

Exact simulation of non-Markovian quantum systems with reduced cost

Zhen Huang, 2024/10

with Gunhee Park, Yuanran Zhu, Lin Lin, Chao Yang and Garnet K-L Chan

Physical formulation: G. Park, **Z. Huang**, Y. Zhu, L. Lin, C. Yang, G. K-L Chan, **Physics Review B**, **110**, 195148.

Mathematical theory: **Z. Huang**, L. Lin, G. Park, Y. Zhu, **arXiv: 2411.08741**

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 - Error analysis beyond Gronwall
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Open quantum systems

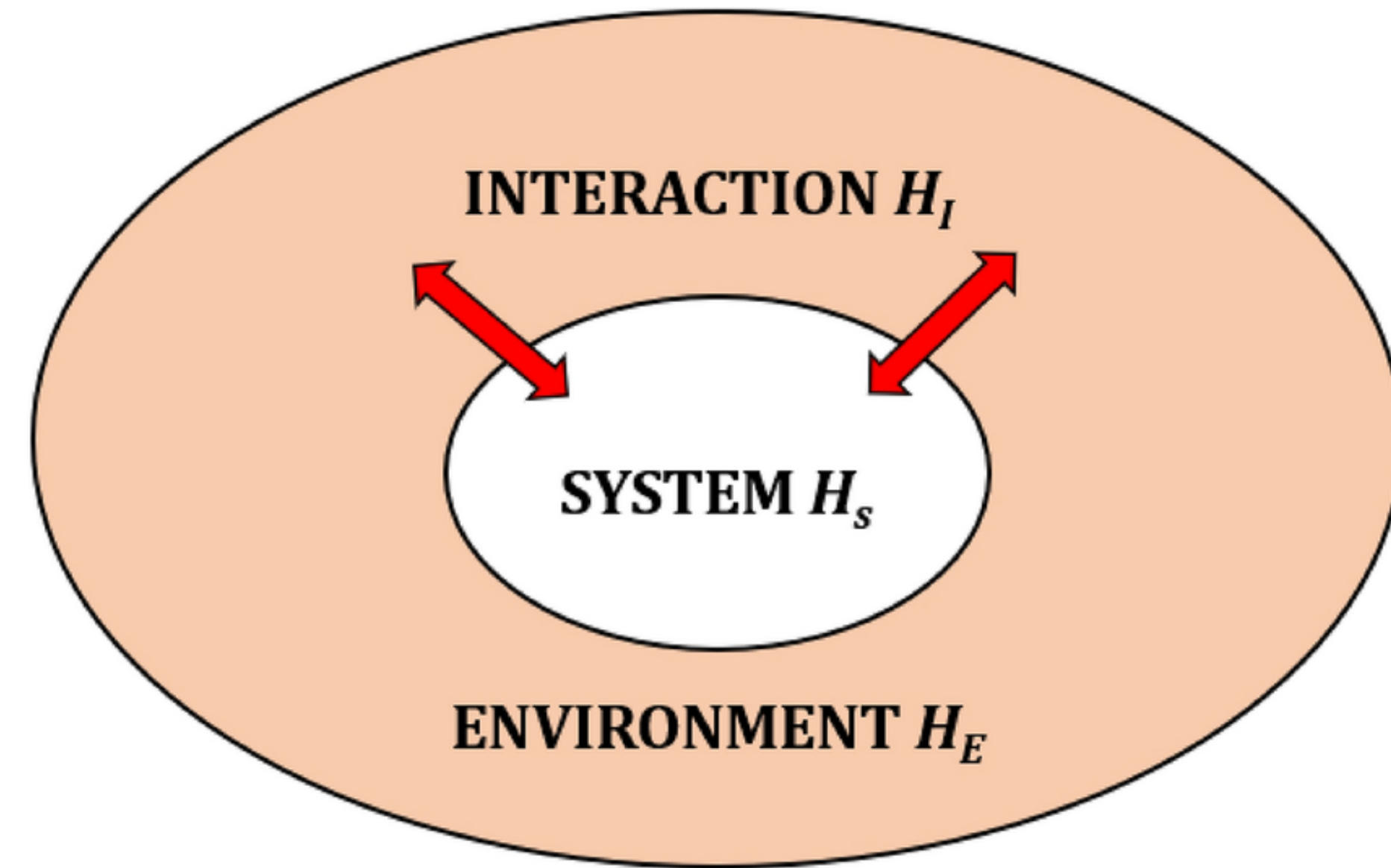
- Open quantum systems:

- $\partial_t \hat{\rho} = -i[\hat{H}_S + \hat{H}_E + \hat{H}_I, \hat{\rho}],$

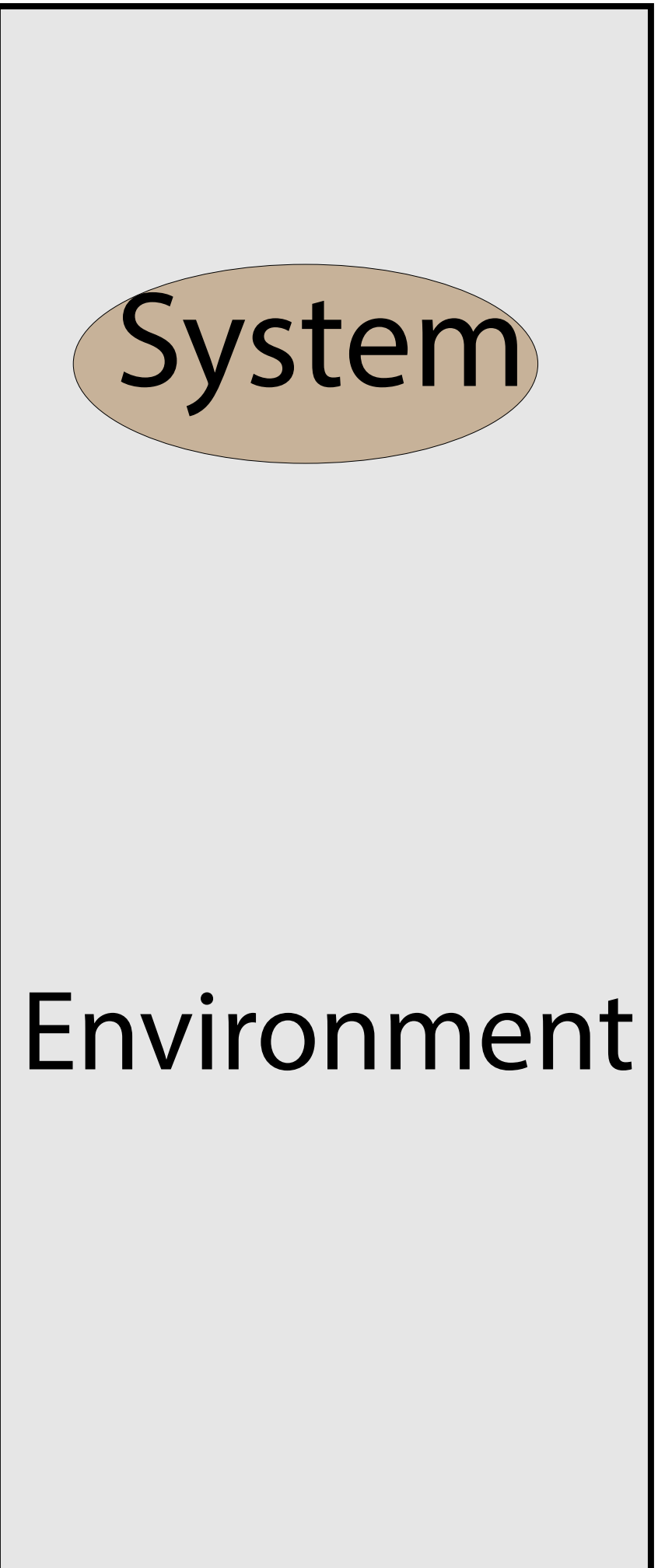
- Goal: obtain $\hat{\rho}_S(t) = \text{Tr}_E(\hat{\rho}(t)).$

- Making separability assumption + truncation + secular approximation, one can obtain Markovian approximation, i.e. Lindblad equation.

- How to simulate the non-Markovian dynamics $\hat{\rho}_S(t)$ exactly?



