

① Lecture 2 - Deformations of complex structures; Hodge theory

Recall: mirror symmetry predicts existence of mirror pairs of Calabi-Yau manifolds (= complex mfd with $\Omega^{n,0} \cong \mathcal{O}_X$, i.e. $\exists \Omega$ holom. vol. form)

$$X, X^\vee \text{ mirror pair} \Rightarrow H^q(X, \wedge^p T_X) \cong H^q(X^\vee, \Omega^p_{X^\vee})$$

+ isomorphism between "Yukawa couplings" on $H^{1,1}$ and on $H^1(X, T_X)$.

On $H^{1,1}(X)$: $\langle \omega_1, \omega_2, \omega_3 \rangle := \int_X \omega_1 \wedge \omega_2 \wedge \omega_3$

$$+ \sum_{\substack{\beta \in H_2(X, \mathbb{Z}) \\ \beta \neq 0}} \eta_\beta \int_\beta \omega_1 \int_\beta \omega_2 \int_\beta \omega_3 \frac{e^{2\pi i \int_\beta \omega}}{1 - e^{2\pi i \int_\beta \omega}}$$

where η_β = "number of genus 0 complex curves in X representing the homology class β " △
Z
(defined in terms of Gromov-Witten invariants!)

On $H^1(X, T_X) \cong H^{2,1}(X)$: $\langle \theta_1, \theta_2, \theta_3 \rangle := \int_X \Omega \wedge (\theta_1 \cdot \theta_2 \cdot \theta_3 \cdot \Omega)$

where $H^1(X, T_X)^{\otimes 3} \otimes H^0(X, \Omega^3) \rightarrow H^3(X, \wedge^3 T_X \otimes \Omega^3) = H^3(X, \mathcal{O}_X) = H^{0,3}(X)$
 $\theta_1 \otimes \theta_2 \otimes \theta_3$ Ω $\theta_1 \cdot \theta_2 \cdot \theta_3 \cdot \Omega$
 (0,1)-forms in T_X (3,0)-form (3,3)-form w/ values in $\wedge^3 T_X$

or, in fact, thinking of θ_i as deformations of complex structure J
 $\rightarrow \Omega$ changes as we deform $J \rightsquigarrow \langle \theta_1, \theta_2, \theta_3 \rangle = \int_X \Omega \wedge \mathcal{P}_{\theta_1} \mathcal{P}_{\theta_2} \mathcal{P}_{\theta_3} \Omega$

Need of course to assume Ω normalized for this to be ok.

Mirror symmetry prediction: if X & X^\vee are mirror, then

$$\langle \cdot, \cdot, \cdot \rangle \text{ on } H^{1,1}(X) \cong H^1(X^\vee, T_{X^\vee}), \langle \cdot, \cdot, \cdot \rangle$$

match under a certain change of coordinates, induced by mirror map $\mathcal{M}_{\text{Kähler}}(X) \rightarrow \mathcal{M}_{\text{complex}}(X^\vee)$ on tangent spaces

• Focus on: deformation of complex structure, coupling on $H^1(X, T_X)$

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Reference: Gross-Huybrechts-Joyce, "CY moduli & related geometries", ch. 14

- (X, J) almost complex $(J^2 = -1) \rightsquigarrow TX \otimes \mathbb{C} = TX^{1,0} \oplus TX^{0,1}$
 $v^{1,0} = \frac{1}{2}(v - iJv), v^{0,1} = \frac{1}{2}(v + iJv)$
 similarly, $T^*X \otimes \mathbb{C} \simeq T^*X^{1,0} \oplus T^*X^{0,1}$
 $\text{span}(dz_i) \quad \text{span}(d\bar{z}_i)$
 $\Lambda^k T^*X = \bigoplus_{p+q=k} \Lambda^{p,q} T^*X = \Omega^{p,q}(X)$
 relation
 $(TX, J) \simeq TX^{1,0}$ complex vector bundle

