

Even worse, $V_\varepsilon: \mathbb{C}P^1 \rightarrow \mathbb{C}P^2, [x:y] \mapsto [\varepsilon x^2: \sqrt{\varepsilon} y^2: xy]$ parametrizes C_ε

but converges to $V_0[x:y] = [0:0:1]$ except for $|\mathrm{d}V_\varepsilon|$ -blowup at $[1:0]$ and $[0:1]$.

One blowup can be avoided by reparametrization $V_\varepsilon \circ \varphi_\varepsilon = u_\varepsilon$ ($\varphi_\varepsilon[x:y] = [\sqrt{\varepsilon}x: y/\sqrt{\varepsilon}]$)

(or $V_\varepsilon \circ (\varphi_\varepsilon[x:y] = [x/\sqrt{\varepsilon}: \sqrt{\varepsilon}y]) = [x^2: \varepsilon y^2: xy] \xrightarrow{\varepsilon \rightarrow 0} [x:0:y]$) but not both.

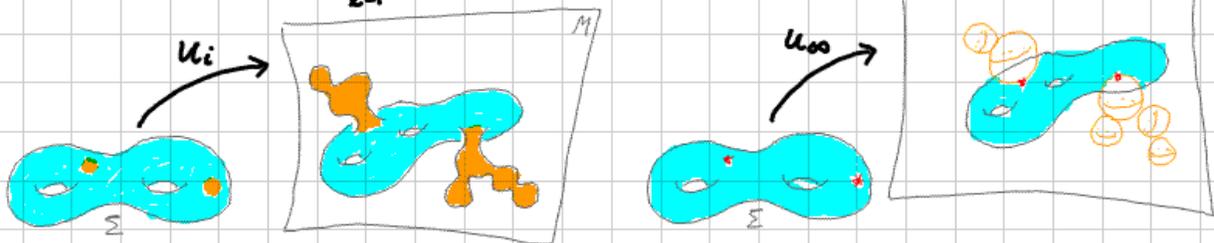
Gromov compactness: Any sequence $u_i: \Sigma \rightarrow M, \bar{\partial}_{J_i} u_i = 0$
! j on Σ fixed!

with fixed energy $\sup_{i \in \mathbb{N}} E(u_i) = \int_\Sigma u_i^* \omega = E$ and $J_i \xrightarrow{e \rightarrow \infty} J_\infty$

has a subsequence that "Gromov-converges" to a J_∞ -hol. map $u_\infty: \Sigma \rightarrow M$

and "stable trees of sphere bubbles" $V_1, \dots, V_L: S^2 \rightarrow M$ J_∞ -hol.

$$\text{s.t. } \int u_\infty^* \omega + \sum_{\ell=1}^L \int V_\ell^* \omega = E.$$



More precisely: For a subsequence (again denoted $(u_i)_{i \in \mathbb{N}}$)

① The energy densities $|du_i|^2: \Sigma \rightarrow \mathbb{R}$ converge as measure

$$e_i = |du_i|^2 \xrightarrow{i \rightarrow \infty} e_\infty + \sum_{k=1}^N \frac{E_k}{V_k} \delta_{p_k} \quad \text{for } p_1, \dots, p_N \in \Sigma$$

$e_\infty \in \mathcal{C}^0(\Sigma, [0, \infty))$ $V_k > 0$

② For every blow-up point $p \in \Sigma$ (i.e. $\exists p_i \xrightarrow{i \rightarrow \infty} p, |du_i(p_i)| \rightarrow \infty$)
 in local coordinates $0 \in \Omega \subset \mathbb{C} \quad \exists z_i \rightarrow 0, |du_i(z_i)| = R_i \rightarrow \infty$

s.t. the rescaled maps $v_i(z) = u_i(z_i + z/R_i) : B_{\varepsilon_i R_i} \rightarrow M$
 $\mathbb{C} \cap B_{\varepsilon_i R_i} \subset \Omega$

converge to a "sphere" $v_\infty : \mathbb{C} \rightarrow M$

(by removal of singularity) $\tilde{v}_\infty : \mathbb{C}P^1 = (S^2) \rightarrow M$ nonconstant

(Note this does not yet capture the full energy / full tree of spheres.)

Cor: For nonsqueezing, the moduli spaces of J-hol. $u : S^2 \rightarrow S^2 \times T$

with energy $E(u) = \int_{S^2} u^* \omega = \langle [S^2 \times pt], [\omega_0^2 \times K \omega_0^T] \rangle = \pi$ are compact

modulo $\text{Aut}(S^2)$.

→ any bubble $\tilde{v}_\infty : S^2 \rightarrow S^2 \times T$ has energy $> 0, \in \pi \mathbb{Z}$
 $\Rightarrow \geq \pi \quad \langle \pi_2(S^2 \times T), [\omega_0^2 \times K \omega_0^T] \rangle$

⇒ $u_i \xrightarrow{e^{\alpha_i}} u_\infty$ or $u_i \xrightarrow{e^{\alpha_i}} u_\infty \equiv \text{const}$ & one bubble \tilde{v}_∞

rescaling gives $u_i \circ \varphi_i \xrightarrow{e^{\alpha_i}} \tilde{v}_\infty$

Tool kit for bubbling analysis

- **Removable singularity**: $u: B \setminus \{0\} = \{z \in \mathbb{C} \mid 0 < |z| < 1\} \rightarrow M$

$$\bar{\partial}_j u = 0, \int_{B \setminus \{0\}} |du|^2 < \infty \Rightarrow \lim_{z \rightarrow 0} u(z) =: \tilde{u}(0) \text{ exists}$$

and extends u to $\tilde{u} \in C^\infty(B, M)$, $\bar{\partial}_j \tilde{u} = 0$.

- **Hofer trick**: (X, d) complete metric space, $f \in C^0(X, [0, \infty))$

$$\forall p \in X, \varepsilon > 0 \exists z \in B_{2\varepsilon}(p), 0 < \varepsilon' \leq \varepsilon: \sup_{B_{\varepsilon'}(z)} f \leq 2f(z), \varepsilon' f(z) \geq \varepsilon f(p)$$

- **mean value inequality**: $e \in C^2(\mathbb{R}^2, [0, \infty))$

$$\left. \begin{array}{l} \Delta e \geq -A_0 - ae^2 \\ \int_{B_r(0)} e < h \\ \quad \quad \quad "h(a) > 0 \end{array} \right\} \Rightarrow r^2 e(0) \leq C \left(\int_{B_r(0)} e + A_0 r^4 \right)$$

- $\partial_s u + J(u) \partial_t u = 0 \Rightarrow e = |du|^2 = 2|\partial_s u|^2$ satisfies

$$(\partial_s^2 + \partial_t^2) e = 4|\partial_s \partial_s u|^2 + 4|\partial_t \partial_s u|^2 + 4g_j(\Delta \partial_s u, \partial_s u)$$

$$\geq \dots \geq -4a|\partial_s u|^4 = -ae^2$$