Theory of open quantum systems

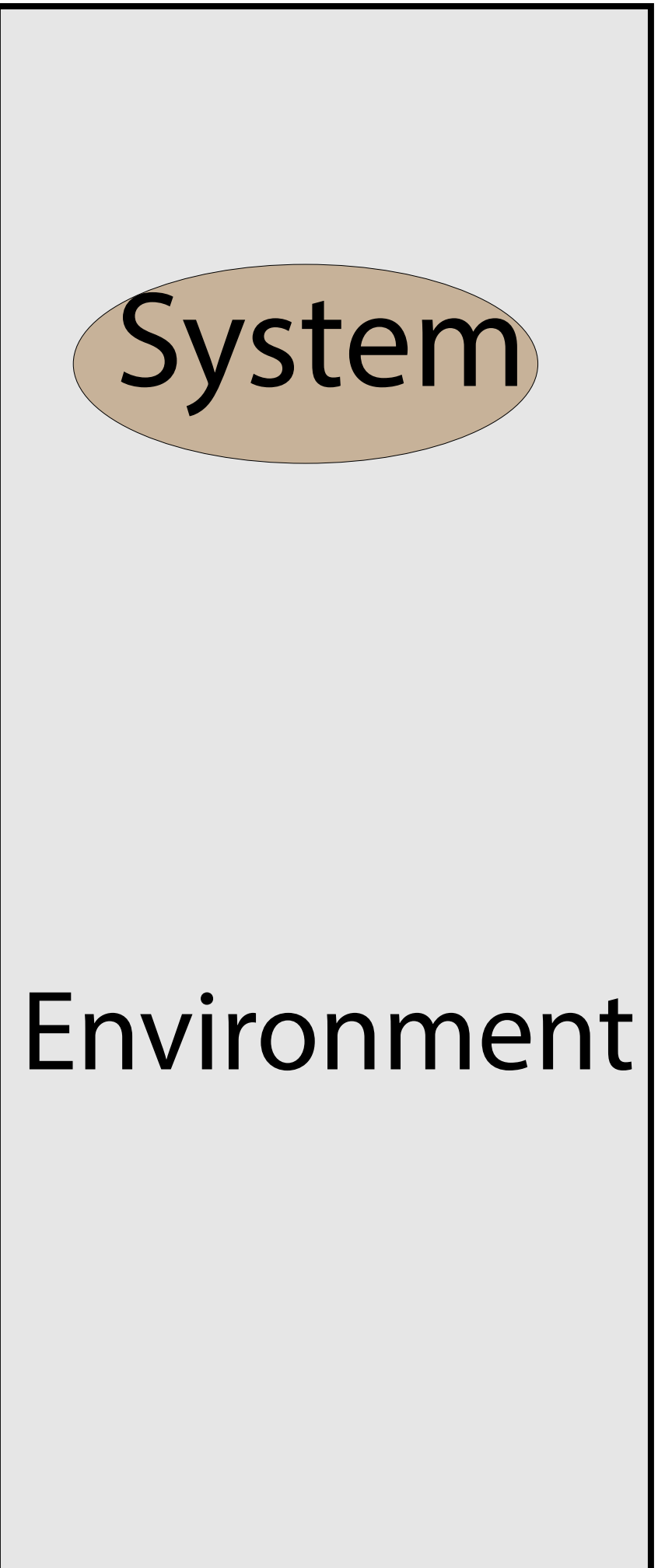


A diagram illustrating the theory of open quantum systems. It features a large, light gray rectangular box with a black border. Inside this box, at the top, is a smaller, light brown oval with a black border containing the word "System". At the bottom of the gray box, the word "Environment" is written in black text. The overall layout is simple and clear, emphasizing the relationship between the system and its environment.

