

:

MSRI, subfactors, knots and the development of quantum topology.

Vaughan Jones

25th Anniversary celebration of MSRI, Jan 2008



MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

Before the MSRI year:

MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

Before the MSRI year: Rediscovery in subfactors of the "Temperley-Lieb" Algebra.

$$e_i^2 = e_i^* = e_i, \quad i = 1, 2, \dots, n$$

(n orthogonal projections onto subspaces of Hilbert Space.)

MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

Before the MSRI year: Rediscovery in subfactors of the "Temperley-Lieb" Algebra.

$$e_i^2 = e_i^* = e_i, \quad i = 1, 2, \dots, n$$

(n orthogonal projections onto subspaces of Hilbert Space.)

$$e_i e_j = e_j e_i \quad \text{if } |i - j| \geq 2$$

(subspaces orthogonal modulo their intersection if $|i - j| \geq 2$)

MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

Before the MSRI year: Rediscovery in subfactors of the "Temperley-Lieb" Algebra.

$$e_i^2 = e_i^* = e_i, \quad i = 1, 2, \dots, n$$

(n orthogonal projections onto subspaces of Hilbert Space.)

$$e_i e_j = e_j e_i \quad \text{if } |i - j| \geq 2$$

(subspaces orthogonal modulo their intersection if $|i - j| \geq 2$)

$$e_i e_{i\pm 1} e_i = \tau e_i \quad \text{for } 1 \leq i < n$$

("angle" between i th. and $(i + 1)$ th. subspaces determined by the number $\tau \in \mathbb{R}$.)

MSRI 1984-1985

Operator algebras year side by side with Low dimensional topology year.

Many themes, *subfactors* $N \subseteq M$ being one of them.

Before the MSRI year: Rediscovery in subfactors of the "Temperley-Lieb" Algebra.

$$e_i^2 = e_i^* = e_i, \quad i = 1, 2, \dots, n$$

(n orthogonal projections onto subspaces of Hilbert Space.)

$$e_i e_j = e_j e_i \quad \text{if } |i - j| \geq 2$$

(subspaces orthogonal modulo their intersection if $|i - j| \geq 2$)

$$e_i e_{i\pm 1} e_i = \tau e_i \quad \text{for } 1 \leq i < n$$

("angle" between i th. and $(i + 1)$ th. subspaces determined by the number $\tau \in \mathbb{R}$.)

And a trace tr on the algebra generated by the e_i 's uniquely defined by

$$tr(w e_{n+1}) = \tau tr(w) \quad \text{if } w \text{ is a word on } e_1, e_2, \dots, e_n$$

Restrictions on τ

For infinitely many e_i 's to exist as above, $\frac{1}{\tau}$ must be either ≥ 4 or one of the numbers

$$4 \cos^2 \pi/k \quad \text{for } k = 3, 4, 5, \dots$$

Restrictions on τ

For infinitely many e_i 's to exist as above, $\frac{1}{\tau}$ must be either ≥ 4 or one of the numbers

$$4 \cos^2 \pi/k \quad \text{for } k = 3, 4, 5, \dots$$

Wenzl- these restrictions apply even without the trace.

Restrictions on τ

For infinitely many e_i 's to exist as above, $\frac{1}{\tau}$ must be either ≥ 4 or one of the numbers

$$4 \cos^2 \pi/k \quad \text{for } k = 3, 4, 5, \dots$$

Wenzl- these restrictions apply even without the trace.

Implications for subfactors:

- 1) index $[M : N]$ must be ≥ 4 or in the above set of numbers.
- 2) Subfactors exist with all these index values (constructed from the e_i 's and the trace).

Restrictions on τ

For infinitely many e_i 's to exist as above, $\frac{1}{\tau}$ must be either ≥ 4 or one of the numbers

$$4 \cos^2 \pi/k \quad \text{for } k = 3, 4, 5, \dots$$

Wenzl- these restrictions apply even without the trace.

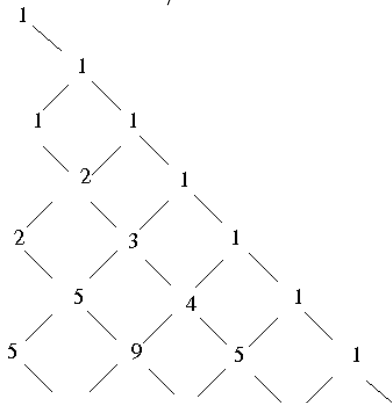
Implications for subfactors:

- 1) index $[M : N]$ must be ≥ 4 or in the above set of numbers.
- 2) Subfactors exist with all these index values (constructed from the e_i 's and the trace).

