

Exercises:

1. Let  $T: \mathbb{R}^2 \rightarrow \mathbb{R}^3$  be the smooth function defined by

$$T(z_1, z_2) = (z_1^2 + z_2^2, z_1 z_2, z_1 - z_2^3).$$

For  $\omega = y_2 dy_{(1,3)} \in \Omega^2(\mathbb{R}^3)$ , compute the ascending presentation of  $T^*\omega$ .

2. Let  $\varphi \in C_2(\mathbb{R}^3)$  be defined by

$$\varphi(u) = (\cos(2\pi u_1), \sin(2\pi u_1), u_2).$$

Compute the dipoles  $\delta^1\varphi$  and  $\delta^2\varphi$ , and the boundary  $\partial\varphi$ . Describe these (using words or pictures), making sure to note the proper orientations.

3. Let  $\omega \in \Omega^1(\mathbb{R}^2)$ . Assume  $\omega$  is **closed**:  $d\omega = 0$ .

- (a) For  $p, q \in \mathbb{R}^2$ , let  $\varphi, \psi \in C_1(\mathbb{R}^2)$  be such that  $\varphi(0) = \psi(0) = p$  and  $\varphi(1) = \psi(1) = q$ . Show that  $\int_\varphi \omega = \int_\psi \omega$ .

[**Hint:** consider  $\sigma \in C_2(\mathbb{R}^2)$  defined by  $\sigma(s, t) = (1-s)\varphi(t) + s\psi(t)$ .]

- (b) Fix  $p \in \mathbb{R}^2$ . Define  $h: \mathbb{R}^2 \rightarrow \mathbb{R}$  by  $h(q) := \int_\varphi \omega$  where  $\varphi$  is a 1-cell in  $\mathbb{R}^2$  satisfying  $\varphi(0) = p$  and  $\varphi(1) = q$ . Show that  $h$  is a smooth function with  $dh = \omega$ .

Solutions:

1. According to the formula we proved in class,

$$T^*\omega = y_2 \circ T dT_1 \wedge dT_3 = T_2 dT_1 \wedge dT_3 = z_1 z_2 dT_1 \wedge dT_3.$$

Now

$$\begin{aligned} dT_1 &= \frac{\partial T_1}{\partial z_1} dz_1 + \frac{\partial T_1}{\partial z_2} dz_2 = 2z_1 dz_1 + 2z_2 dz_2 \\ dT_3 &= \frac{\partial T_3}{\partial z_1} dz_1 + \frac{\partial T_3}{\partial z_2} dz_2 = dz_1 - 3z_2^2 dz_2. \end{aligned}$$

Thus using the distributive property and signed commutativity we have

$$\begin{aligned} T^*\omega &= z_1 z_2 (2z_1 dz_1 + 2z_2 dz_2) \wedge (dz_1 - 3z_2^2 dz_2) \\ &= z_1 z_2 (0 - 6z_1 z_2^2 dz_{(1,2)} + 2z_2 dz_{(2,1)} + 0) \\ &= z_1 z_2 (-6z_1 z_2^2 - 2z_2) dz_{(1,2)} \\ &= (-6z_1^2 z_2^3 - 2z_1 z_2^2) dz_{(1,2)} \end{aligned}$$

□

2. We first compute:

$$\begin{aligned} \delta^1\varphi(t) &= \varphi \circ \iota^{1,1}(t) - \varphi \circ \iota^{1,1}(t) \\ &= \varphi(1, t) - \varphi(0, t) \\ &= (\cos(2\pi), \sin(2\pi), t) - (\cos(0), \sin(0), t) \\ &= (1, 0, t) - (1, 0, t), \end{aligned}$$

which is equal to zero even as a formal difference. Let us visualize this formal difference: as  $t$  ranges over  $[0, 1]$ , the first term yields the segment  $[(1, 0, 0), (1, 0, 1)]$  (oriented upwards), while the second term yields the same segment but oriented downwards.

Next we have compute

$$\begin{aligned}\delta^2\varphi(t) &= \varphi \circ \iota^{2,1}(t) - \varphi \circ \iota^{2,0}(t) \\ &= \varphi(t, 1) - \varphi(t, 0) \\ &= (\cos(2\pi t), \sin(2\pi t), 1) - (\cos(2\pi t), \sin(2\pi t), 0).\end{aligned}$$

As  $t$  ranges over  $[0, 1]$ , the first term traces out a circle with center  $(0, 0, 1)$  and radius 1 in the plane  $z = 1$  in the counter-clockwise direction when viewed from above. The second term traces out a circle with center  $(0, 0, 0)$  in the plane  $z = 0$  in the clockwise direction (because of the negative sign).

