H} = \text{Hom}(\mathcal{B}_{cc}, \mathcal{B})$$

[Gukov-Witten]
[Nekrasov-Witten]

Hitchin Moduli Space (n=2)

[Gukov]

[PK Gukov Nawata Saberi]

SU(2) theory \longrightarrow sl(2) flat connections

$x = \text{Tr} A$
electric

$y = \text{Tr} B$
magnetic

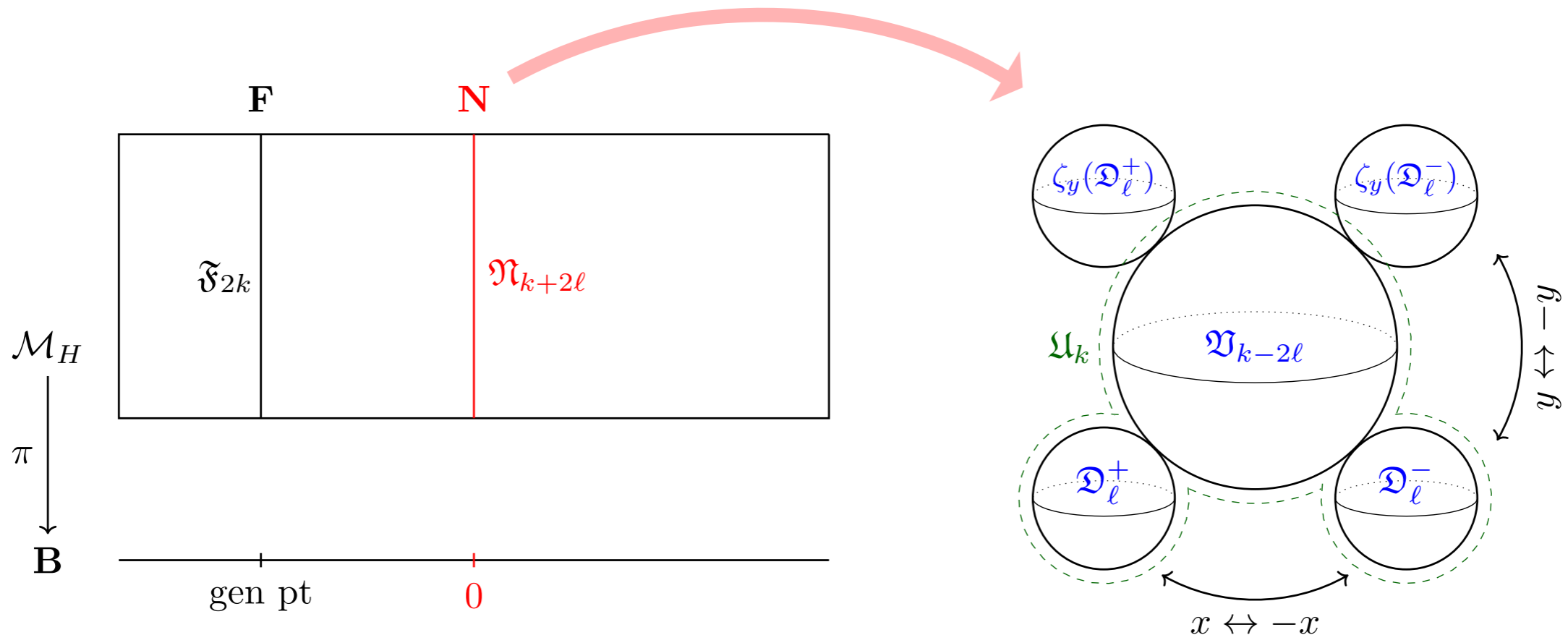
$z = \text{Tr} AB$
dyonic

Nonabelian Hodge
correspondence:

$$\mathcal{M}_{\text{flat}}(SL(2; \mathbb{C}), T^2 \setminus \{\text{pt}\}) \simeq \mathcal{M}_H(SU(2), T^2 \setminus \{\text{pt}\})$$

$$\mathcal{M}_H : x^2 + y^2 + z^2 + xyz = \hbar + \hbar^{-1} + 2 \quad \text{for } \hbar=1 \quad \mathcal{M}_n \simeq \frac{\mathbb{C}^\times \times \mathbb{C}^\times}{\mathbb{Z}_2}$$