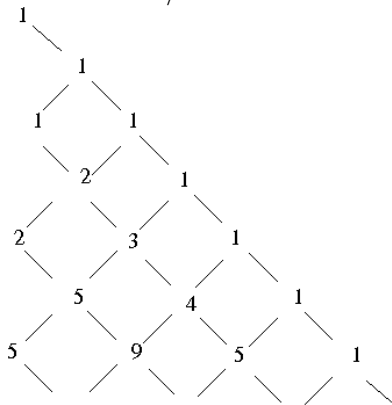
Representation theory of the e_i 's.

Given by the following diagram:

For $0 < \tau < 1/4$:



For $0 < \tau < 1/4$:



or for $1/\tau = 4 \cos^2 \pi/k$:

Braid group: on generators $\sigma_1, \sigma_2, \dots, \sigma_{n-1}$,

$$\sigma_i \sigma_j = \sigma_j \sigma_i \text{ if } |i - j| \geq 2$$

$$\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1} \text{ for } 1 \leq i < n$$

Braid group: on generators $\sigma_1, \sigma_2, \dots, \sigma_{n-1}$,

$$\sigma_i \sigma_j = \sigma_j \sigma_i \text{ if } |i - j| \geq 2$$

$$\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1} \text{ for } 1 \leq i < n$$

Compare:

$$e_i e_j = e_j e_i \text{ if } |i - j| \geq 2$$

$$e_i e_{i\pm 1} e_i = \tau e_i \text{ for } 1 \leq i < n$$

Braid group: on generators $\sigma_1, \sigma_2, \dots, \sigma_{n-1}$,

$$\sigma_i \sigma_j = \sigma_j \sigma_i \text{ if } |i - j| \geq 2$$

$$\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1} \text{ for } 1 \leq i < n$$

Compare:

$$e_i e_j = e_j e_i \text{ if } |i - j| \geq 2$$

$$e_i e_{i\pm 1} e_i = \tau e_i \text{ for } 1 \leq i < n$$

To get a representation of the braid group send σ_i to

$$te_i - (1 - e_i)$$

With $\tau = \frac{t}{(1+t)^2}$.

Braid group: on generators $\sigma_1, \sigma_2, \dots, \sigma_{n-1}$,

$$\sigma_i \sigma_j = \sigma_j \sigma_i \text{ if } |i - j| \geq 2$$

$$\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1} \text{ for } 1 \leq i < n$$

Compare:

$$e_i e_j = e_j e_i \text{ if } |i - j| \geq 2$$

$$e_i e_{i\pm 1} e_i = \tau e_i \text{ for } 1 \leq i < n$$

To get a representation of the braid group send σ_i to

$$te_i - (1 - e_i)$$

With $\tau = \frac{t}{(1+t)^2}$. If e_i were diagonalised, σ_i look like:

$$\begin{pmatrix} t & 0 & 0 & & & \\ 0 & t & 0 & \dots & & \\ 0 & 0 & t & & & \\ \dots & & & & -1 & 0 \\ & & & & 0 & -1 \end{pmatrix}$$

Note: $0 < \tau \leq 1/4$ means $0 < t < \infty$ while

$$1/\tau = 4 \cos^2 \pi/k \text{ means } t = e^{\frac{2\pi i}{k}}$$

$$\begin{pmatrix} t & 0 & 0 & & & \\ 0 & t & 0 & \dots & & \\ 0 & 0 & t & & & \\ \dots & & & & -1 & 0 \\ & & & & 0 & -1 \end{pmatrix}$$

Note: $0 < \tau \leq 1/4$ means $0 < t < \infty$ while

$$1/\tau = 4 \cos^2 \pi/k \text{ means } t = e^{\frac{2\pi i}{k}}$$

So these braid group representations are *manifestly unitary* if $t = e^{\frac{2\pi i}{k}}$ and not otherwise.

$$\begin{pmatrix} t & 0 & 0 & & & \\ 0 & t & 0 & \dots & & \\ 0 & 0 & t & & & \\ \dots & & & & -1 & 0 \\ & & & & 0 & -1 \end{pmatrix}$$

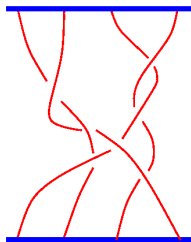
Note: $0 < \tau \leq 1/4$ means $0 < t < \infty$ while

$$1/\tau = 4 \cos^2 \pi/k \text{ means } t = e^{\frac{2\pi i}{k}}$$

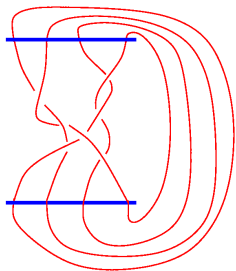
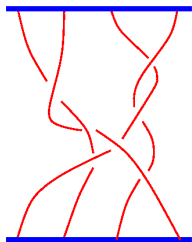
So these braid group representations are *manifestly unitary* if $t = e^{\frac{2\pi i}{k}}$ and not otherwise.