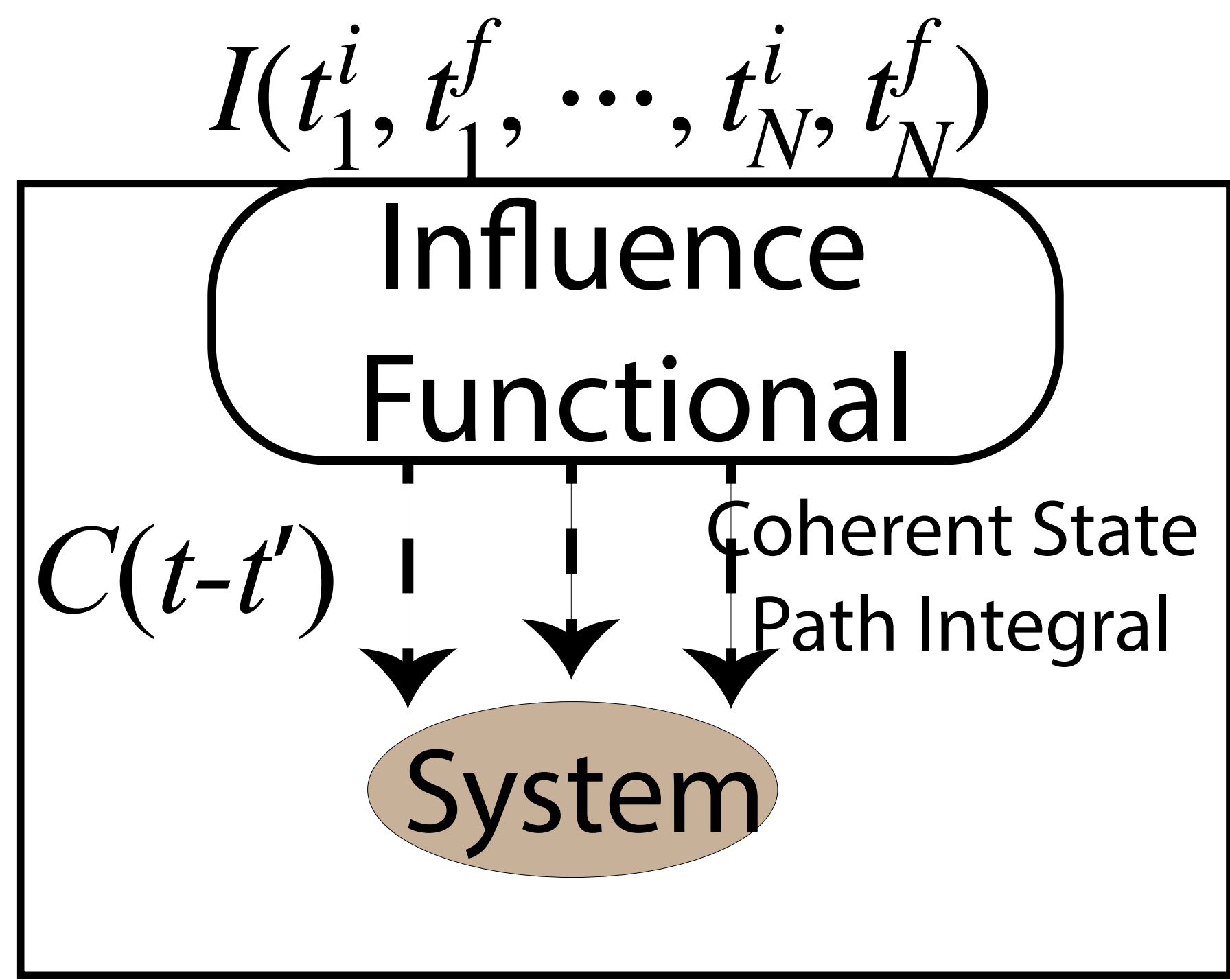
System

Environment

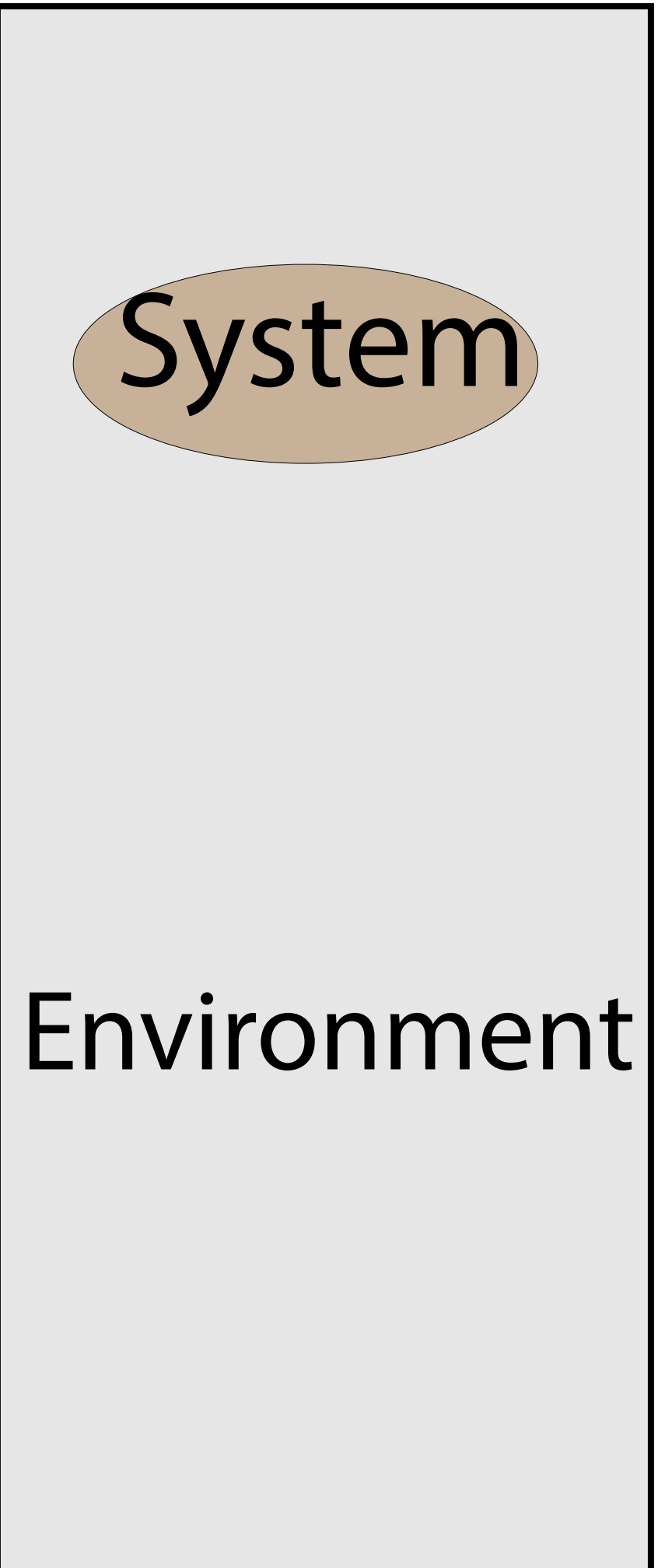
Theory of open quantum systems



Trace out environment exactly

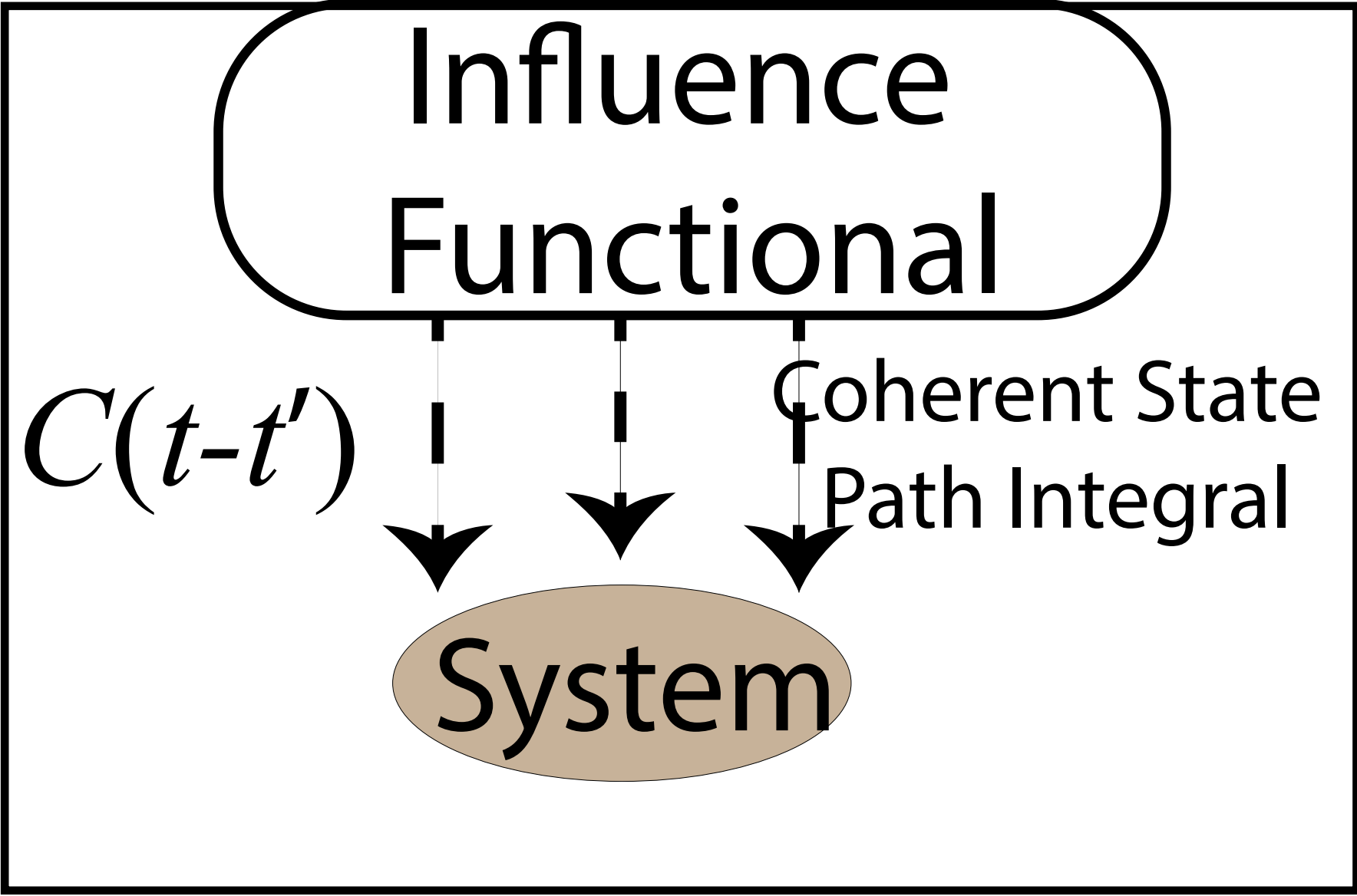


Theory of open quantum systems

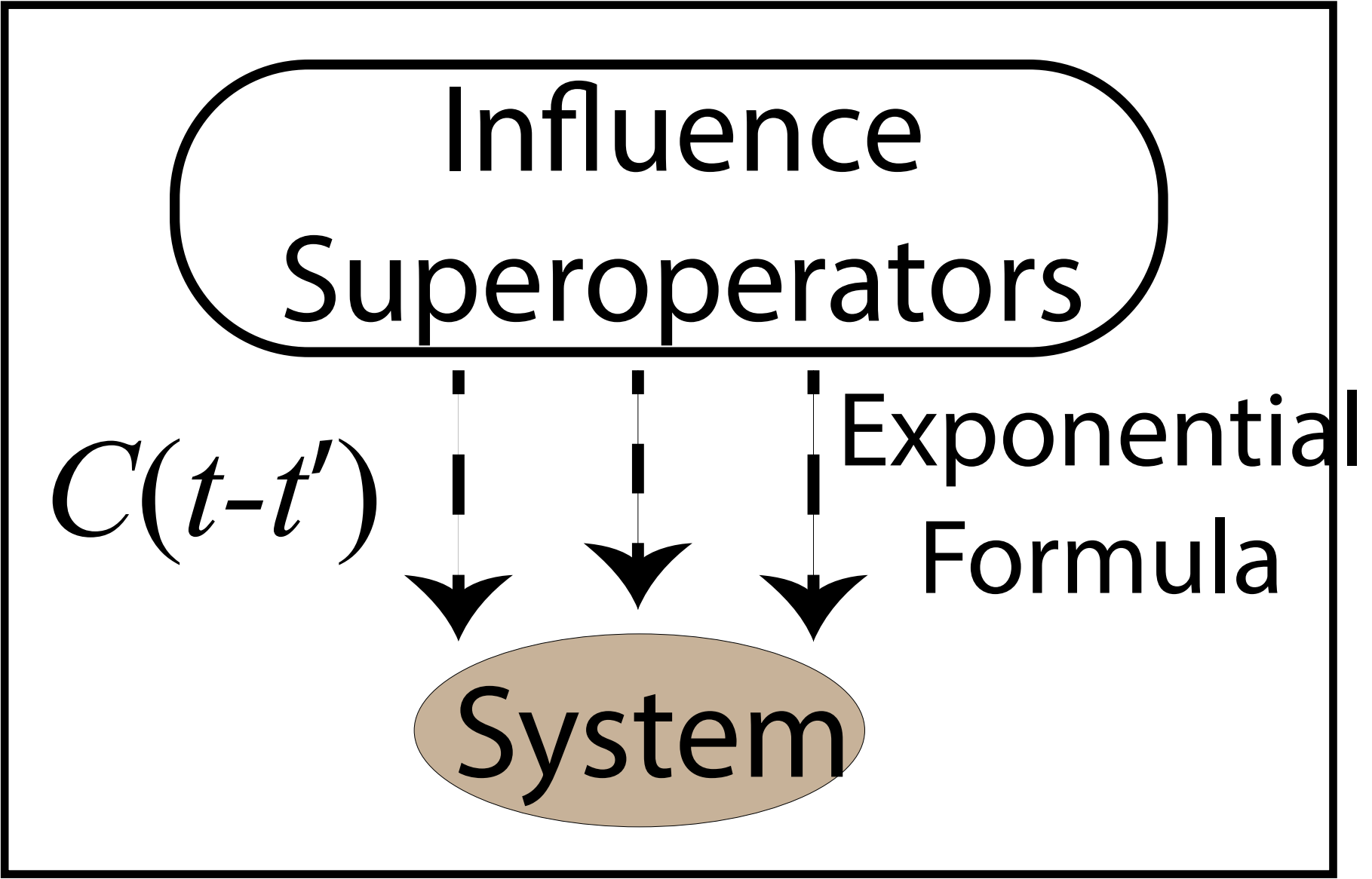


Trace out environment exactly

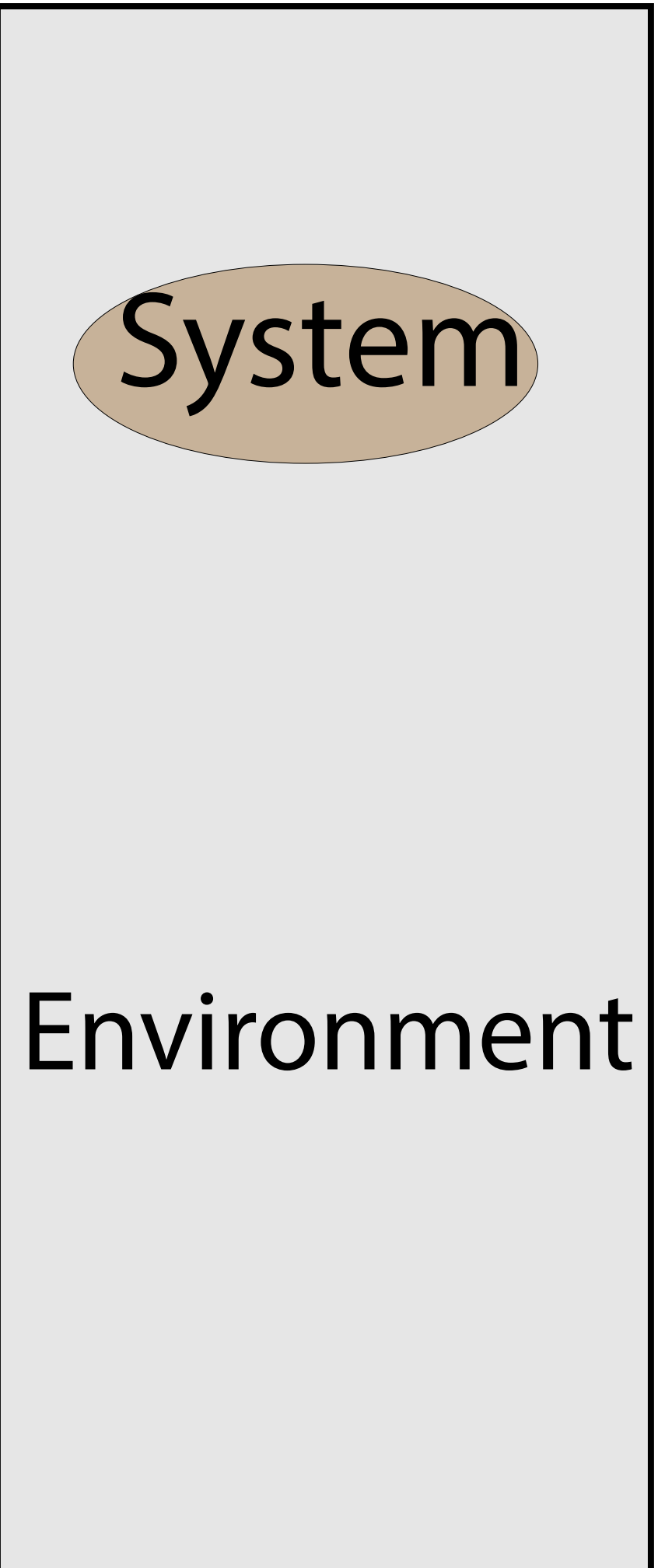
$$I(t_1^i, t_1^f, \dots, t_N^i, t_N^f)$$



