

① Lagrangian Floer homology:

(M, ω) symplectic manifold $\supset L_0, L_1$ compact Lagrangian submanifolds

Formally, Floer homology = Morse theory for "action functional" on path space $\mathcal{P}(L_0, L_1)$, where crit pts are contract paths & gradient flowlines = J-hol. strips.

More precisely: $A: \widetilde{\mathcal{P}}(L_0, L_1) \rightarrow \mathbb{R}$, $(\gamma, [u]) \mapsto \int u^* \omega$
univ. cover
 $u: [0,1]^2 \rightarrow M$ homotopy $\ast \rightarrow \gamma$

$$dA(\gamma) \cdot v = \int_{[0,1]} \omega(\dot{\gamma}, v) dt = \int_{[0,1]} g(J\dot{\gamma}, v) dt = \langle J\dot{\gamma}, v \rangle_{L_2}$$

\uparrow vect-field along γ , $v(0) \in T_{\gamma(0)} L_0$, $v(1) \in T_{\gamma(1)} L_1$

hence crit pts = contr. paths $\dot{\gamma} = 0$; gradient traj. = J-hol. maps $\frac{\partial \gamma}{\partial s} = -J\dot{\gamma}$

Difficult to define rigorously ∞ -dim Morse theory, so use holom. curves instead.

Actual setup: Assume $L_0 \pitchfork L_1 \Rightarrow L_0 \cap L_1$ finite set.

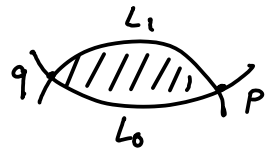
Recall Novikov ring $\Lambda = \{ \sum a_i T^{\lambda_i} \mid \lambda_i \rightarrow +\infty \}$

Floer complex $CF(L_0, L_1) = \Lambda^{|L_0 \cap L_1|}$ free Λ -module gen^d by $L_0 \cap L_1$.

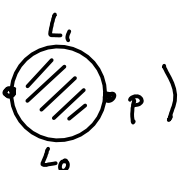
Goal: define a differential ∂ by counting holomorphic discs:

Look at: $u: \mathbb{R} \times [0,1] \rightarrow M$ equipped with J ω -compat. a.c.s.

$$s.t. \begin{cases} \bullet \bar{\partial}_J u = 0, \text{ i.e. } \frac{\partial u}{\partial s} + J \frac{\partial u}{\partial t} = 0 \\ \bullet u(s,0) \in L_0, u(s,1) \in L_1 \\ \bullet \lim_{s \rightarrow +\infty} u(s,t) = p, \lim_{s \rightarrow -\infty} u(s,t) = q \\ \bullet \text{the energy } E(u) = \int u^* \omega = \iint \left| \frac{\partial u}{\partial s} \right|^2 < \infty \end{cases}$$



(NB: $\mathbb{R} \times [0,1] \xrightarrow[\text{biholom.}]{\cong} \mathbb{D}^2 \setminus \{\pm 1\}$ so also think of maps $q \circ \text{disc} \rightarrow p$)



$\mathcal{M}(p, q, [u], J) = \{ u \text{ solns of } (\ast) \}$ moduli space of holom. discs.
 \uparrow
 $\pi_2(M; L_0, L_1)$ homotopy class

(\ast) is a Fredholm problem, exp. dim. $\mathcal{M} = \text{ind}(\bar{\partial}_J)$

$\text{ind}(\bar{\partial}_J) = \text{Maslov index}$

comes from $\pi_1(\Lambda Gr) = \mathbb{Z}$ for Lagrangian grassmannian in \mathbb{R}^{2n} .

② Maaslov index:

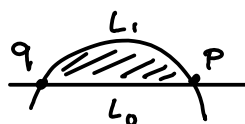
Let $L_0, L_1(t)_{t \in [0,1]}$ Lagr. subspaces of \mathbb{R}^{2n} , s.t. $L_1(0), L_1(1) \pitchfork L_0$

Then Maaslov index of the path $L_1(t) := \#$ times that $L_1(t)$ fails to be transverse to L_0 (counted with signs & multiplicities)

Ex: path $(e^{i\theta_1} \mathbb{R}) \times \dots \times (e^{i\theta_n} \mathbb{R})$ if $\theta_i \nearrow$ through 0, then $\mu(L_0, L_1(t)) = n$.

Now: given a strip u , initialize $u^* \text{TM} \rightarrow u^* \text{TL}_0, \text{TL}_1$ paths of Lagrangians. Can initialize so that TL_0 remain constant

Then $\text{ind}(u) :=$ Maaslov index of path TL_1 relative to TL_0 as one goes from p to q

Ex:  in \mathbb{R}^2 has $\text{ind}(u) = 1$. sympl. area of $\phi \downarrow (= \int u^* \omega)$

• Want to define: $\partial(p) = \sum_{\substack{q \in L_0 \cap L_1 \\ \phi \in \pi_2 / \text{ind}(\phi) = 1}} (\# \mathcal{M}(p, q, \phi, J) / \mathbb{R}) T^{u(\phi)} q$
 \uparrow
translation

- Issues:
- transversality
 - compactness, bubbling
 - orientation of \mathcal{M} (\Rightarrow signed counts? else work over $\mathbb{Z}/2$)
 - $\partial^2 = 0$?