Braids to knots:

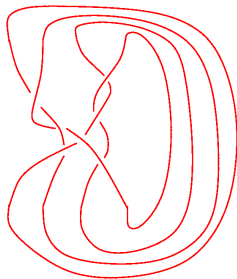
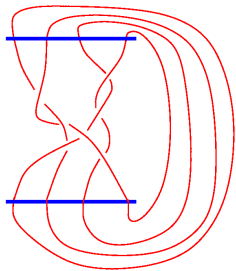
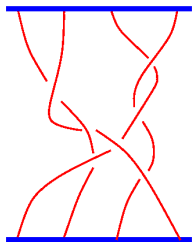
Braids to knots:



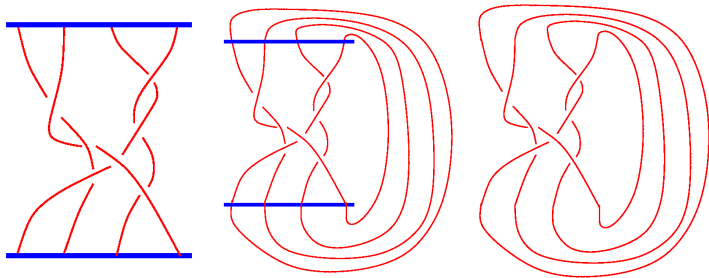
Braids to knots:



Braids to knots:

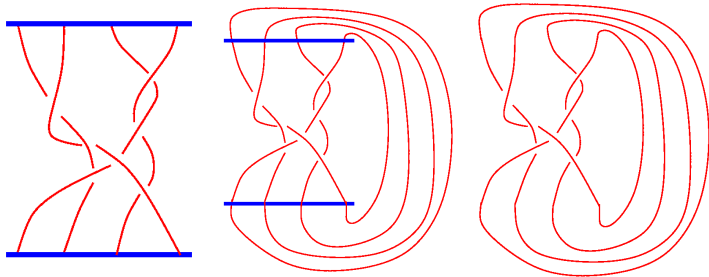


Braids to knots:



$$\alpha \in B_n \rightarrow \bar{\alpha}$$

Braids to knots:



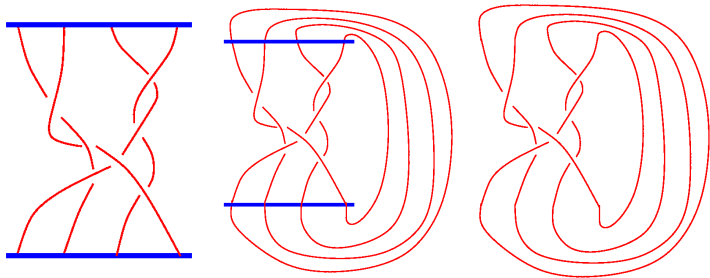
$$\alpha \in B_n \rightarrow \bar{\alpha}$$

Observation: If α is represented in the Temperley-Lieb algebra as above, by a theorem of Markov,

(normalisation) $tr(\alpha)$

gives an invariant of the knot (or link) \overline{alpha} .

Braids to knots:



$$\alpha \in B_n \rightarrow \bar{\alpha}$$

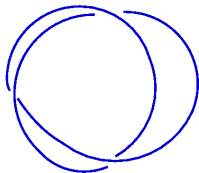
Observation: If α is represented in the Temperley-Lieb algebra as above, by a theorem of Markov,

(normalisation) $tr(\alpha)$

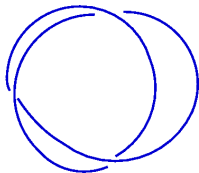
gives an invariant of the knot (or link) $\overline{\alpha}$. It's a Laurent polynomial in t , written $V_L(t)$ for a link L .

For the trefoil knot $K =$

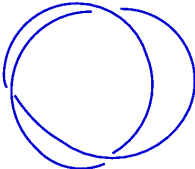
For the trefoil knot $K =$



For the trefoil knot $K =$



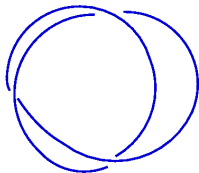
$$V_K(t) = t + t^3 - t^4$$

For the trefoil knot $K =$ 

$$V_K(t) = t + t^3 - t^4$$

This summarises what was known going into the year 1984-1985.

For the trefoil knot $K =$



$$V_K(t) = t + t^3 - t^4$$

This summarises what was known going into the year 1984-1985.
Then things moved FAST.

