

MIRROR SYMMETRY: LECTURE 2

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Reference for today: M. Gross, D. Huybrechts, D. Joyce, “Calabi-Yau Manifolds and Related Geometries”, Chapter 14.

1. DEFORMATIONS OF COMPLEX STRUCTURES

An (almost) complex structure (X, J) splits the complexified tangent and (wedge powers of) cotangent bundles as

$$\begin{aligned}
 TX \otimes \mathbb{C} &= TX^{1,0} \oplus TX^{0,1}, v^{0,1} = \frac{1}{2}(v + iJv) \\
 (1) \quad T^*X \otimes \mathbb{C} &= T^*X^{1,0} \oplus T^*X^{0,1}, T^*X^{1,0} = \text{Span}(dz_i), T^*X^{0,1} = \text{Span}(d\bar{z}_i) \\
 \bigwedge^k T^*X \otimes \mathbb{C} &= \bigoplus_{p+q=k} \bigwedge^{p,q} T^*X = \Omega^{p,q}(X)
 \end{aligned}$$

If J is almost complex, these are \mathbb{C} -vector bundles. J is integrable (i.e. a complex structure)

$$\begin{aligned}
 (2) \quad [T^{1,0}, T^{1,0}] \subset T^{1,0} &\Leftrightarrow d = \partial + \bar{\partial} \text{ maps } \Omega^{p,q} \rightarrow \Omega^{p+1,q} \oplus \Omega^{p,q+1} \\
 &\Leftrightarrow \bar{\partial}^2 = 0 \text{ on diff. forms}
 \end{aligned}$$

We obtain a Dolbeault cohomology for holomorphic vector bundles E :

$$\begin{aligned}
 (3) \quad C_{\bar{\partial}}^q(X, E) &= \{C^\infty(X, E) \xrightarrow{\bar{\partial}} \Omega^{0,1}(X, E) \xrightarrow{\bar{\partial}} \Omega^{0,2}(X, E) \rightarrow \dots\} \\
 H_{\bar{\partial}}^q(X, E) &= \ker \bar{\partial} / \text{im } \bar{\partial}
 \end{aligned}$$

Deforming J to a “nearby” J' gives

$$(4) \quad \Omega_{J'}^{1,0} \subseteq T^*\mathbb{C} = \Omega_J^{1,0} \oplus \Omega_J^{0,1}$$

is a graph of a linear map $(-s) : \Omega_J^{1,0} \rightarrow \Omega_J^{0,1}$. J' is determined by $\Omega_{J'}^{1,0}$ (acted on by i) and $\Omega_{J'}^{0,1}$ (acted on by i'). s is a section of $(\Omega_J^{1,0})^* \otimes \Omega_J^{0,1} = \mathbb{T}_J^{1,0} \otimes \Omega_J^{0,1}$ i.e. a $(0,1)_J$ -form with values in $T_J^{1,0}X$. If z_1, \dots, z_n are local holomorphic

coordinates for J , then $s = \sum s_{ij} \frac{\partial}{\partial z_i} \otimes d\bar{z}_j$. A basis of $(1, 0)$ -forms for J' is given by $dz_i - \underbrace{\sum_j s_{ij} d\bar{z}_j}_{s(dz_i)}$ and $(0, 1)$ -vectors for J' by $\frac{\partial}{\partial \bar{z}_k} + \underbrace{\sum_\ell s_{\ell k} \frac{\partial}{\partial z_\ell}}_{s(\partial/\partial \bar{z}_k)}$.

We can use this to test the integrability of J' . The Dolbeault complex $(\bigoplus_q \Omega_X^{0,q} \otimes TX^{1,0}, \bar{\partial})$ ($\bar{\partial}$ acts “on forms”) carries a Lie bracket

$$(5) \quad [\alpha \otimes v, \alpha' \otimes v'] = (\alpha \wedge \alpha') \otimes [v, v']$$

giving it the structure of a *differential graded Lie algebra*.

Proposition 1. J' is integrable $\Leftrightarrow \bar{\partial}s + \frac{1}{2}[s, s] = 0$.