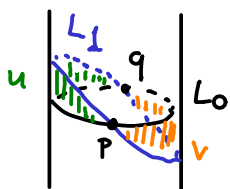
A special case where bubbling isn't an issue:

Thm (Floer): $\left\| \begin{array}{l} \text{if } [w] \cdot \pi_2(M) = 0 \text{ and } [w] \cdot \pi_2(M, L_i) = 0 \\ \text{then } \partial \text{ is well-def'd, } \partial^2 = 0, \text{ and HF is} \\ \text{indep't of chosen } J \text{ \& invariant under Hamiltonian deformations} \\ \text{of } L_0 \text{ and/or } L_1 \end{array} \right.$

Corollary: $\left\| [w] \cdot \pi_2(M, L) = 0, \psi \text{ Ham. diffeo, } \psi(L) \pitchfork L \Rightarrow |\psi(L) \cap L| \geq \sum b_i(L) \right.$

(special case of Arnold's conjecture: idea: $\text{HF}(L, \psi(L)) \cong H^*(L)$; $\text{rank CF} \geq \text{rank HF}$)

Example: $T^*S^1 = \mathbb{R} \times S^1$



$$\begin{aligned} \text{CF}(L_0, L_1) &= \Lambda_p \oplus \Lambda_q \\ \partial p &= (T^{\text{area}(u)} - T^{\text{area}(v)}) q \\ \partial q &= 0. \end{aligned}$$

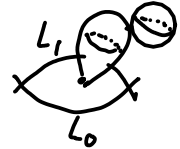
③ In this case \exists well-def'd \mathbb{Z} -grading (Maslov index only depends on p, q , not on homotopy class of strip), e.g. $\deg(p) = 0, \deg(q) = 1$.

- 2 cases: \ast $\text{area}(u) = \text{area}(v)$ (L_0, L_1 hom. isotopic): $\text{MF}(L_0, L_1) \cong H^*(S^1; \Lambda)$
- \ast $\text{area}(u) \neq \text{area}(v)$ (L_0, L_1 can be disjointed): $\text{MF}(L_0, L_1) = 0$.

Back to issues with the definition of ∂ :

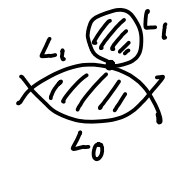
- transversality is achieved for simple maps by picking generic J .
For multiply covered maps (or configs with complicated bubbling), need various tricks... (domain-dependent J 's, multivalued perturbations, virtual cycles, ...)
- orientation on moduli space: need auxiliary data, & top assumption on L_i .
Namely: if equip L_i with spin structures (i.e. double cover of frame bundle) then we get an orientation on moduli spaces.
- compactness: again relies on Gromov's compactness theorem (for a fixed energy bound)
there are now 3 types of phenomena.

\rightarrow bubbling of spheres: $|du_n| \rightarrow \infty$ at interior point
limiting configuration looks like



Treatment similar to case of closed curves / GW invariants
In good cases (if transversality can be achieved), configs with sphere bubbles are a $\text{codim} \geq 2$ subset of the compactified $\overline{\mathcal{M}}$

\rightarrow bubbling of discs: $|du_n| \rightarrow \infty$ at boundary point
limiting configuration looks like



This is a serious technical issue because, in transverse case, bubbled configurations have real $\text{codim} \geq 1$, and contribute to $\partial \overline{\mathcal{M}}$.

\rightarrow breaking of strips: (or: energy escapes towards $s \rightarrow \pm\infty$ i.e. reparametrizing $u_n(\cdot - S_n, \cdot)$ have \neq limits)
(or, if using domain = $\mathbb{D}^2 - \{\pm 1\}$, this is bubbling at ± 1).

limiting configuration looks like



This is as in Morse theory, where a sequence of gradient flow lines can converge to a broken flow line



④ ★ How to prove $\partial^2 = 0$ assuming no bubbling:

consider $\mathcal{M}(p, q, \phi, \mathcal{J})/\mathbb{R}$ for $\phi \in \pi_2$ of index 2
 \mathcal{J} generic

This is expected to be a 1-dim^l mfd, which can be compactified by
adding in broken trajectories $\coprod_{\substack{r \in L_0 \cup L_1 \\ \phi_1, \# \phi_2 = \phi}} (\mathcal{M}(p, r, \phi_1, \mathcal{J})/\mathbb{R}) \cup (\mathcal{M}(r, q, \phi_2, \mathcal{J})/\mathbb{R})$

(this is assuming no bubbling \Rightarrow no other limiting scenarios!)

Gluing them \Rightarrow the resulting $\overline{\mathcal{M}(p, q, \phi, \mathcal{J})/\mathbb{R}}$ is a manifold w/ boundary.

Now: #ends of a compact 1-manifold, oriented (or counting mod 2), = 0.

Ends = contributions to coefficient of $T^{\omega(\phi)}_q$ in $\partial^2(p)$.

★ Bubbling of discs is not just a technical issue to overcome, it's an actual obstruction to defining Floer homology