- integrability of complex structure $\Leftrightarrow [T^{1,0}, T^{1,0}] \subseteq T^{1,0}$
 $\Leftrightarrow d = \partial + \bar{\partial}$ maps $\Omega^{p,q} \rightarrow \Omega^{p+1,q} \oplus \Omega^{p,q+1}$
 $\Leftrightarrow \bar{\partial}^2 = 0$

then TX and assoc^d bundles are holom. vector bundles

Dolbeault cohomology: E holom. vect bundle \Rightarrow

$$C^\infty(X, E) \xrightarrow{\bar{\partial}} \Omega^{0,1}(X, E) \xrightarrow{\bar{\partial}} \Omega^{0,2}(X, E) \rightarrow \dots$$

$$\rightarrow H_{\bar{\partial}}^q(X, E) = \frac{\ker \bar{\partial}}{\text{Im } \bar{\partial}}$$

- Deforming J to a nearby J' :

$\Omega_{J'}^{1,0} \subseteq T^*X \otimes \mathbb{C} = \Omega_J^{1,0} \oplus \Omega_J^{0,1}$ is the graph of a linear map

$(-s): \Omega_J^{1,0} \rightarrow \Omega_J^{0,1}$. Conversely, given J' from s : if s small enough
 then $\Omega_{J'}^{1,0} := \text{graph}(-s), \Omega_{J'}^{0,1} = \overline{\Omega_J^{1,0}}$ satisfy $T^*X \otimes \mathbb{C} = \Omega_{J'}^{1,0} \oplus \Omega_{J'}^{0,1}$
 & set $J' = \begin{pmatrix} i & \\ & -i \end{pmatrix}$

Can also view s as section of $(\Omega_J^{1,0})^* \otimes \Omega_J^{0,1} \simeq T_J^{1,0} \otimes \Omega_J^{0,1}$
 ie. $(0,1)$ form with values in $T^{1,0}X$

z_1, \dots, z_n local holom. coordinates for $(X, J) \Rightarrow s = \sum s_{ij} \frac{\partial}{\partial z_i} \otimes d\bar{z}_j$

then basis of $(1,0)$ -forms for J' : $dz_i - s(d\bar{z}_i) = dz_i - \sum_j s_{ij} d\bar{z}_j$
 $(0,1)$ -vector fields $\frac{\partial}{\partial \bar{z}_k} + s \left(\frac{\partial}{\partial \bar{z}_k} \right) = \frac{\partial}{\partial \bar{z}_k} + \sum_l s_{lk} \frac{\partial}{\partial \bar{z}_l}$ } pair trivially

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Integrability?

$(\bigoplus_1 \Omega_x^{0,1} \otimes TX^{1,0}, \bar{\partial})$ Dolbeault complex for $TX^{1,0}$ on (X, J)
($\bar{\partial}$ acts on forms only)

carries a Lie bracket $[\alpha \otimes v, \alpha' \otimes v'] := (\alpha \wedge \alpha') \otimes [v, v']$

\rightarrow diff! graded Lie algebra (dgl).

Prop: J' is integrable $\Leftrightarrow \bar{\partial}s + \frac{1}{2}[s, s] = 0$

Pf: want: $\left[\frac{\partial}{\partial \bar{z}_i} + \sum_l s_{li} \frac{\partial}{\partial z_l}, \frac{\partial}{\partial \bar{z}_j} + \sum_l s_{lj} \frac{\partial}{\partial z_l} \right] \in TX_{J'}^{0,1}?$

$$= \sum_l \left(\frac{\partial s_{lj}}{\partial \bar{z}_i} \frac{\partial}{\partial z_l} - \frac{\partial s_{li}}{\partial \bar{z}_j} \frac{\partial}{\partial z_l} \right) + \sum_{k,l} \left(s_{ki} \frac{\partial s_{lj}}{\partial z_k} \frac{\partial}{\partial z_l} - s_{kj} \frac{\partial s_{li}}{\partial z_k} \frac{\partial}{\partial z_l} \right)$$