Elliptic fibration with one singular fiber of Kodaira type I_0^*

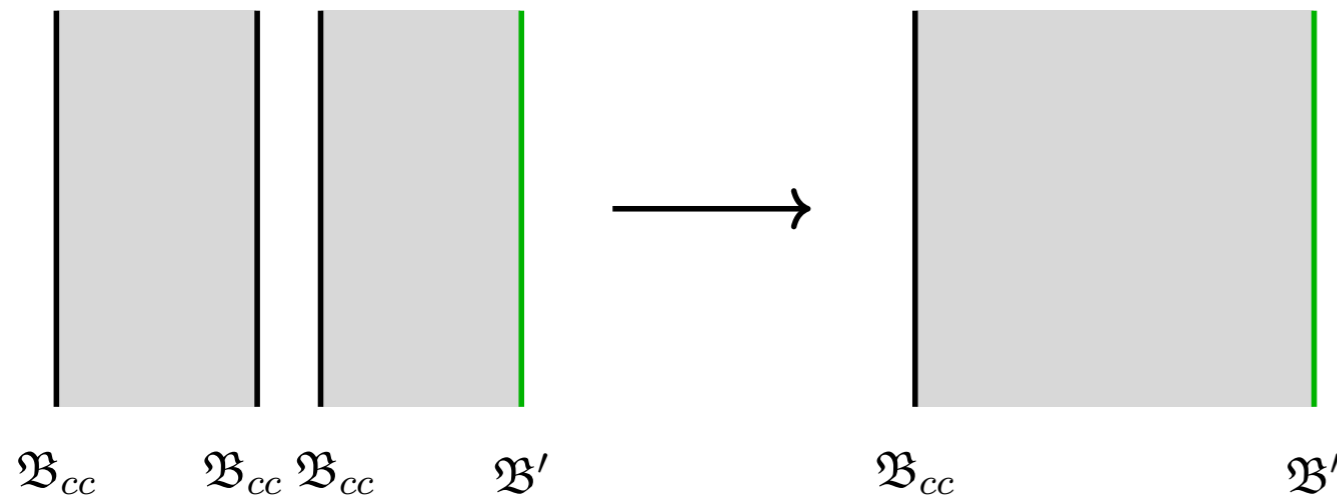


DAHA Modules

[Kapustin Orlov]
[Kapustin Witten]

Algebra acts naturally by attaching open strings to closed strings

Hilbert space comes from $(\mathcal{B}_{cc}, \mathcal{B}')$ strings



$$\mathcal{A} = \text{Hom}(\mathcal{B}_{cc}, \mathcal{B}_{cc})$$

$$\mathcal{H} = \text{Hom}(\mathcal{B}_{cc}, \mathcal{B}')$$

$\mathcal{B}' \rightarrow \text{Hom}(\mathcal{B}_{cc}, \mathcal{B}')$ gives a functor $\text{Hom}(\mathcal{B}_{cc}, \cdot)$

$$\text{Fuk}(\mathcal{M}, \Omega) \simeq \text{Rep}(\mathcal{A})$$

$$\mathcal{B}_{cc} : \mathcal{L} \rightarrow \mathcal{M}_H$$

Algebra-deformation
quantization of functions
on \mathcal{M}_H

$$F + B = \frac{i}{\log q} \Omega_J$$

$$[x, y]_q = (q - q^{-1})z$$

$$[z, x]_q = (q - q^{-1})y$$

$$[y, z]_q = (q - q^{-1})x$$

$$\Omega_J = \frac{dx \wedge dy}{2z - xy}$$

Lagrangian
A-brane



Module
of DAHA

Dimension of a module

$$\dim V = \int_{\mathcal{M}} \text{ch}(\mathcal{B}') \wedge \text{ch}(\mathcal{B}_{cc}) \wedge \text{Td}(\mathcal{M})$$

compact
branes

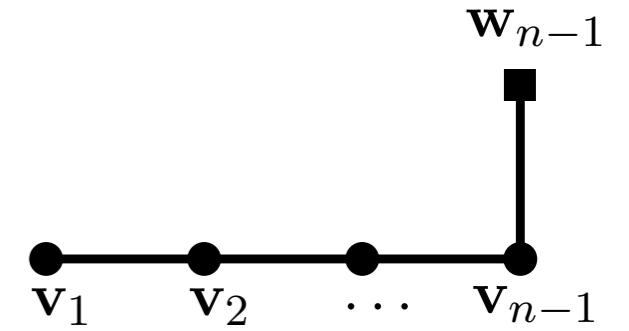


Finite dim
reps

DAHA Reps

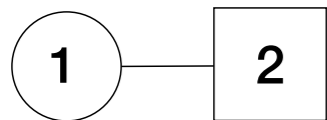
Start with a vertex function for T^*F_n

Specify equivariant parameters $a_k = q^{\lambda_k} \hbar^{n-k}$



q-hypergeometric series \longrightarrow Macdonald polynomials with $\hbar = t^{-1}$