Proof. We want to check that the bracket of two $0, 1$ tangent vectors is still $0, 1$, i.e. that

$$(6) \quad \left[\frac{\partial}{\partial \bar{z}_k} + \sum_\ell s_{\ell k} \frac{\partial}{\partial z_\ell}, \frac{\partial}{\partial \bar{z}_k} + \sum_\ell s_{\ell k} \frac{\partial}{\partial z_\ell} \right] \in TX_{J'}^{0,1}$$

Evaluating this bracket gives

$$(7) \quad \sum_\ell \left(\frac{\partial s_{\ell j}}{\partial \bar{z}_i} - \frac{\partial s_{\ell i}}{\partial \bar{z}_j} \right) \frac{\partial}{\partial z_\ell} + \sum_{k, \ell} \left(s_{ki} \frac{\partial s_{\ell j}}{\partial z_k} - s_{kj} \frac{\partial s_{\ell i}}{\partial z_k} \right) \frac{\partial}{\partial z_\ell}$$

We want this to be 0, i.e. for all i, j, ℓ ,

$$(8) \quad 0 = \underbrace{\frac{\partial s_{\ell j}}{\partial \bar{z}_i} - \frac{\partial s_{\ell i}}{\partial \bar{z}_j}}_{\text{coefficient of } \frac{\partial}{\partial z_\ell} \otimes (d\bar{z}_i \wedge d\bar{z}_j) \text{ in } (\bar{\partial}s)} + \sum_k \underbrace{\left(s_{ki} \frac{\partial s_{\ell j}}{\partial z_k} - s_{kj} \frac{\partial s_{\ell i}}{\partial z_k} \right)}_{\text{in } \frac{1}{2}[s, s]}$$

We leave the rest as an exercise. \square

We would now like to use this to understand the moduli space of complex structures. Define

$$(9) \quad \mathcal{M}_{CX}(X) = \{J \text{ integrable complex structures on } X\} / \text{Diff}(X)$$

(or, assuming that $\text{Aut}(X, J)$ is discrete, we want that near J , \exists a universal family $\mathcal{X} \rightarrow \mathcal{U} \subset \mathcal{M}_{CX}$ (complex manifolds, holomorphic fibers $\cong X$) s.t. any family of integrable complex structures $\mathcal{X}' \rightarrow S$ induces a map $S \rightarrow \mathcal{U}$ s.t. \mathcal{X} pulls back to \mathcal{X}'). We have an action of the diffeomorphisms of X : for $\phi \in \text{Diff}(X)$ close to id ,

$$(10) \quad \begin{aligned} d\phi &: TX \otimes \mathbb{C} \xrightarrow{\sim} \phi^* TX \otimes \mathbb{C} \\ \partial\phi &: TX^{1,0} \rightarrow \phi^* TX^{1,0} \\ \bar{\partial}\phi &: TX^{0,1} \rightarrow \phi^* TX^{1,0} \end{aligned}$$

so

$$(11) \quad \begin{aligned} \phi^* dz_i &= dz_i \circ d\phi = dz_i \circ \partial\phi + dz_i \circ \bar{\partial}\phi \\ &= \underbrace{(dz_i \circ \partial\phi)}_{(1,0) \text{ for } J} (\text{id} + (\partial\phi)^{-1} \cdot \bar{\partial}\phi) \end{aligned}$$

Deformation by $s \in \Omega^{0,1}(X, TX^{1,0})$ gives $\Omega_{J'}^{1,0} = \{\alpha - s(\alpha) | \alpha \in \Omega^{1,0}\}$ (the graph of $-s$): taking $s = -(\partial\phi)^{-1} \cdot \bar{\partial}\phi : TX^{0,1} \rightarrow \phi^* TX^{1,0} \rightarrow TX^{1,0}$ gives the desired element of $\Omega^{0,1}(TX^{1,0})$.

1.1. First-order infinitesimal deformations. Given a family $J(t)$, $J(0) = J$ gives $s(t) \in \Omega^{0,1}(X, TX^{1,0})$, $s(0) = 0$. By the above, this should satisfy

$$(12) \quad \bar{\partial}s(t) + \frac{1}{2}[s(t), s(t)] = 0$$

In particular, $s_1 = \frac{ds}{dt}|_{t=0}$ solves $\bar{\partial}s_1 = 0$. We obtain an infinitesimal action of $\text{Diff}(X)$: for (ϕ_t) , $\phi_0 = \text{id}$, $\frac{d\phi}{dt}|_{t=0} = v$ a vector field,

$$(13) \quad \frac{d}{dt}|_{t=0}(-(\partial\phi_t)^{-1} \circ \bar{\partial}\phi_t) = -\frac{d}{dt}|_{t=0}(\bar{\partial}\phi_t) = -\bar{\partial}v$$

This implies that first-order deformations are given as

$$(14) \quad \text{Def}_1(X, J) = \frac{\text{Ker } (\bar{\partial} : \Omega^{0,1}(TX^{1,0}) \rightarrow \Omega^{2,0}(TX^{1,0}))}{\text{Im}(\bar{\partial} : C^\infty(TX^{1,0}) \rightarrow \Omega^{0,1}(TX^{1,0}))}$$

We can write this more compactly using Dolbeault cohomology, namely $H_{\bar{\partial}}^1(X, TX^{1,0})$. Furthermore, given a family

$$(15) \quad \begin{array}{ccc} X & \longrightarrow & \mathcal{X} \\ \downarrow & & \downarrow \\ * & \longrightarrow & S \end{array}$$

of deformations of (X, J) parameterized by S , we get a map $T_*S \rightarrow H^1(X, TX^{1,0})$ called the *Kodaira Spencer map*