Trace out environment exactly

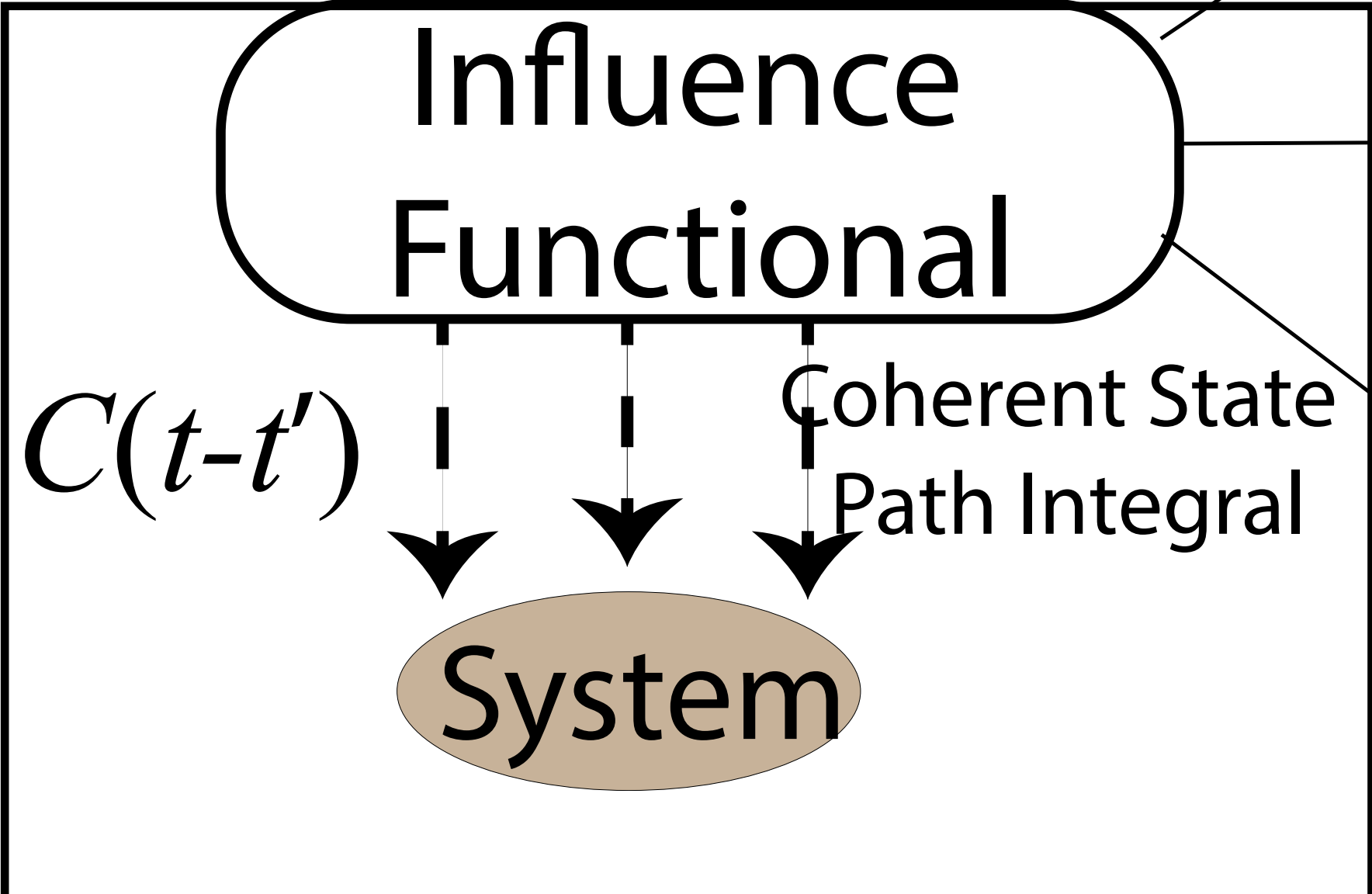


Theory of open quantum systems



Trace out environment exactly

$$I(t_1^i, t_1^f, \dots, t_N^i, t_N^f)$$



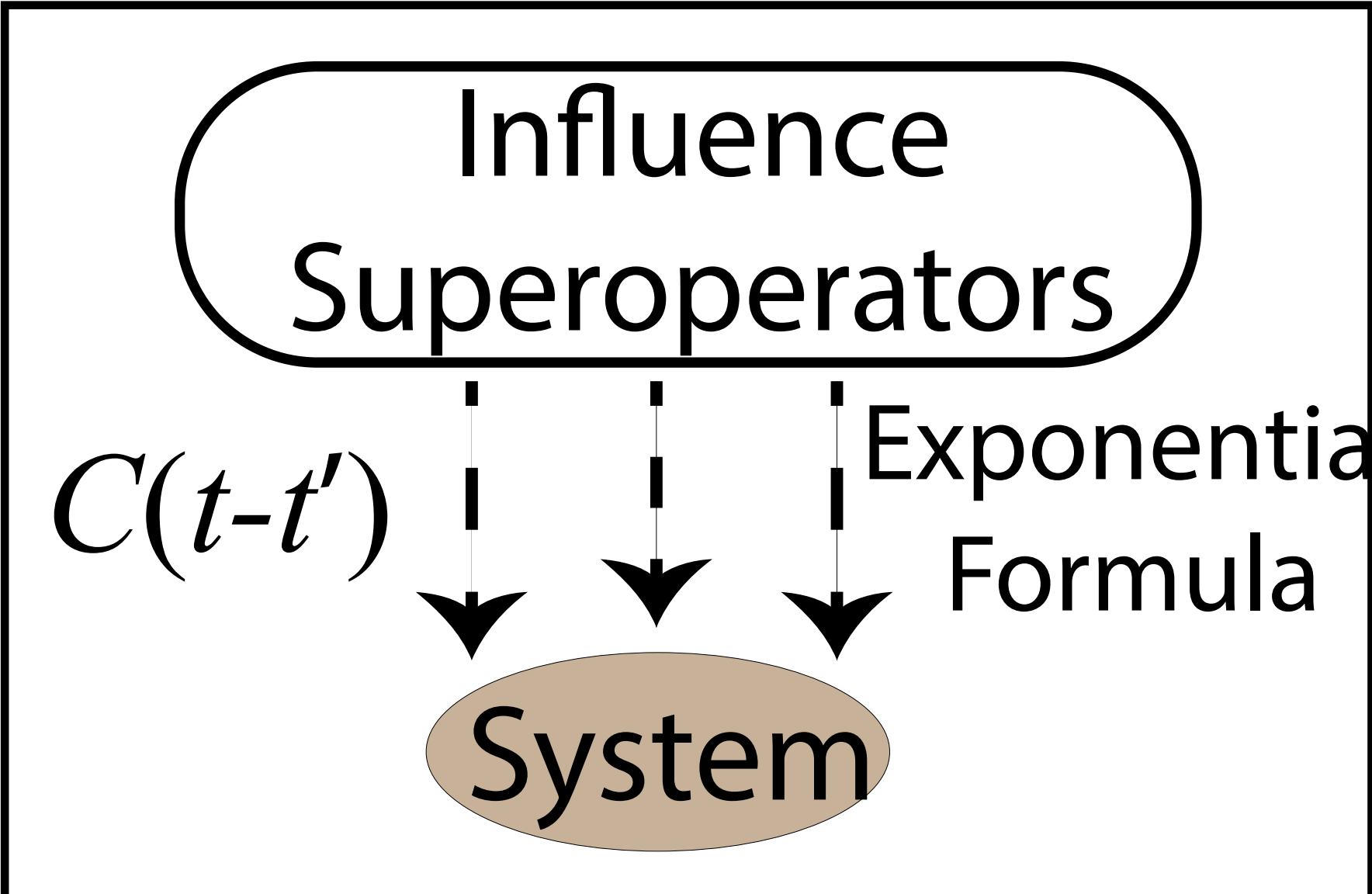
Quasi-adiabatic Path Integral method (QuAPI)

Hierarchical equation of motion (HEOM)

