

**Workshop on Spectral Theory and Harmonic
Analysis, ANU, Canberra**

Solitary wave dynamics in a slowly varying potential.

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UC Berkeley

16 July 2009

In this talk we will study the **focusing** nonlinear Schrödinger equation (NLS)

$$i\partial_t u = -\frac{1}{2}\Delta u - F(|u|^2)u,$$

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In both examples above, the nonlinear term acts as an **attractive effective potential**, which causes the focusing effect.

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For the remainder of the talk, for definiteness, we will focus on the [Hartree](#) case, although many analogous results are true for the Gross-Pitaevskii and for other nonlinearities.

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From this **soliton state** we have a family of solutions

$$e^{ix \cdot v} e^{i|v|^2 t/2} e^{i\gamma} e^{i\lambda t} \mu^2 \eta(\mu(x - a - vt))$$

to

$$i\partial_t u = -\frac{1}{2}\Delta u - (|x|^{-1} * |u|^2)u.$$

Here $(a, v, \gamma, \mu) \in \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{R} \times \mathbb{R}^+$.

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Notice the center of mass obeys a **classical equation of motion**.

In this talk we will see how adding an **external potential**
 $V: \mathbb{R}^3 \rightarrow \mathbb{R}$ can give a richer example of the same phenomenon:

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Compare this to

$$u(x, t) = e^{ix \cdot v} e^{i|v|^2 t/2} e^{i\gamma} e^{i\lambda t} \mu^2 \eta(\mu(x - a - vt))$$

in the case $V \equiv 0$.

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In a recent preprint available on [arXiv](#), Ivan Ventura and I prove the following

Theorem. Let $V(x) = W(hx)$ with $W \in C^3$. Fix $0 < c_1$ and $(v_0, a_0) \in \mathbb{R}^3 \times \mathbb{R}^3$. Suppose u solves

$$i\partial_t u = -\frac{1}{2}\Delta u - (|x|^{-1} * |u|^2)u + V(x)u, \quad u(0) = e^{iv_0(x-a_0)}\eta(x-a_0).$$

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Then for

$$0 \leq t \leq \frac{c_1}{h} + \frac{\delta |\log h|}{c_2 h}, \quad , 0 \leq \delta < 1/2 \quad 0 < h \leq h_0,$$

we have

$$\|u(x, t) - e^{iv(t)(x-a(t))} e^{i\gamma(t)} \eta(x-a(t))\|_{H^1} \leq c_2 h^{2-\delta}.$$

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(We will see presently that the natural unit time scale in this problem is h^{-1} .)

In the estimate

$$\|u(x, t) - e^{iv(t)(x-a(t))} e^{i\gamma(t)} \eta(x - a(t))\|_{H^1} \leq c_2 h^{2-\delta},$$

the functions $(a(t), v(t), \gamma(t))$ solve

$$\dot{a} = v, \quad \dot{v} = -\frac{1}{2} \int \nabla V(x + a) \eta^2(x) dx,$$

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Moreover,

$$-\frac{1}{2} \int \nabla V(x + a) \eta^2(x) dx = -\nabla V(a) + \mathcal{O}(h^3).$$

In summary, up to time $t \sim h^{-1} + \delta |\log h| h^{-1}$, we have

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Remark. The same result is valid if we take

$$\|u(0) - e^{iv_0(x-a)} \eta(x - a_0)\|_{H^1} \leq c_1 h^2.$$

We also have a slightly weaker result in the case

$$\|u(0) - e^{iv_0(x-a)} \eta(x - a_0)\|_{H^1} \leq c_1 h^{1/2+\delta_0}.$$

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- More precise equations of motion. Where previous authors had

$$\dot{a} = v + \mathcal{O}(h^2), \quad \dot{v} = -\nabla V(a) + \mathcal{O}(h^2),$$

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- More precise error estimates. Where previous authors had

$$\|u(x, t) - e^{iv(t)(x-a(t))} e^{i\gamma(t)} \eta(x - a(t))\|_{H^1} \leq c_2 h^{1-\delta},$$