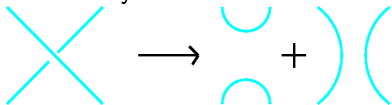
- 1) The polynomial $V_L(t)$ was extended to a two variable generalisation , the HOMFLYPT polynomial, which also contained the Alexander polynomial.
- 2) Several applications including braid index estimates and answers to some questions about Conway's "skein theory".
- 3) Observation from stat mechanical models that $V_L(t)$ is essentially independent of the orientation of L . Followed within a few weeks by Kauffman's great diagrammatic insights:
 - a) $V_L(t)$ is simply the result of replacing a crossing by no crossing in both ways:

1) The polynomial $V_L(t)$ was extended to a two variable generalisation, the HOMFLYPT polynomial, which also contained the Alexander polynomial.

2) Several applications including braid index estimates and answers to some questions about Conway's "skein theory".

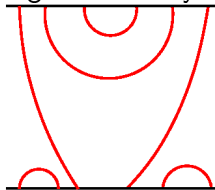
3) Observation from stat mechanical models that $V_L(t)$ is essentially independent of the orientation of L . Followed within a few weeks by Kauffman's great diagrammatic insights:

a) $V_L(t)$ is simply the result of replacing a crossing by no crossing in both ways:

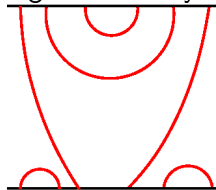


b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:

b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:

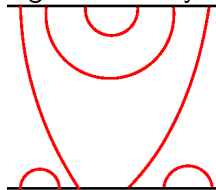


b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:



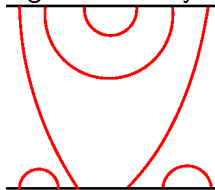
With braid-like concatenation as multiplication.

b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:

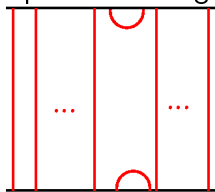


With braid-like concatenation as multiplication.
In particular the generator e_j is represented by:

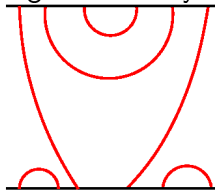
b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:



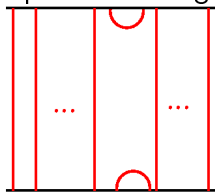
With braid-like concatenation as multiplication. In particular the generator e_j is represented by:



b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:

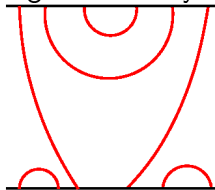


With braid-like concatenation as multiplication. In particular the generator e_j is represented by:

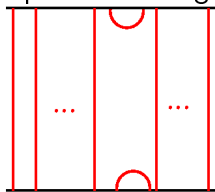


And in the same breath Kauffman discovered another 2-variable knot polynomial.

b) The Temperley-Lieb algebra can be realised entirely diagrammatically as pictures like:



With braid-like concatenation as multiplication. In particular the generator e_j is represented by:



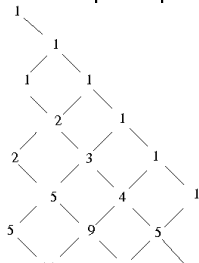
And in the same breath Kauffman discovered another 2-variable knot polynomial. His diagrammatics soon led to the solution of a Tait conjecture about alternating knots. The Tait conjectures were completely solved more recently by Menasco and Thistlethwaite.

On the subfactor side progress was made by Ocneanu and Wenzl. Ocneanu's approach to the HOMFLYPT polynomial saw it as coming from the Hecke algebra and isolated a sequence of 1-variable polynomials understood to have something to do with $SU(n)$, the polynomial $V_L(t)$ being the case $n = 2$. Used the corresponding Hecke algebras to construct new subfactors.

On the subfactor side progress was made by Ocneanu and Wenzl. Ocneanu's approach to the HOMFLYPT polynomial saw it as coming from the Hecke algebra and isolated a sequence of 1-variable polynomials understood to have something to do with $SU(n)$, the polynomial $V_L(t)$ being the case $n = 2$. Used the corresponding Hecke algebras to construct new subfactors. The *principal graph* of a subfactor emerged, being the graph A_n for the subfactors constructed from the Temperley-Lieb algebra. It encodes induction-restriction of bimodules between N and M . Ocneanu gave a complete classification of subfactors in index < 4 by Coxeter graphs of types A_n , D_{2n} , E_6 and E_8 .