E.g. $k=2, n=2$



$$V(z; \hbar q, q) = P_{(1,1)}(z|q, \hbar)$$

$$V(z; \hbar q^2, q) = P_{(2,0)}(z|q, \hbar)$$

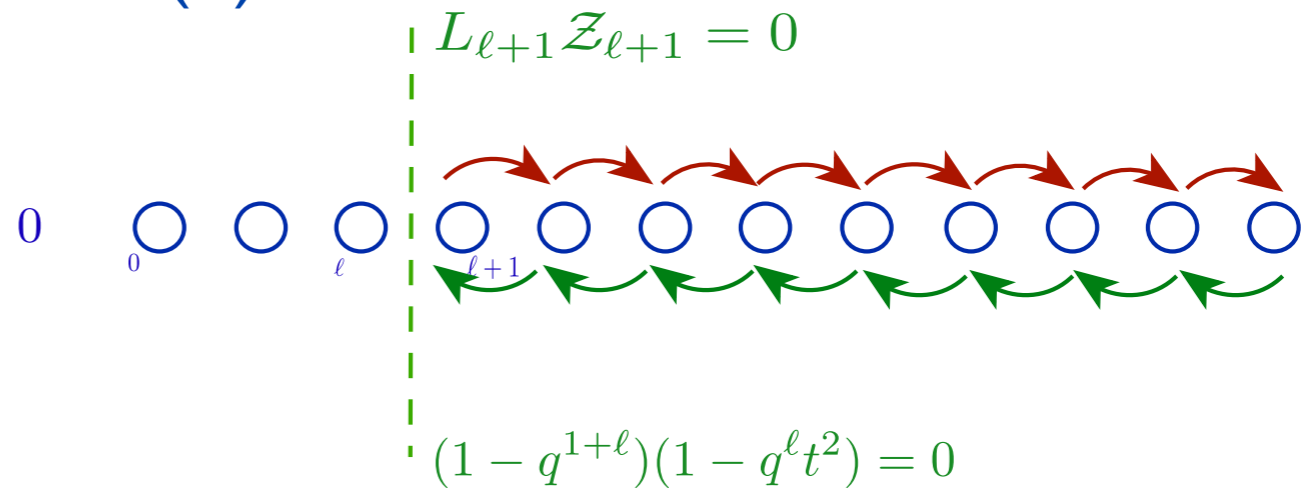
Raising and lowering operators of $\mathfrak{sl}(2)$ DAHA

$$R_a = x + a_k^{-1} z$$

$$L_a = x + a_k z$$

$$R_a \mathcal{Z}_a = r_a \mathcal{Z}_{a+1}$$

$$L_a \mathcal{Z}_a = l_a \mathcal{Z}_{a-1}$$



$$0 \longrightarrow \iota(\mathfrak{D}_k^+ \oplus \mathfrak{D}_k^-) \longrightarrow \mathfrak{U}_N \longrightarrow \mathfrak{Y}_{n+1} \longrightarrow 0$$

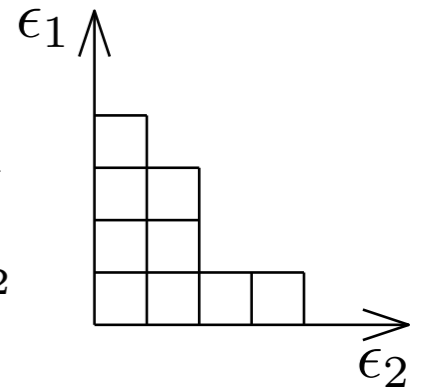
Fock Space

Power-symmetric variables

$$p_m = \sum_{l=1}^n z_l^m$$

$$q = e^{\epsilon_1}$$

$$\hbar = e^{\epsilon_2}$$



Macdonald polynomials depend only on k and the partition

$$P_{\square\square} = \frac{1}{2}(p_1^2 - p_2), \quad P_{\begin{smallmatrix} \square \\ \square \end{smallmatrix}} = \frac{1}{2}(p_1^2 - p_2) + \frac{1 - qt}{(1 + q)(1 - t)} p_2$$