We then have for the boundary:

$$\partial\varphi = \delta^1\varphi - \delta^2\varphi = 0 - \delta^2\varphi = (\cos(2\pi t), \sin(2\pi t), 0) - (\cos(2\pi t), \sin(2\pi t), 1),$$

which can be visualized exactly as  $\delta^2\varphi$  but with the orientations reversed.  $\square$

3. (a) Letting  $\sigma$  be as in the hint, we have

$$\begin{aligned}\delta^1\sigma(t) &= \sigma \circ \iota^{1,1}(t) - \sigma \circ \iota^{1,0}(t) \\ &= \sigma(1, t) - \sigma(0, t) \\ &= \psi(t) - \varphi(t),\end{aligned}$$

while

$$\begin{aligned}\delta^2\sigma(s) &= \sigma \circ \iota^{2,1}(s) - \sigma \circ \iota^{2,0}(s) \\ &= \sigma(s, 1) - \sigma(s, 0) \\ &= [(1-s)\varphi(1) + s\psi(1)] - [(1-s)\varphi(0) + s\psi(0)] \\ &= [(1-s)q + sq] - [(1-s)p + sp] \\ &= \{q\} - \{p\}.\end{aligned}$$

Now, since  $d\omega = 0$ , we have by the general Stokes' formula:

$$0 = \int_{\sigma} d\omega = \int_{\partial\sigma} \omega = \int_{\delta^1\varphi} \omega - \int_{\delta^2\varphi} \omega = \int_{\psi} \omega - \int_{\varphi} \omega - \int_{\delta^2\varphi} \omega.$$

Since  $\delta^2\varphi$  is a difference of two constant 1-cells, the last term in the last expression above vanishes. Hence  $\int_{\psi} \omega = \int_{\varphi} \omega$  as claimed.  $\square$

- (b) For  $q \in \mathbb{R}^2$ , part (a) implies  $h(q)$  does not actually depend on  $\varphi$  and hence we may replace it with any 1-cell in  $\mathbb{R}^2$ . So for  $q \in \mathbb{R}^2$  consider  $\psi_q \in C_1(\mathbb{R}^2)$  defined by

$$\psi_q(t) = (1-t)p + tq.$$

Note that  $\frac{\partial(\psi_q)_i(t)}{\partial t} = q_i - p_i$ ,  $i = 1, 2$ , where  $p = (p_1, p_2)$  and  $q = (q_1, q_2)$ .

Now, since  $\omega \in \Omega^1(\mathbb{R}^2)$ , we can write  $\omega = f dx + g dy$  for some smooth functions  $f, g: \mathbb{R}^2 \rightarrow \mathbb{R}$ . We then compute

$$\begin{aligned}h(q) &= \int_{\psi_q} \omega = \int_{\psi_q} f dx + g dy \\ &= \int_0^1 f \circ \psi_q(t)(q_1 - p_1) + g \circ \psi_q(t)(q_2 - p_2) dt.\end{aligned}$$

By a theorem from class, we can take partial derivatives under the integral. Hence,  $f$  and  $g$  being smooth functions implies all orders of partial derivatives of  $h$  exist. In particular, for each  $r \in \mathbb{N}$  all the order  $r$  partial derivatives of  $h$  exist and are continuous (since the order  $r + 1$  partials exist), which implies  $h$  is of class  $C^r$ . Since this holds for all  $r \in \mathbb{N}$ ,  $h$  is smooth.

Finally, we compute  $dh$ .

$$\begin{aligned} \frac{\partial h}{\partial q_1} &= \left( \int_0^1 \frac{\partial}{\partial q_1} [f \circ \psi_q(t)(q_1 - p_1)] + \frac{\partial}{\partial q_2} [g \circ \psi_q(t)(q_2 - p_2)] dt \right) dq_2 \\ &= \int_0^1 \frac{\partial f(\psi_q(t))}{\partial x} t(q_1 - p_1) + f \circ \psi_q(t) + \frac{\partial g(\psi_q(t))}{\partial x} t(q_2 - p_2) dt \end{aligned}$$

Note that  $\frac{\partial g}{\partial x} = \frac{\partial f}{\partial y}$  by virtue of  $d\omega = 0$ . Thus we can continue the above computation with

$$\begin{aligned} \frac{\partial h}{\partial q_1} &= \int_0^1 (Df)_{\psi_q(t)}(q - p)t + f \circ \psi_q(t) dt \\ &= \int_0^1 \frac{d}{dt} [f \circ \psi_q(t)t] dt = f \circ \psi_q(1) - 0 = f(q), \end{aligned}$$

where we have invoked the fundamental theorem of calculus in the second-to-last equality. A similar computation yields  $\frac{\partial h}{\partial q_2} = g(q)$ . Thus  $dh = \omega$ .  $\square$