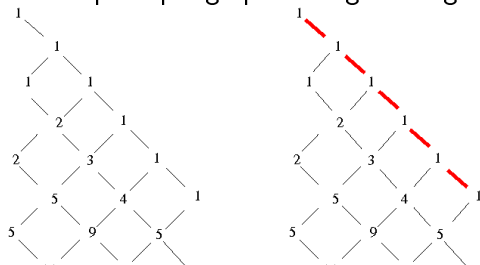
On the subfactor side progress was made by Ocneanu and Wenzl. Ocneanu's approach to the HOMFLYPT polynomial saw it as coming from the Hecke algebra and isolated a sequence of 1-variable polynomials understood to have something to do with $SU(n)$, the polynomial $V_L(t)$ being the case $n = 2$. Used the corresponding Hecke algebras to construct new subfactors. The *principal graph* of a subfactor emerged, being the graph A_n for the subfactors constructed from the Temperley-Lieb algebra. It encodes induction-restriction of bimodules between N and M . Ocneanu gave a complete classification of subfactors in index < 4 by Coxeter graphs of types A_n , D_{2n} , E_6 and E_8 .

To see the principal graph arising look again at :



On the subfactor side progress was made by Ocneanu and Wenzl. Ocneanu's approach to the HOMFLYPT polynomial saw it as coming from the Hecke algebra and isolated a sequence of 1-variable polynomials understood to have something to do with $SU(n)$, the polynomial $V_L(t)$ being the case $n = 2$. Used the corresponding Hecke algebras to construct new subfactors. The *principal graph* of a subfactor emerged, being the graph A_n for the subfactors constructed from the Temperley-Lieb algebra. It encodes induction-restriction of bimodules between N and M . Ocneanu gave a complete classification of subfactors in index < 4 by Coxeter graphs of types A_n , D_{2n} , E_6 and E_8 .

To see the principal graph arising look again at :



At the end of the MSRI year 1984-1985 there was plenty of action in this area but the following aspects were unsatisfactory:

i) A "context" was needed for subfactors.

At the end of the MSRI year 1984-1985 there was plenty of action in this area but the following aspects were unsatisfactory:

- i) A "context" was needed for subfactors.

- ii) A topological interpretation was needed for $V_L(t)$ and the other polynomials.

1985 – 1988

Organisational period. Jimbo, Woronowicz, Drinfeld - quantum group.

Picture emerged:

1985 – 1988

Organisational period. Jimbo, Woronowicz, Drinfeld - quantum group.

Picture emerged:

To every irreducible finite dimensional representation of a simple lie group/algebra there is

- 1) A one parameter family of subfactors. (Sawin)
- 2) A knot polynomial with direct statistical mechanical sum formula. For links can put representations on individual components.
- 3) A braid group representation with Markov trace.

1985 – 1988

Organisational period. Jimbo, Woronowicz, Drinfeld - quantum group.

Picture emerged:

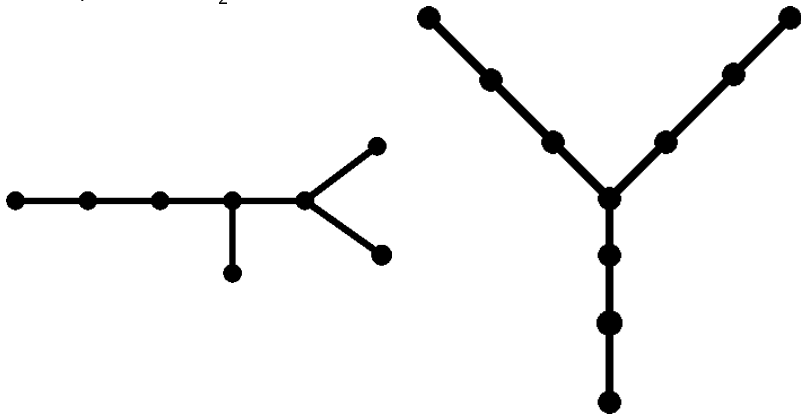
To every irreducible finite dimensional representation of a simple lie group/algebra there is

- 1) A one parameter family of subfactors. (Sawin)
- 2) A knot polynomial with direct statistical mechanical sum formula. For links can put representations on individual components.
- 3) A braid group representation with Markov trace.