Starting with Fock vacuum

$$|0\rangle$$

Construct Hilbert space

$$a_{-\lambda}|0\rangle \longleftrightarrow p_\lambda$$

for each partition

$$a_{-\lambda}|0\rangle = a_{-\lambda_1} \cdots a_{-\lambda_l}|0\rangle$$

Commutators

$$[a_m, a_n] = m \frac{1 - q^{|m|}}{1 - \hbar^{|m|}} \delta_{m, -n}$$

DAHA Action

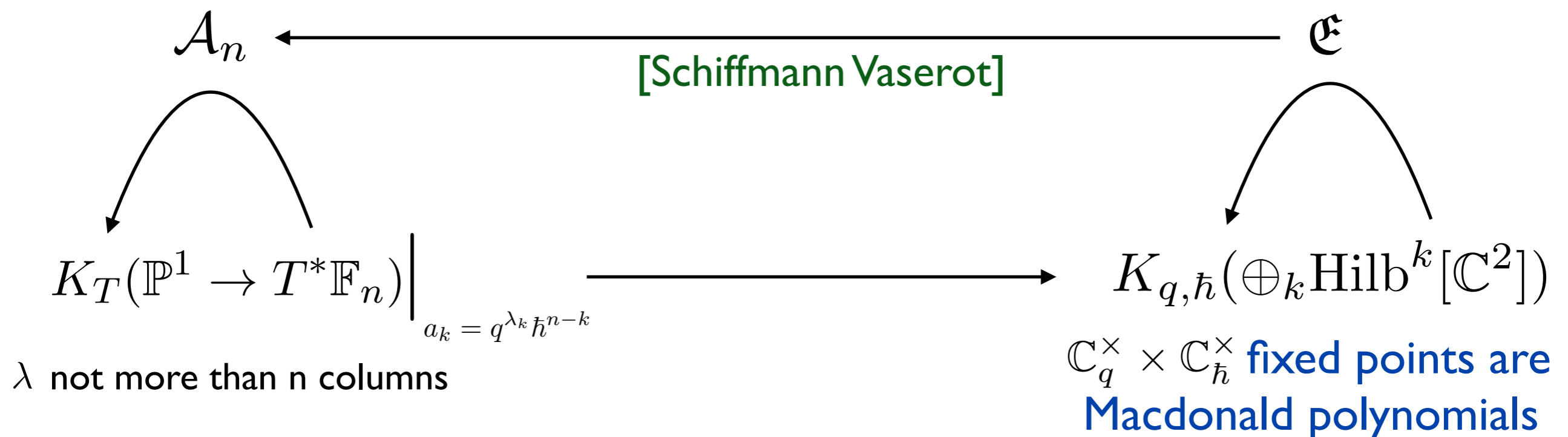
[PK 1805.00986]

Vertex functions or quantum classes for X are elements of quantum K-theory of X . Equivalently we can view them as elements of equivariant K-theory of the space of quasimaps from \mathbb{P}^1 to X

$V \in K_T(\mathbb{P}^1 \rightarrow T^*\mathbb{F}_n)$ with maximal torus $T = \mathbb{T}(U(n) \times U(1)_{\hbar} \times U(1)_q)$.

Specification $a_k = q^{\lambda_k} \hbar^{n-k}$ restricts us to the Fock space representation of (q, \hbar) -Heisenberg algebra which is a DAHA module

In other words, we can define the following action

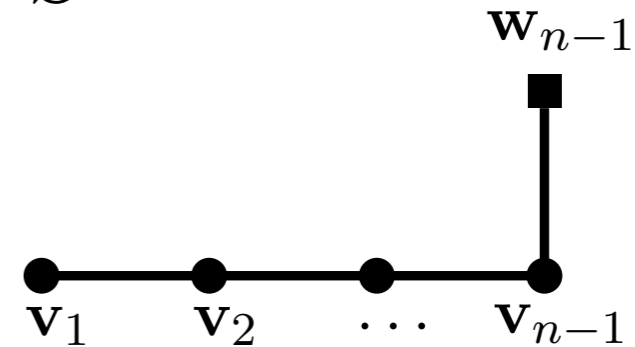


Large-n transition

Recall that $\text{Hilb}^k[\mathbb{C}^2] = \mathcal{M}_{1,k}^{\text{inst}}$ How did U(1) 5d SYM appear?