Remark. A complex manifold (X, J) is a union of complex charts U_i with biholomorphisms $\phi_{ij} : U_{ij} \xrightarrow{\sim} U_{ji}$ s.t. $\phi_{ij} = \phi_{ji}^{-1}$ and $\phi_{ij}\phi_{jk} = \phi_{ik}$ on U_{ijk} . Deformations of (X, J) come from deforming the gluing maps ϕ_{ij} among the space of holomorphic maps. To first order, this is given by holomorphic vector fields v_{ij} on $U_i \cap U_j$ s.t. $v_{ij} = -v_{ji}$ and $v_{ij} + v_{jk} = v_{ik}$ on U_{ijk} . This is precisely the Čech 1-cocycle conditions in the sheaf of holomorphic tangent vector fields. Modding out by holomorphic functions $\psi_i : U_i \xrightarrow{\sim} U_i$ (which act by $\phi_{ij} \mapsto \psi_j \phi_{ij} \psi_i^{-1}$) is precisely modding by the Čech coboundaries. Thus, $\text{Def}_1(X, J) = \check{H}^1(X, TX^{1,0})$.

1.2. Obstructions to Deformation. Given a first-order deformation s_1 , one can ask if one can find an actual deformation $s(t) = s_1 t + O(t^2)$ (or even a formal deformation, i.e. non-convergent power series). Expand

$$(16) \quad s(t) = s_1 t + s_2 t^2 + \cdots \in \Omega^{0,1}(X, TX^{1,0})$$

Then the condition $\bar{\partial}s(t) + \frac{1}{2}[s(t), s(t)] = 0$ implies that $\bar{\partial}s_1 = 0, \bar{\partial}s_2 + \frac{1}{2}[s_1, s_1] = 0, \bar{\partial}s_3 + [s_1, s_2] = 0, \dots$. Now, we need $[s_1, s_1] \in \text{im}(\bar{\partial}) \subset \Omega^{0,2}(TX^{1,0})$. We know that $[s_1, s_1] \in \text{Ker}(\bar{\partial})$. Thus, the primary obstruction to deforming is the class of $[s_1, s_1]$ in $H^2(X, TX^{1,0})$. If it is zero, then there is an s_2 s.t. $\bar{\partial}s_2 + \frac{1}{2}[s_1, s_1] = 0$, and the next obstruction is the class of $[s_1, s_2] \in H^2(X, TX^{1,0})$. We are basically attempting to apply by brute force the implicit function theorem.

If it happens that $H^2(X, TX) = 0$, then the deformations are unobstructed and the moduli space of complex structures is locally a smooth orbifold (not a manifold, because we may have to quotient by automorphisms) with tangent space $H^1(X, TX^{1,0})$. For Calabi-Yau manifolds, this will not be true: however, we still have

Theorem 1 (Bogomolov-Tian-Todorov). *For X a compact Calabi-Yau ($\Omega_X^{n,0} \cong \mathcal{O}_X$) with $H^0(X, TX) = 0$ (automorphisms are discrete), deformations of X are unobstructed and, assuming $\text{Aut}(X, J) = \{1\}$, \mathcal{M}_{CX} is locally a smooth manifold with $T\mathcal{M}_{CX} = H^1(X, TX)$.*

Theorem 2 (Griffiths Transversality). *For a family (X, J_t) , $\alpha_t \in \Omega^{p,q}(X, J_t) \implies \frac{d}{dt}|_{t=0} \alpha_t \in \Omega^{p,q} + \Omega^{p+1,q-1} + \Omega^{p-1,q+1}$.*

Proof. J_t is given by $s(t) \in \Omega^{0,1}(TX^{1,0}), s(0) = 0$. In local coordinates, we have $T^*X_{J_t}^{1,0} = \text{Span}\{dz_i^{(t)} = dz_i - \sum s_{ij}(t)d\bar{z}_j\}$

$$(17) \quad \alpha_t = \sum_{I, J || I|=p, |J|=q} \alpha_{IJ}(t) dz_{i_1}^{(t)} \wedge \cdots \wedge dz_{i_p}^{(t)} \wedge d\bar{z}_{j_1}^{(t)} \wedge \cdots \wedge d\bar{z}_{j_q}^{(t)}$$

Taking $\frac{d}{dt}|_{t=0}$, the result follows from the product rule. We mostly get (p, q) terms and a few $(p+1, q-1), (p-1, q+1)$ forms (the latter from $\frac{d}{dt}|_{t=0}(dz_{i_k}^{(t)})$). \square