$\in \text{span} \left(\frac{\partial}{\partial z_l} \right) \dots$

\Rightarrow should be zero: want: $\forall i, j, l,$

$$\frac{\partial s_{lj}}{\partial \bar{z}_i} - \frac{\partial s_{li}}{\partial \bar{z}_j} + \sum_k \left(s_{ki} \frac{\partial s_{lj}}{\partial z_k} - s_{kj} \frac{\partial s_{li}}{\partial z_k} \right) = 0$$

coeff^t of $(d\bar{z}_i \wedge d\bar{z}_j) \otimes \frac{\partial}{\partial z_l}$
in $\bar{\partial}s$

$\frac{1}{2}$ -coeff of $d\bar{z}_i \wedge d\bar{z}_j \otimes \frac{\partial}{\partial z_l}$
in $[s, s]$

• We'd like to understand $\mathcal{M}_{cx}(X) = \{ J \text{ integrable cx. str. on } X \} / \text{Diff}(X)$
or rather its germ near X

(or: assuming $\text{Aut}(X, J)$ is discrete, near $J \exists$ universal family

$\mathcal{X} \xrightarrow{\pi} U \subset \mathcal{M}_{cx}$, \mathcal{X}, U complex manifolds, π holomorphic,
fibers of π are $\simeq X$

any other family near J is induced by a classifying map
& pullback from \mathcal{X} .

④ • $\{ \text{integrable } J's \} \underset{\text{locally}}{\cong} \{ s \in \Omega^{0,1}(X, TX^{1,0}) / \bar{\partial}s + \frac{1}{2}[s,s] = 0 \}$

but need to quotient by $\text{Diff}(X)$: $J \sim \phi^* J$

If ϕ is close to Id , $\phi^* J$ is close to $J \Rightarrow$ can be described as above.

$d\phi = \partial\phi + \bar{\partial}\phi$ where

$\partial\phi: TX^{1,0} \rightarrow \phi^* TX^{1,0}$ parts of $d\phi$ that commute / anticommute w/ J
 $\bar{\partial}\phi: TX^{0,1} \rightarrow \phi^* TX^{1,0}$

$\Rightarrow \phi^* dz_i = \underbrace{dz_i \circ \partial\phi}_{(1,0)} + \underbrace{dz_i \circ \bar{\partial}\phi}_{(0,1)} = \underbrace{(dz_i \circ \partial\phi)}_{(1,0)} \cdot (\text{Id} + (\partial\phi)^{-1} \bar{\partial}\phi)$

i.e. $s = -(\partial\phi)^{-1} \bar{\partial}\phi$

• Tangent space - infinitesimal deformations ("over $\text{Spec } \mathbb{C}[t]/t^2$ ")

$J(t), J(0) = J \rightarrow s(t) \in \Omega^{0,1}(X, TX^{1,0}), \bar{\partial}s(t) + \frac{1}{2}[s(t), s(t)] = 0$

$\rightarrow s_{\perp} = \frac{ds}{dt}|_{t=0}$ satisfies $\bar{\partial}s_{\perp} = 0$

Infinitesimal action of diffeomorphisms:

$(\phi_t), \phi_0 = \text{Id}, \frac{d\phi}{dt}|_{t=0} = v$ vector field \leadsto

$\frac{d}{dt}|_{t=0} \left(-(\partial\phi_t)^{-1} \bar{\partial}\phi_t \right) = -\frac{d}{dt}|_{t=0} (\bar{\partial}\phi_t) = -\bar{\partial}v$

So: $\left\| \begin{array}{l} \text{first order deformations} \\ \text{Def}_{\perp}(X, J) = \frac{\ker(\bar{\partial}: \Omega^{0,1}(X, TX^{1,0}) \rightarrow \Omega^{0,2})}{\text{Im}(\bar{\partial}: C^{\infty}(X, TX^{1,0}) \rightarrow \Omega^{0,1})} = H^1(X, TX^{1,0}) \end{array} \right.$

In particular, given a family $\begin{matrix} X \supset X \\ \downarrow \downarrow \\ S \ni 0 \end{matrix}$ of deformations of (X, J) parametrized by S

get a map $T_0 S \rightarrow H^1(X, TX)$ by looking at 1st order variations of J ...

Kodaira-Spencer map