Starting with M-theory on $S^1 \times \mathbb{C}_q \times \mathbb{C}_{\hbar} \times T^*S^3$
 n M5 branes wrapping $S^1 \times \mathbb{C}_q \times S^3 \subset T^*S^3$

Upon compactification on three sphere will get 3d quiver gauge theory on T^*F_n

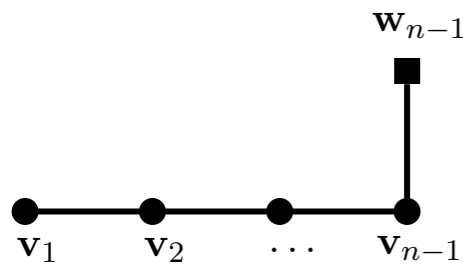


When n becomes large the background undergoes through the **conifold transition** and the *resolved* conifold becomes a *deformed* conifold Y: $S^1 \times \mathbb{C}_q \times \mathbb{C}_t \times Y$

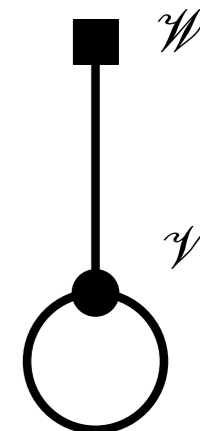
Reduction on Y leads us to a 5d U(1) theory with 8 supercharges

Flags vs ADHM

[PK Sciarappa]



$K_T(\text{QM}(\mathbb{P}^1, X))$	$K_{q, \hbar}(\text{Hilb}(\mathbb{C}^2))$
Kähler/quantum parameters of X z_1, z_2, \dots	Ring generators x_1, x_2, \dots
Vertex function V_q	Classes of $(\mathbb{C}^\times)^2$ fixed points $[\mathcal{J}]$
\mathbb{C}_q^\times acting on base curve	\mathbb{C}_q^\times acting on $\mathbb{C} \subset \mathbb{C}^2$
\mathbb{C}_\hbar^\times acting on cotangent fibers of X	\mathbb{C}_\hbar^\times acting on another $\mathbb{C} \subset \mathbb{C}^2$
Eigenvalues e_r of tRS operators T_r	Chern polynomials \mathcal{E}_r of $\Lambda^r \mathcal{U}$



quantum deformation:

Eigenvalues of **elliptic** RS model at large n

Eigenvalues of **quantum** multiplication by



$$\mathcal{U} = \mathcal{W} + (1 - q)(1 - \hbar)\mathcal{V}|_{\mathcal{J}_{\vec{x}}}$$

$$E_r(\vec{\zeta}) = \sum_{\substack{\mathcal{J} \subset \{1, \dots, n\} \\ |\mathcal{J}|=r}} \prod_{\substack{i \in \mathcal{J} \\ j \notin \mathcal{J}}} \frac{\theta_1(\hbar \zeta_i / \zeta_j | \mathfrak{p})}{\theta_1(\hbar \zeta_i / \zeta_j | \mathfrak{p})} \prod_{i \in \mathcal{J}} p_k$$

Chern roots obey

$$\prod_{l=1}^N \frac{s_a - a_l}{s_a - q^{-1} \hbar^{-1} a_l} \cdot \prod_{\substack{b=1 \\ b \neq a}}^k \frac{s_a - q s_b}{s_a - q^{-1} s_b} \frac{s_a - \hbar s_b}{s_a - \hbar^{-1} s_b} \frac{s_a - q^{-1} \hbar^{-1} s_b}{s_a - q \hbar s_b} = \mathfrak{z}$$

ADHM & 1/2 ADHM

[PK Koroteeva
Gorsky Vainshtein]

$K_{\hbar}(T^*\mathbb{F}l_n) \longleftrightarrow$ ADHM (instanton moduli space)

$$\lim_{n \rightarrow \infty} \left[\hbar^{n-1} (1 - \hbar) \left\langle W_{\square}^{U(n)} \right\rangle \right] \Big|_{\lambda} = a - (1 - q)(1 - \hbar) e_1(s_1, \dots, s_k) \Big|_{\lambda}$$

Claim: $\hbar \rightarrow \infty$ retracting the fibers, dimensional transmutation

[Hanany Tong]

$K(\mathbb{F}l_n) \longleftrightarrow$ 1/2 ADHM (vortex moduli space)

Eigenvalues of **affine**
qToda lattice at large n

Eigenvalues of **quantum**
multiplication by

$$\mathcal{E}_1^{\Lambda}(\lambda) = a - (1 - q) e_1(s_1, \dots, s_k)$$

$$H_1^{\text{aff}} = \mathfrak{p}_1 \left(1 - \mathfrak{p}^{\Lambda} \frac{\partial n}{\partial \mathfrak{z}_1} \right) + \sum_{i=2}^n \mathfrak{p}_i \left(1 - \frac{\partial i-1}{\partial \mathfrak{z}_i} \right)$$

Chern roots obey

$$\prod_{l=1}^N (s_a - a_l) \cdot \prod_{\substack{b=1 \\ b \neq a}}^k \frac{q s_a - s_b}{s_a - q s_b} = \tilde{\mathfrak{p}}^{\Lambda}$$