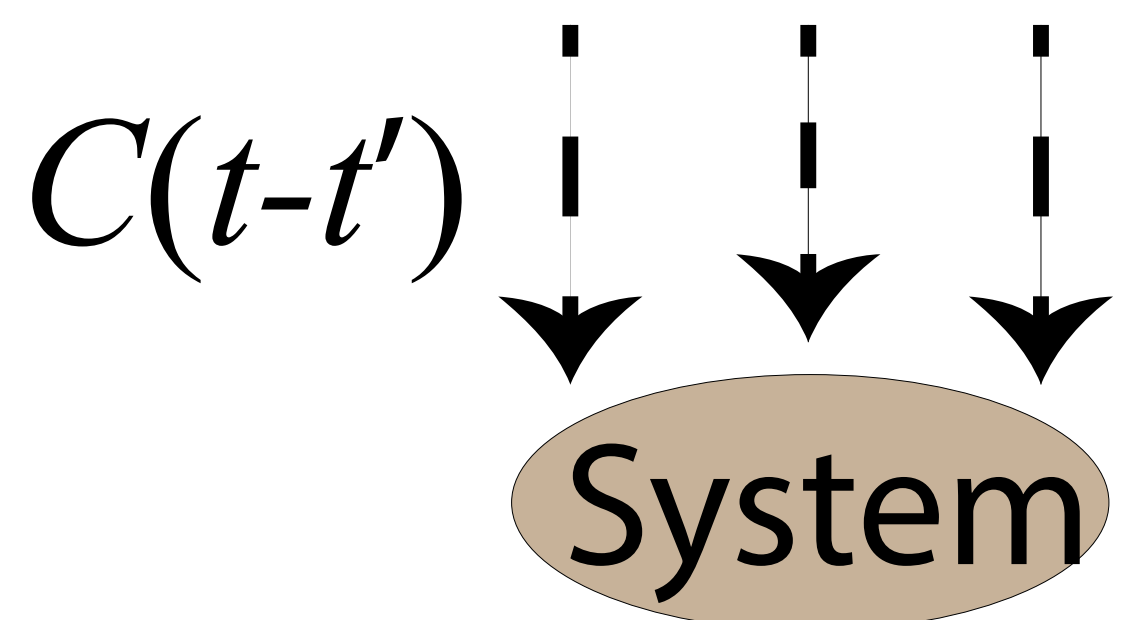
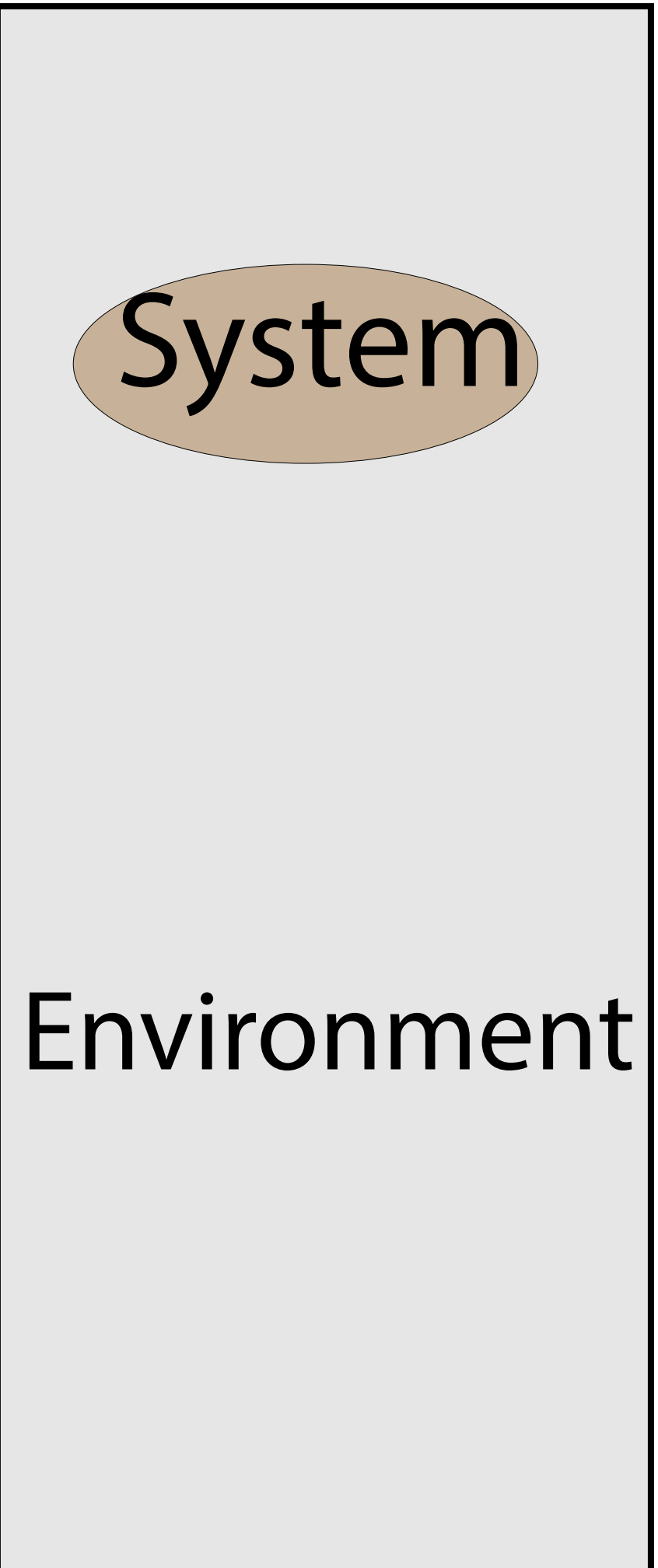
Tensor-network-based Methods

Methods based on approximating influence functional

Trace out environment exactly

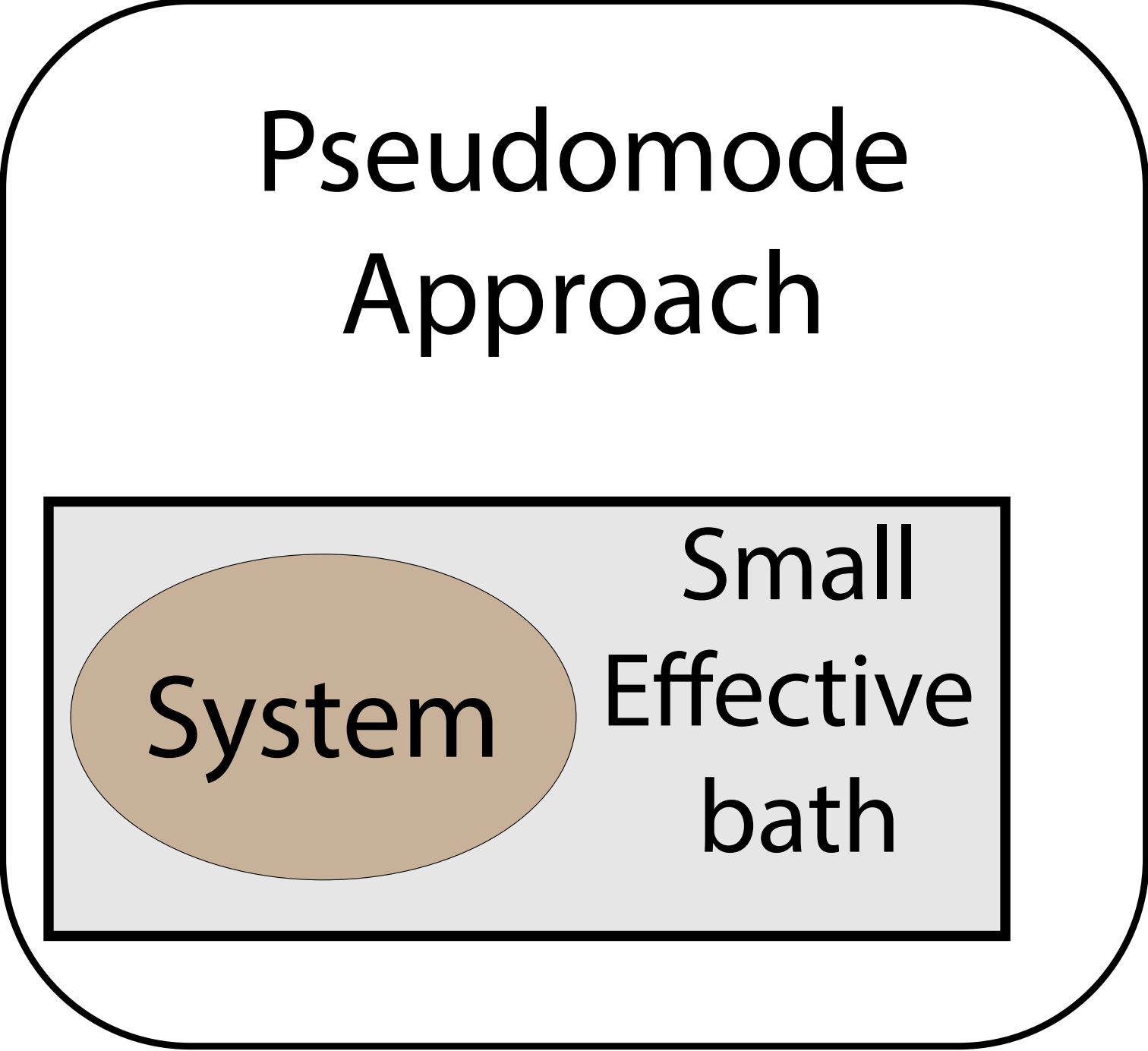


Theory of open quantum systems



The environment's influence on system is only through correlation function $C(t - t')$