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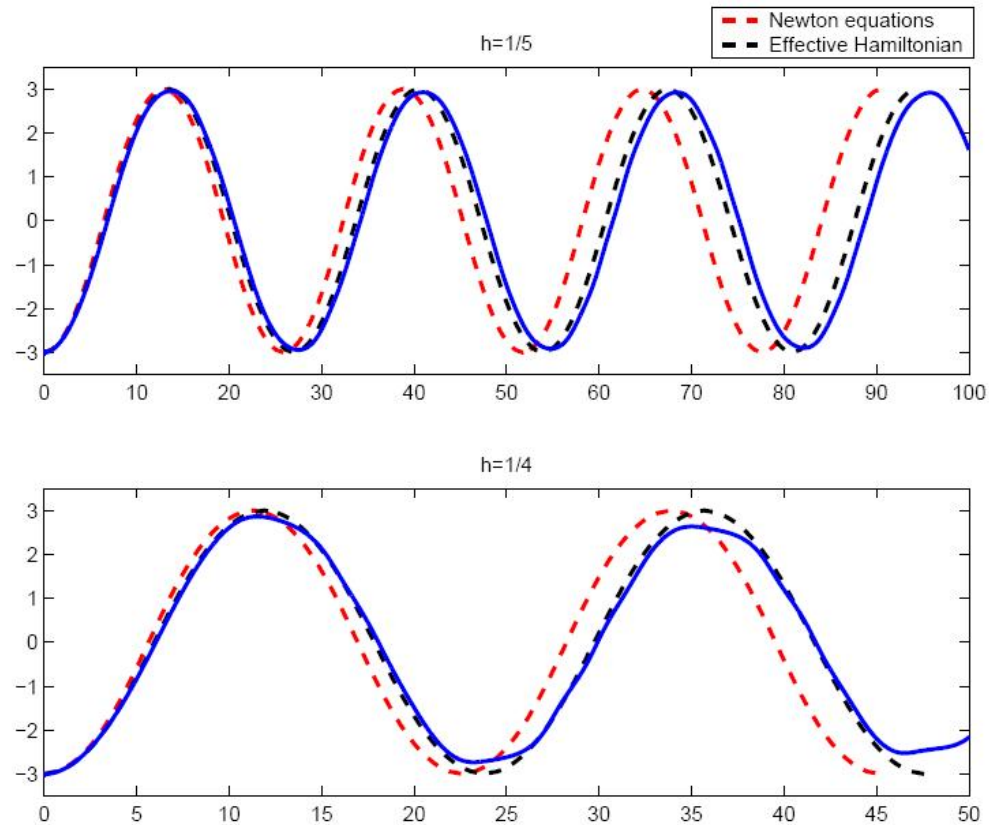
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is the **Hamiltonian flow** corresponding to the Hamiltonian

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We have used here

$$\omega(v, \Xi_{H_V}(u)) = d_u H(v), \quad \partial_t u = \Xi_{H_V}(u).$$

Manifold of solitons.

We define the manifold of solitons $M \subset H^1$ to be the set

$$M = \{e^{iv(x-a)} e^{i\gamma} \mu^2 \eta(\mu(x-a)); (a, v, \gamma, \mu) \in \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{R} \times \mathbb{R}^+\}.$$

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In (a, v, γ, μ) coordinates we have

$$\omega|_M = \mu dv \wedge da + v d\mu \wedge da + d\gamma \wedge d\mu,$$

$$H_V|_M = \frac{|v|^2 \mu}{2} + \lambda \frac{\mu^3}{3} + \frac{\mu^4}{2} \int V(x) \eta^2(\mu(x-a)).$$

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It is from this that we obtain the equations of motion

$$\dot{a} = v, \quad \dot{v} = -\frac{1}{2} \int \nabla V(x+a) \eta^2(x) dx, \quad \dot{\gamma} = \dots$$

In summary, in order to determine the equations of motion for (a, v, γ) in

$$\|u(x, t) - e^{iv(t)(x-a(t))} e^{i\gamma(t)} \eta(x - a(t))\|_{H^1} \leq c_2 h^{2-\delta},$$

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- We observe that u evolves according to a **Hamiltonian flow**

$$\partial_t u = -i \left(-\frac{1}{2} \Delta u - (|x|^{-1} * |u|^2)u + V(x)u \right) = \Xi_{H_V}(u)$$

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(μ ends up varying so little that we can neglect it).

Ideas from the proof.

Using the [implicit function theorem](#), we write (initially for small time)

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It then remains to prove that w is small.

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To do this we employ a variant of the **Lyapounov functional**, used by **Michael Weinstein** (1986):

$$\|w\|_{H_x^1}^2 \leq c \operatorname{Re} \int \mathcal{L}w \bar{w} dx, \quad (2)$$

where \mathcal{L} is the linearization of the Hartree differential operator.

$$\mathcal{L}w \stackrel{\text{def}}{=} -\frac{1}{2}\Delta w - \left(\frac{1}{|x|} * \eta(w + \bar{w}) \right) \eta - \left(\frac{1}{|x|} * \eta^2 \right) w + \lambda w.$$

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That the estimate (2) is valid for w symplectically orthogonal to M is a result of **Enno Lenzmann** (2008).

We build most closely off of work of Holmer-Zworski (2007), who prove a similar theorem for the cubic NLS in 1 dimension:

$$i\partial_t u = -\frac{1}{2}\partial_x^2 u - |u|^2 u + V(x)u.$$

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- We study more general initial data. Where H-Z required

$$\|u(0) - e^{iv_0(x-a)}\eta(x - a_0)\|_{H^1} \leq c_1 h^{3/2+\delta_0},$$

we only require

$$\|u(0) - e^{iv_0(x-a)}\eta(x - a_0)\|_{H^1} \leq c_1 h^{1/2+\delta_0}.$$

Extensions to other nonlinearities.

This phenomenon can also be observed for other nonlinearities: [Fröhlich-Tsai-Yau \(2002\)](#), [Fröhlich-Gustafson-Jonsson-Sigal \(2004\)](#), and [Abou-Salem \(2008\)](#) study general equations of the form

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This assumption is the most delicate, and poses the most substantial barrier to the further extension of these results.

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