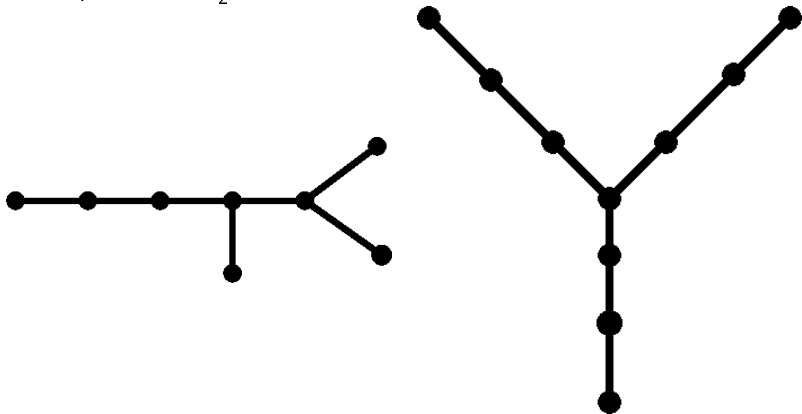
Roots of unity:

Tricky in quantum group world but here the subfactor picture reigns supreme (Wenzl, Feng Xu) as all the non semi-simple clutter is pared away by positivity.

To jump ahead into the '90s, it was in the subfactor arena that enduring "sporadic" objects were discovered by Haagerup (and Asaeda). Index $\frac{5+\sqrt{13}}{2}$ and principal graphs :



To jump ahead into the '90s, it was in the subfactor arena that enduring "sporadic" objects were discovered by Haagerup (and Asaeda). Index $\frac{5+\sqrt{13}}{2}$ and principal graphs :



These objects remain "exotic" creatures in the zoo.

The 600 pound gorilla: CONFORMAL FIELD THEORY

Belavin Polyakov Zamolodchikov.

A paper by Tsuchiya and Kanie obtaining the braid group reps (at roots of unity) in a WZW model.

The 600 pound gorilla: CONFORMAL FIELD THEORY

Belavin Polyakov Zamolodchikov.

A paper by Tsuchiya and Kanie obtaining the braid group reps (at roots of unity) in a WZW model.

Even more murky, Fredenhagen, Rehren, Schroer, and Longo, suggesting that braid groups and subfactors should crop up in general low dimensional QFT in the Haag-Kastler algebraic QFT framework following Dopplcher, Haag, Roberts.

1988: Witten- TQFT (Atiyah)

An interpretation of $V_L(t)$, at $t = e^{2\pi i/k}$, as a functional integral in a 2+1 dimensional gauge theory, gauge group $SU(2)$ (as expected), with Chern -Simons action, the knot invariant being the expected value of the trace of the monodromy along the components of the link which now became known as "Wilson loops".

1988: Witten- TQFT (Atiyah)

An interpretation of $V_L(t)$, at $t = e^{2\pi i/k}$, as a functional integral in a 2+1 dimensional gauge theory, gauge group $SU(2)$ (as expected), with Chern -Simons action, the knot invariant being the expected value of the trace of the monodromy along the components of the link which now became known as "Wilson loops".

EXPLICIT formula for links in an arbitrary 3-manifold via surgery on links or Heegard splitting!

Complete verificaiton by Reshetikhin-Turaev.

Turaev-Viro Ocneanu-3 manifold invariants from general categorical data as provided for instance by a subfactor.

MSRI meeting January 1989. Lots of key actors present, Atiyah, Bott, Witten, Jimbo, Miwa, Kauffman, Faddeev, Wassermann.... Greatly clarified the situation. Finally got over the feeling that the Russians and the Japanese knew something that we didn't.

MSRI meeting January 1989. Lots of key actors present, Atiyah, Bott, Witten, Jimbo, Miwa, Kauffman, Faddeev, Wassermann.... Greatly clarified the situation. Finally got over the feeling that the Russians and the Japanese knew something that we didn't. Subfactor outcome: Wassermann's construction of subfactors from loop group representations: take positive energy projective representations of $LSU(2)$, parametrized by a central extension (level) and an irrep of $SU(2)$, then if I is an interval in the circle with I^c the complementary interval,

$$L_I SU(2)'' \subseteq L_{I^c} SU(2)'$$

is a subfactor realising the indices $4 \cos^2 \pi / (\text{level})$ (more generally $\frac{\sin^2(k\pi/\ell)}{\sin^2\pi/\ell}$).

And in fact this is an instance of the Fredenhagen-Rehren-Schroer theory, with braiding as predicted!

MSRI meeting January 1989. Lots of key actors present, Atiyah, Bott, Witten, Jimbo, Miwa, Kauffman, Faddeev, Wassermann.... Greatly clarified the situation. Finally got over the feeling that the Russians and the Japanese knew something that we didn't. Subfactor outcome: Wassermann's construction of subfactors from loop group representations:

take positive energy projective representations of $LSU(2)$, parametrized by a central extension (level) and an irrep of $SU(2)$, then if I is an interval in the circle with I^c the complementary interval,

$$L_I SU(2)'' \subseteq L_{I^c} SU(2)'$$

is a subfactor realising the indices $4 \cos^2 \pi / (\text{level})$ (more generally $\frac{\sin^2(k\pi/\ell)}{\sin^2\pi/\ell}$).