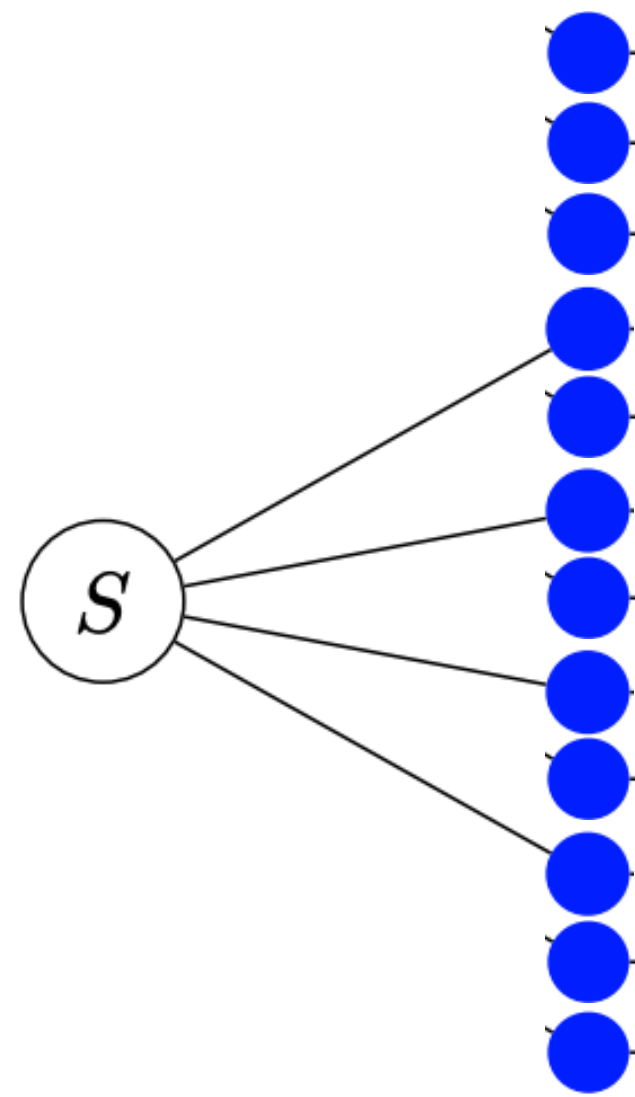
Our approach



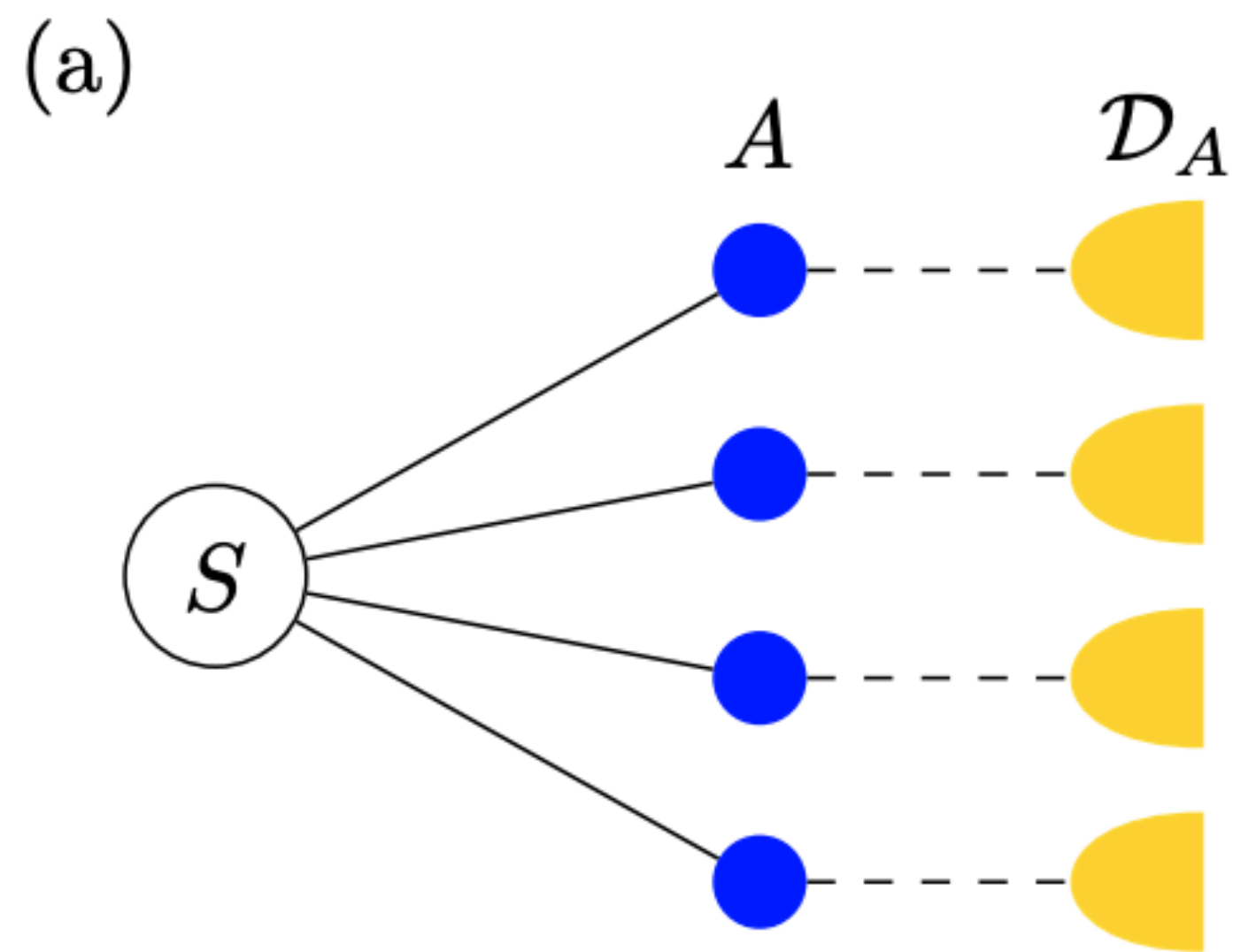
Construct a small effective bath that has the same influence on the system dynamics.

Pseudomode theory

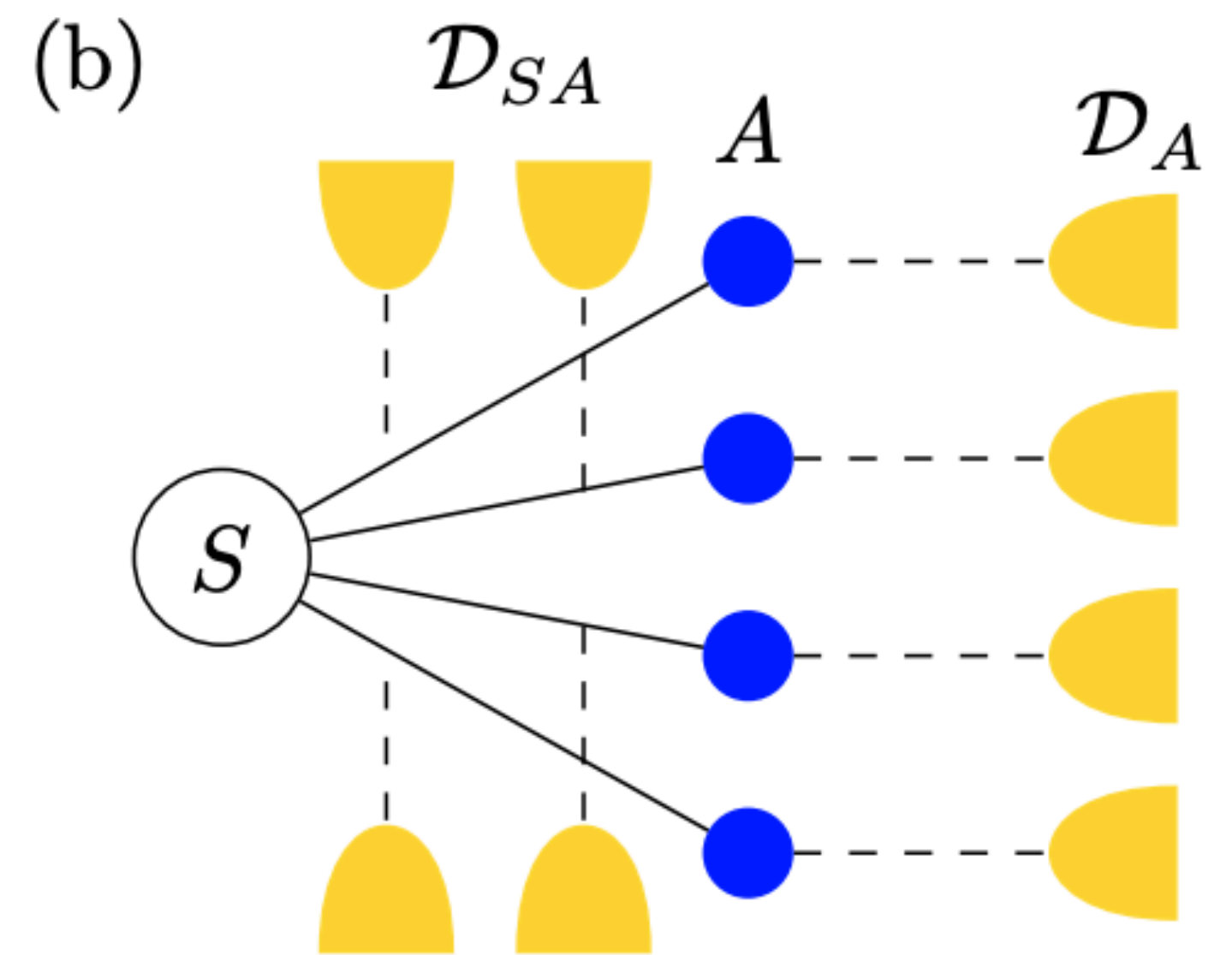
- Pseudomode theory: construct a fictitious/auxiliary bath to replace the original bath.
- Unitary case: discretize the continuous spectrum in some way: equidistance grid, Gauss-Legendre, adaptive quadrature, ...
- Lorentzian pseudomode theory: replace the unitary bath with diagonal Lindbladian bath.
- Quasi-Lindbladian pseudomode theory: replace the unitary bath with Lindbladian bath, replace the unitary system-bath coupling with Lindbladian coupling.
- Conditions for exact simulation: make sure correlation functions are the same!