And in fact this is an instance of the Fredenhagen-Rehren-Schroer theory, with braiding as predicted!

The main technical and conceptual tool used by Wassermann is the Connes tensor product of bimodules over a von Neumann algebra, which in turn relies on the Tomita-Takesaki theory.

Turning again to topology, a new ingredient entered the game in the early 1990's - Vassiliev theory.

Thinking of knots as smooth functions from S^1 to \mathbb{R}^3 , Vassiliev began the study of the algebraic topology of this space and obtained new knot invariants. Birman and Lin furthered their study and the concept of "finite type invariant" emerged. The idea is to extend any additive knot invariant Inv to an invariant of immersed curves with at worst double point singularities by locally setting

$$Inv \begin{array}{c} \nearrow \\ \times \\ \searrow \end{array} = Inv \begin{array}{c} \nearrow \\ \times \\ \searrow \end{array} - Inv \begin{array}{c} \nearrow \\ \times \\ \searrow \end{array}$$

Then Inv is of finite type n if it vanishes on immersed curves with n double points.

The knot polynomials are a source of finite type invariants but it was shown by Vogel that there are others. Kontsevich discovered a beautiful general integral formula for finite type invariants, thinking of them more as a perturbative expansion of a functional integral.

Bar-Natan put everything together in a stunningly comprehensible paper which made the whole theory broadly accessible and showed that one could construct knot invariants by studying "chord diagrams" modulo certain linear relations.

There emerged a basic relationship with (not necessarily simple) Lie algebras and understanding of how these invariants fit together, even for the simplest knots, gives information about Lie algebras.

(Bar-Natan, Le, Thurston-Alexeev.) Vogel showed that not all finite type invariants come from Lie algebras.

All this has to a certain extent been extended to knots and links in arbitrary 3-manifolds though I believe that, beyond homology 3-spheres, things remain a little murky.

Another major problem in this area is to understand how the Witten-Reshetikhin-Turaev invariant depends on the root of unity. In the best of all worlds there would be a holomorphic function of q of which it is the boundary value but this appears to be too naive. Recent progress by Garoufalidis.

There are various conjectures concerning the asymptotic behaviour of the WRT invariants and how they relate to other geometric or topological invariants. Let me mention Kashaev's volume conjecture (just for knots). If K is a knot with hyperbolic complement, one looks at

$$V_{SU_2, k\text{dimensionalrepresentation}}(e^{\pi i/k})$$

(the value of t is the value for which the invariant of an unknot is zero (!))

There are various conjectures concerning the asymptotic behaviour of the WRT invariants and how they relate to other geometric or topological invariants. Let me mention Kashaev's volume conjecture (just for knots). If K is a knot with hyperbolic complement, one looks at

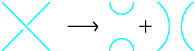
$$V_{SU_2, k\text{dimensional representation}}(e^{\pi i/k})$$

(the value of t is the value for which the invariant of an unknot is zero (!))

Then as $k \rightarrow \infty$, the asymptotic growth of this number is controlled by the hyperbolic volume.

Result is known for some small knots by explicit computations. The experts are not entirely in agreement as to whether it should be true.

Other asymptotic expansions proved recently by J. Anderson. Marcos Marino, Aganagic, Klemm, Vafa and others made connections with Calabi-Yau three-folds via matrix models.

A more recent development has been Khovanov homology. This uses the basic Kauffman resolution  of a crossing but "categorifies" it in that the basic invariant associated to a link is a graded complex and the relation between the complexes at a crossing is an exact sequence between the three complexes. The invariant ends up, at least in the hands of Bar-Natan, being a complex up to homotopy equivalence. The Euler characteristic of the complex is $V_L(t)$. Bar-Natan's version is highly calculable by a "planar algebra" approach.

Khovanov and others have extended categorification to many other knot invariants and beyond.

Khovanov homology has yielded proofs of topological results (Rasmussen), it is strictly more powerful than the corresponding Euler Characteristic polynomials, and it looks like it is the arena for connections with gauge theory, e.g. Floer homology.

What about subfactors?

The 90's saw deep classification results by Sorin Popa who showed that in "amenable" cases (both the factor, the subfactor and the principal graph), the subfactor is completely classified by what is known as the "standard invariant" (an enrichment of the principal graphs containing not just the bimodules and their tensor powers but intertwiners between them). These results extend Connes' breakthrough results from the 70's classifying automorphisms and group actions.