Original bath

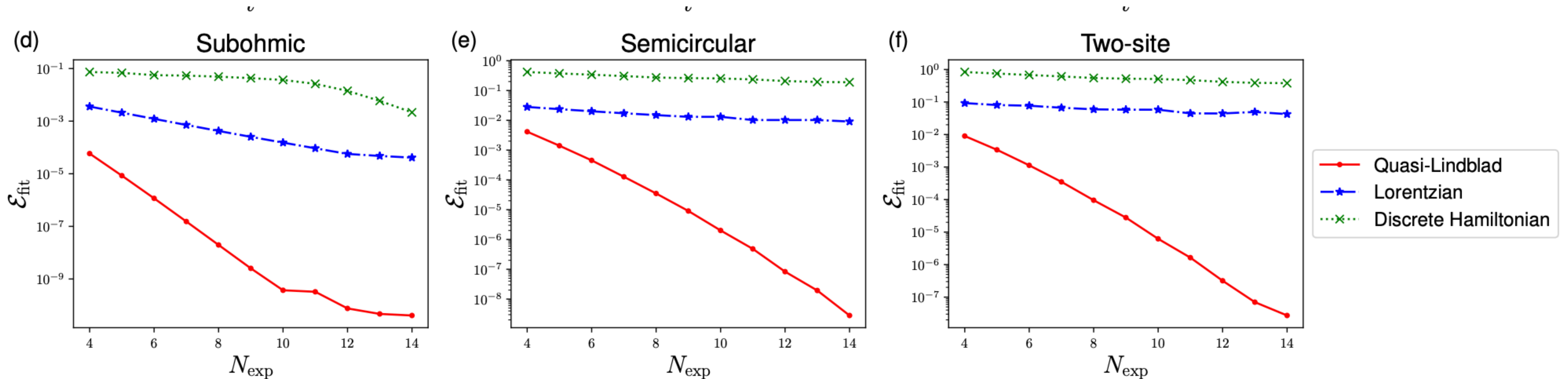


Lindbladian pseudomode theory



quasi-Lindbladian pseudomode theory

Number of baths required for getting accurate correlation functions



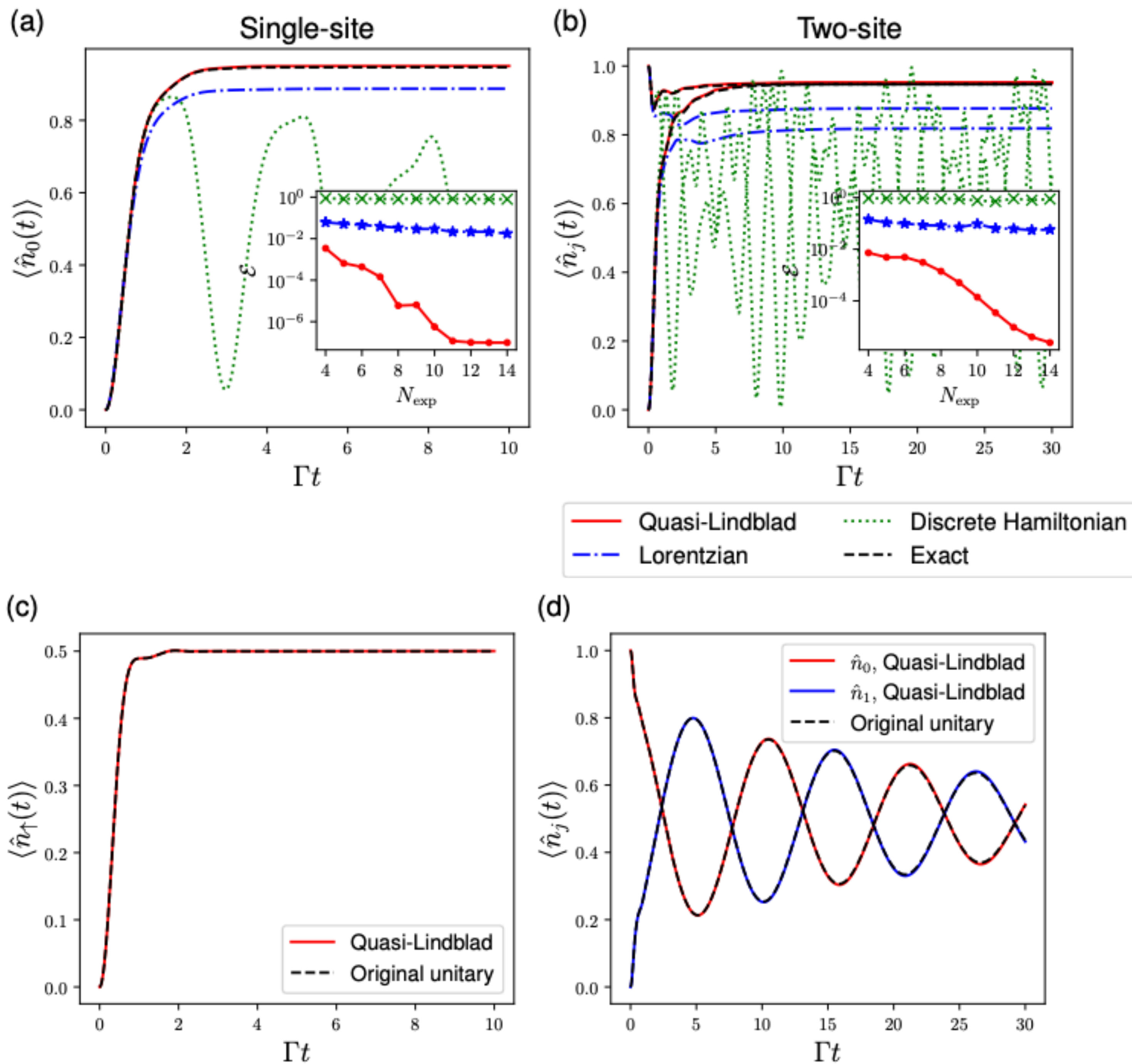
- Number of bath required is much smaller in quasi-Lindblad pseudocode theory, compared to Lorentzian and unitary.
- Similar things, known in the HEOM community (Free pole HEOM).

Quasi-Lindblad theory for spin-boson model

- Original unitary dynamics: $\hat{H}_{\text{phy}} = \hat{H}_S + \hat{H}_B + \hat{H}_{SB} = \hat{H}_S + \int d\omega \omega \hat{c}_\omega^\dagger \hat{c}_\omega + \sum_j \hat{S}_j \hat{B}_j$
- Bath correlation functions: $C_{jj'}^B(t, t') = \text{Tr}_B \left[\hat{B}_j(t) \hat{B}_{j'}(t') \hat{\rho}_B(0) \right]$
- Quasi-Lindblad Pseudomode theory:
- $\partial_t \hat{\rho} = -i[\hat{H}_{\text{aux}}, \hat{\rho}] + D_A(\hat{\rho}) + D_{SA}(\hat{\rho}), \quad \hat{H}_{\text{aux}} = \hat{H}_S + \hat{H}_A + \hat{H}_{SA}$
- $D_A \bullet = 2 \sum_{kk'} \Gamma_{kk'} \left(\hat{F}_{k'} \bullet \hat{F}_k^\dagger - \frac{1}{2} \{ \hat{F}_k^\dagger \hat{F}_{k'}, \bullet \} \right), \quad D_{SA} \bullet = \sum_j \hat{L}'_j \bullet \hat{S}_j + \hat{S}_j \bullet \hat{L}'_j^\dagger - \frac{1}{2} \{ \hat{S}_j \hat{L}'_j + \hat{L}'_j^\dagger \hat{S}_j, \bullet \},$
- Let $L_{SA}(\hat{\rho}) = -i[\hat{H}_{SA}, \hat{\rho}] + D_{SA}(\hat{\rho}), L_A(\hat{\rho}) = -i[\hat{H}_A, \hat{\rho}] + D_A(\hat{\rho}).$
- L_{SA} has the decomposition $\mathcal{L}_{SA} = -i \sum_j \mathcal{S}_j \mathcal{F}_j + i \sum_j \tilde{\mathcal{S}}_j \tilde{\mathcal{F}}_j,$
- Then BCF is $C_{jj'}^A(t, t') = \text{Tr}_A \left[\mathcal{F}_j e^{\mathcal{L}_A(t-t')} \mathcal{F}_{j'} e^{\mathcal{L}_A t'} \hat{\rho}_A(0) \right].$
- Matching condition: $C_{jj'}^B(t, t') = C_{jj'}^A(t, t').$

Result

- Single-site and two-site fermionic Anderson impurity model.



Mathematical foundations (1): Error analysis beyond Gronwall

Lemma 1.1 (Gronwall-type error bound for system observables). Let $C(t - t')$ and $C'(t - t') = C(t - t') + \delta C(t - t')$ be two-point BCFs corresponding to two different environments, and for a bounded system operator \hat{O}_S , let $O_S(t)$, $O'_S(t)$ be the system observables corresponding to the two environments. Then, we have

$$|O_S(t) - O'_S(t)| \leq \|\hat{O}_S\| \epsilon_1 t e^{\mathcal{M}_1 t} \quad \text{for } t \in [0, T] \quad T_0 \sim O\left(\frac{1}{M_1} \log(1/\epsilon_1)\right)$$

Here, $\epsilon_1 = \|\delta C\|_{L^1[0, T]}$, $\mathcal{M}_1 = \|C\|_{L^1[0, T]}$ (see Eq. (2.31)), $\|\hat{O}_S\|$ is the operator norm of \hat{O}_S , and $\|\cdot\|_{L^1[0, T]}$ is the L^1 norm of a function on $[0, T]$.

Theorem 1.1 (Main theorem, for system observables). Let $C(t - t')$, $C'(t - t')$, $O_S(t)$, $O'_S(t)$ and ϵ_1 be the same as in Lemma 1.1. Then, we have

$$|O_S(t) - O'_S(t)| \leq \|\hat{O}_S\| (e^{\epsilon_1 t} - 1), \quad \text{for } t \in [0, T] \quad T \sim O(1/\epsilon_1).$$

Mathematical foundations (2): Stability analysis

- The quasi-Lindblad equation breaks the positivity condition. Yet it appears stable in practice. WHY?
- Turns out this is a coherent-induced stability.
- Recall $\hat{H}_{SA} = \sum_{i \in sys} \sum_{j \in bath} \nu_{ij} \hat{a}_i^\dagger \hat{c}_j + h.c.$
- Theorem (informal): (with some technical assumption), there exists constant m , if $|\nu| \geq m$, the dynamics is stable.
- Rigorous proved for quasi-free case (quadratic Hamiltonian, linear jump).
- Ongoing work: proof for general case.