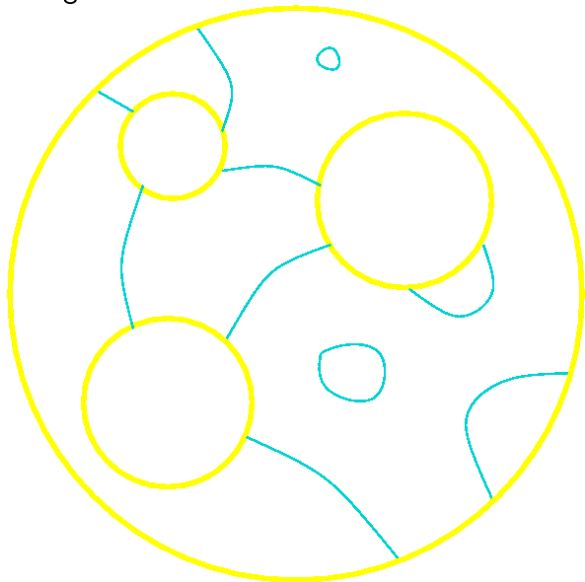
An Einstein manifold is one for which the Ricci tensor is a constant multiple of the metric. The constant is called the cosmological constant and Einstein stated at one point that it was the greatest mistake in his career....

An Einstein manifold is one for which the Ricci tensor is a constant multiple of the metric. The constant is called the cosmological constant and Einstein stated at one point that it was the greatest mistake in his career....

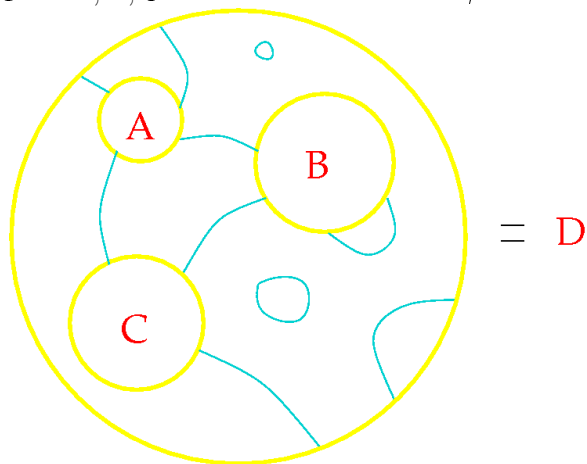
An attempt to use subfactors to provide information on a huge variety of combinatorial problems. Difficulty: Compute the principal graph.

May eventually prove fruitful though no doubt less universal than hoped, but one thing it did lead to was the development of *planar algebras* which provided a very convenient axiomatisation of the standard invariant of a subfactor and a useful technique for analysing their structure. Similar in spirit to Conway's "skein theory".

A tangle:



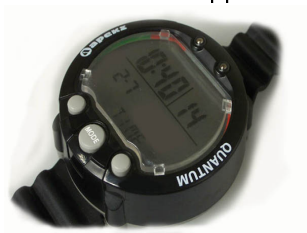
Given A, B, C in the standard invariant,



Gives D also in the standard invariant. Recent interaction with random matrices and free probability (Voiculescu). (with Guionnet and Shlyakhtenko.)

Freedman et al approach to quantum computer.

Freedman et al approach to quantum computer.



Freedman et al approach to quantum computer.



Freedman et al approach to quantum computer.



The idea is to spread out the quantum system in some 2-dimensional quantum fluid.

The qubits are then replaced by quasiparticles. According to FRS, if one carries out a motion in such a fluid, returning n points to their original position, the Hilbert space should evolve according to a braid group representation. The ones described at the beginning of this talk are precisely those one would expect for such a system. One could then encode a calculation as a braid.

The idea is to spread out the quantum system in some 2-dimensional quantum fluid.

The qubits are then replaced by quasiparticles. According to FRS, if one carries out a motion in such a fluid, returning n points to their original position, the Hilbert space should evolve according to a braid group representation. The ones described at the beginning of this talk are precisely those one would expect for such a system.

One could then encode a calculation as a braid.

Freedman et al have shown that for the value $t = e^{2\pi i/5}$ such a computer would be "Universal for quantum computation". The problem then is to find a physical system obeying such non-abelian statistics. Hopes are pinned on the fractional quantum hall effect at very low temperature.

The idea is to spread out the quantum system in some 2-dimensional quantum fluid.

The qubits are then replaced by quasiparticles. According to FRS, if one carries out a motion in such a fluid, returning n points to their original position, the Hilbert space should evolve according to a braid group representation. The ones described at the beginning of this talk are precisely those one would expect for such a system.

One could then encode a calculation as a braid.

Freedman et al have shown that for the value $t = e^{2\pi i/5}$ such a computer would be "Universal for quantum computation". The problem then is to find a physical system obeying such non-abelian statistics. Hopes are pinned on the fractional quantum hall effect at very low temperature.

Freedman's Navier Stokes on Mars analogy.