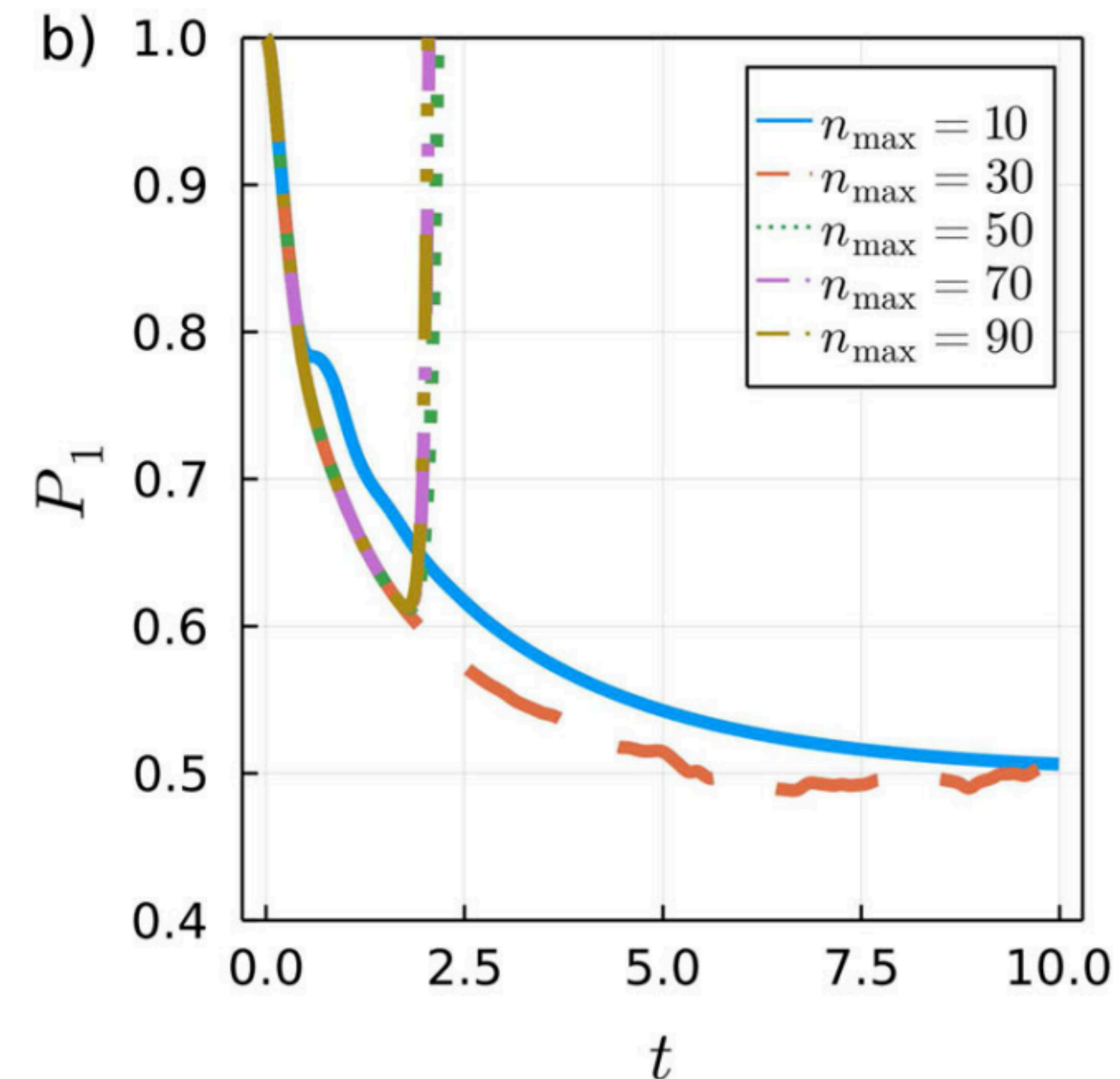
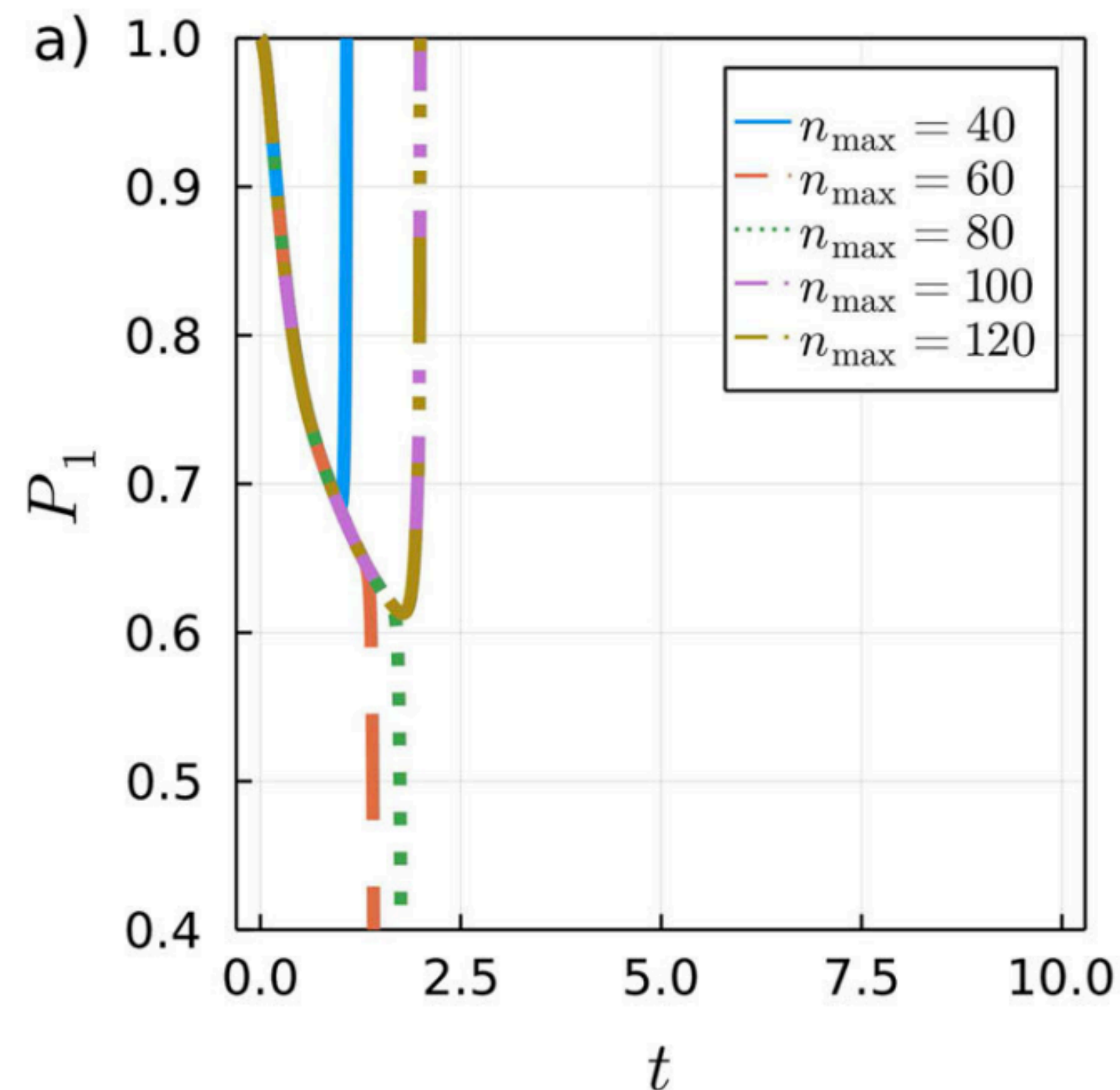


On stability issues of the HEOM method

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- The instability is not improvable by increasing the truncation dimension.

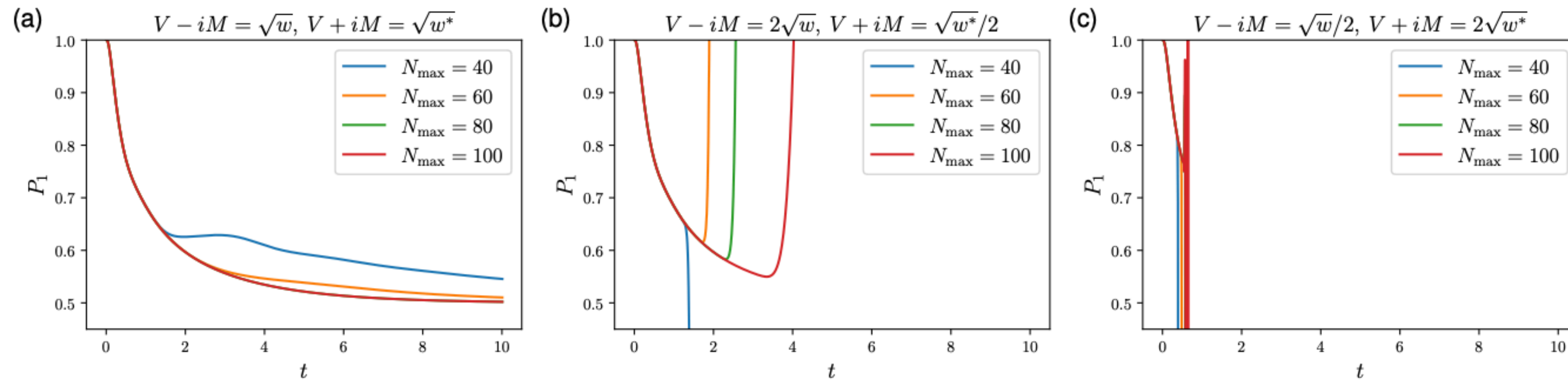
Pseudomode theory could be unstable

Pseudomode theory

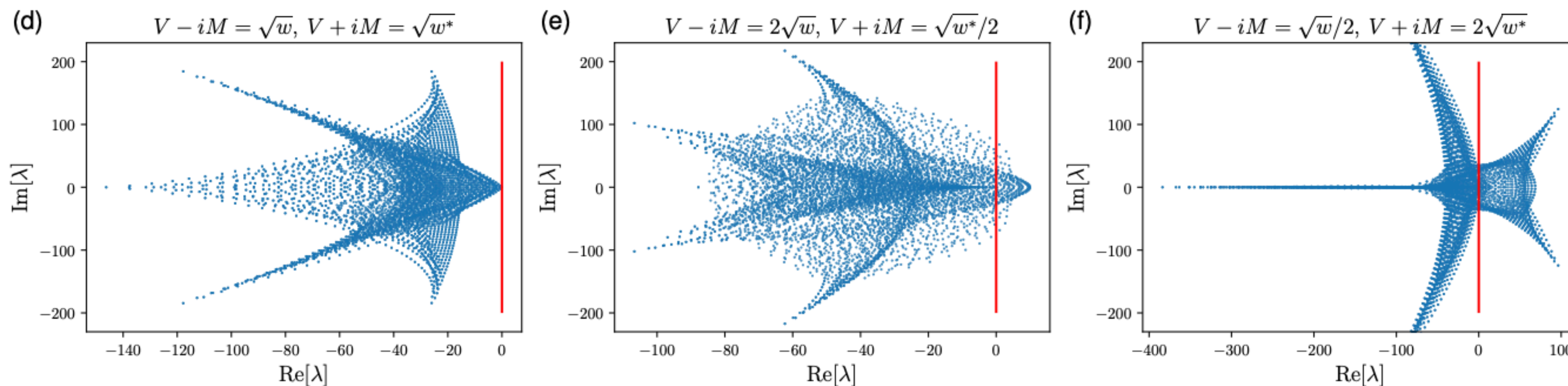
- Various pseudomode theories that break positivity to gain computational advantage.
- Comparison? (Using same parameter settings as the previous HEOM experiment)

Our pseudomode theory

Observables



Eigenvalues of Liouvillians



Discussions

- Mathematics:
 - Stability.
 - Error bound analysis.
- Numerics:
 - Algorithms for solving quasi-Lindblad equations.
- Applications:
 - Nonequilibrium dynamics of density, non-equilibrium Green's function, ...